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THE UNIVERSITY OF ALBERTA

APPLICATION OF THE FINITE ELEMENT PROGRAM 'ADINAT' TO ONE  
DIMENSIONAL CONSOLIDATION



by

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A THESIS

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This thesis presents a study of the application of the finite element program ABAQUS to the analysis of one-dimensional consolidation. A few basic features of the program are reviewed, and linear and nonlinear consolidation analyses are discussed in detail. Finally, a simple general governing equation is derived for the primary use of predicting consolidation pressures.

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Applications of ABAQUS to pore pressure predictions are illustrated through various sample analyses. Illustrations have included a comprehensive review of relevant researches to date. Within the scope of the present study, this dissertation may serve as a state of the art manual for a general treatment of one-dimensional consolidation.

TO MY PARENTS





## ABSTRACT

This thesis presents a study of the application of the finite element program ADINAT (Automatic Dynamic Incremental Nonlinear Analysis of Temperatures) to one dimensional consolidation. A few basic features of the program are reviewed, and linear and nonlinear consolidation analyses are discussed in detail. Specifically, a simple general governing equation is formulated for the primary use of predicting consolidation pore pressures.

Applications of ADINAT in pore pressure predictions are illustrated through various sample analyses. Illustrations have included a comprehensive review of relevant researches to date. Within the scope of the present study, this dissertation may serve as a state of the art manual for a general treatment of one dimensional consolidation.





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## LIST OF SYMBOLS

a	constant; Lagrangian coordinate, measure in the direction of gravity
b	reduced coordinate, measure in the direction of gravity
c	energy capacity; compressibility; reduced coordinate
c'	mode of variation of the modulus of elasticity
c''	specific heat per unit volume
d	distance from water table to soil layer
e	void ratio; superscript indicating "element"
f	subscript for final value; bold type indicates load vector
g	consolidation parameter
$\bar{g}$	load vector
h	total depth of compressible layer
i	subscript refers to nodal value or counter
k	conductivity; permeability; bold type indicates vector
k <sub>H</sub>	thermal conductivity
m	subscript denotes number of elements





$\hat{m}$	mode of variation for C F
$m_v$	coefficient of volume compressibility
$n$	a constant; an outward normal to boundary surface; subscript refers to values obtained with zero void ratio
$o$	subscript refers to the initial state
$p$	pore water pressure; material constant
$q$	energy flow rate per unit volume; fluid flow; applied loading; material constant
$q^B$	rate of energy generated per unit volume
$q^c$	rate of energy stored per unit volume
$q^s$	boundary energy flow input to the surface $S_2$
$r$	radius
$s$	subscript for value at soil surface
$t$	time
$u$	excess pore pressure
$v$	volume flow velocity; subscript indicates vertical direction
$w$	subscript indicates water
$x$	coordinate variable in ADINAT



$z$	Eulerian coordinate, measure in the direction of gravity
$A$	a constant; cross-sectional area; bold type for interpolation matrix
$A$	amplitude of sine function
$B$	Terzaghi's excess pore pressure ratio; bold type for gradient interpolation matrix
$C$	mass or compressibility matrix
$C_0$	material constant
$C_c$	slope of the $e$ - $\log_{10} \sigma'$ line
$C_F$	consolidation parameter, $(1+e_0/1+e)^2 C_v$
$C'_F$	material parameter, $(\lambda e' + 1)/(\lambda + 1)$
$C_k$	slope of the $e$ - $\log_{10} k$ line
$C_k$	material constant
$C_v$	coefficient of consolidation
$C_v^m$	$C_v$ with constant $m$
$C_{vi}$	initial consolidation ratio, $k_0 \sigma'_0 \ln 10 (1+e_0)/\gamma_w C_c$
$C_{vo}$	consolidation ratio, $k_0 (1+e_0) (d\sigma'/de)_0 / \gamma_w$
$E$	modulus of elasticity of soil
$F$	given function; operator





G	given function; operator
H	initial length of vertical drainage path; subscript refers to heat
K	symmetric matrix; conductivity matrix
K	permeability variable in ADINAT
KM	refers to either the permeability or the compressibility
L	differential operator
M	dimensionless number
M	compressibility variable in ADINAT
N	shape function; total number of solution time step
$Q^g$	rate of energy generated
$Q^i$	concentrated energy flow input
$Q^{in}$	energy inflow rate
$Q^{out}$	energy outflow rate
$Q^s$	rate of energy stored
$Q$	total nodal energy input vector
$Q_B$	energy generated at the system
$Q_c$	energy stored at the system



$Q_s$	energy transfer at the boundary surface
$Q_E$	externally applied nodal energy flows
$S$	boundary surface of the problem domain; settlement
$S_1$	boundary on which the unknown is prescribed
$S_2$	boundary on which the unknown gradient is prescribed
$T$	temperature; time factor; superscript denotes transpose of matrix
$T_b$	time factor for bottom layer
$T'$	time factor, $C_0 t/h$
$T_1$	time factor, $(c'k_{\gamma,t})/\gamma_w$
$U$	average degree of consolidation for the whole layer
$W$	weight
$\alpha$	load ratio; space coordinate; time integration scheme
$\beta$	permeability ratio; space coordinate
$\gamma$	unit weight
$\overline{\gamma}_b$	average submerged unit weight of soil
$\gamma_s$	unit weight of soil solids
$\delta$	variational operator





$\epsilon$	mode of variation of material properties; indicates 'an element of'
$\xi$	self-weight ratio
$\bar{\xi}$	average self-weight ratio
$\eta$	mode of variation of material properties
$\eta'$	material constant
$\theta$	unknown function
$\theta_1$	prescribed unknown function
$\theta^i$	unknown function at boundary point
$\theta^s$	unknown function at the boundary surface
$\bar{\theta}$	vector of the unknown function
$\lambda'$	eigenvalue; material parameter, $\hat{m}e_0/(C_0 - \hat{m}e_0)$
$\lambda'$	material parameter
$\mu$	local consolidation ratio or relative degree of compression
$\nu$	ratio of $C_c/C_k$
$\nu_r$	material constant, $1/C_c - 1/C_k$
$\xi$	strain; convective coordinate
$\rho$	density



$\sigma$	total stress
$\sigma'$	effective stress
$\sigma'_0$	initial effective stress in soil, before loading is applied
$\tau$	dimensionless time factor, $C_v i \cdot t / H^2$
$\phi$	subscript for initial value at surface (before loading is applied)
$\chi$	diffusivity
$\omega$	variable of $\log_{10}(\sigma' / \sigma'_f)$
$\varsigma$	material parameter
$\Delta$	prefix indicates a finite increment
$\Theta$	space coordinte denotes either $a$ or $z$
$\Pi$	a stationary functional
$\Sigma$	summation
$\Omega$	bounded problem domain
$\S$	characteristic number
$\ell$	reduced thickness of soil layer
$i$	subscript or superscript refers to the $i$ th layer
$1$	subscript refers to an arbitrary reference state



- ∞ indicates the final value at layer bottom
- ' superscript prime denotes coordinate differentiation
- ' quote mark refers to a ratio of the original value
- superposed dot denotes time differentiation
- upper underscore indicates mean value or trial function





## 1. INTRODUCTION

### 1.1 General

The majority of phenomena arising in engineering and physics are often described by differential equations, especially by partial differential equations, and by boundary conditions imposed on the unknown functions.

Seepage through porous media, heat conduction and torsion are three important phenomena governed essentially by differential equations of similar form, particularly the diffusion, Laplace and Poisson equations. These problems are often known as *field problems*, and the general governing equation is called a *field governing equation*.

Some other examples of this type of problem include gravitational and irrotational fluid flow, surface waves on a fluid, natural frequency of footings, heat convection, electrostatics, magnetostatics and electrical currents. (Desai and Abel, 1972)

ADINAT, a general purpose computer program which employs the parabolic differential equation of the heat conduction type, can be used to solve many of these field problems.



Since finding closed-form solutions for the more general cases can be difficult, various methods of computing approximate solutions have been developed. Among these, the energy method, or variational method, has gained considerable popularity. The finite element formulation, which is modified from the classical variational method, became widely used with the advent of computers.

ADINAT employs the variational procedure in deriving the finite element equations. The high accuracy of the program and its incremental solution procedure are particularly applicable to problems with soils, where nonlinear constitutive relations are fundamental and intrinsic.

Consolidation is the process of gradual reduction in volume of a soil mass of low permeability due to the dissipation of excess pore water pressure. The one dimensional consolidation condition occurs either with a large uniformly loaded area or with lateral physical constraints. One dimensional consolidation theory is extensively used in practice for pore pressure predictions, settlement and stability analyses.

The conventional consolidation analysis is based on the classical Terzaghi theory. Despite the simplifying assumptions embodied in the theory, it is still widely





employed in practice due to its simplicity. It also permits extensions in a simple way to cover other situations.

An example of an earlier extension is the nonhomogeneous problem. It is often approximated as discrete layers of unlike soils, in each of which the material properties are reasonably uniform. Continuous spatial variation of material properties, is another simple method to simulate the varying material properties in soils. The latter is, however, a step closer to a *nonlinear* analysis.

The term 'nonlinear' here refers to the condition of nonlinear constitutive relationships, or of material properties which basically vary with time.

One dimensional nonlinear consolidation studies were developed by 1965. Barden and Berry (1965), Davis and Raymond (1965), Janbu (1965), Mikasa (1965), Raymond (1966 & 1969), Gibson *et al* (1967), Poskitt (1969), Davis (1971) provided prominent nonlinear theories at that time. Gibson *et al* published a general theory removing many of the simplifying assumptions in the conventional analysis, and included the effects of both finite strain and self-weight. Extensions to this theory have been made by other researchers. (Poskitt, 1969; Simons and Beng, 1969; Monte and Krizek, 1976; DeSimone and Viggiani, 1976; Gibson *et al*,



1981)

Due to mathematical difficulties, recourse has been made by most nonlinear theories to further approximations in order to derive solutions.

Finite element solutions are provided by ADINAT for various nonlinear analyses. Detailed discussions of the application of the program to nonlinear consolidation will be presented in this thesis.

## 1.2 Scope of Research

The basic aim of this thesis is a general study of ADINAT application in one dimensional consolidation. Necessary steps to achieve this goal are outlined below:

1. Examination of basic features in ADINAT;
2. Thorough review of consolidation theories;
3. Assessment of available models applicable to ADINAT;
4. Formulation of the appropriate model compatible with the illustrative theory.
5. Simulation of ADINAT input and preparation for the results.
6. Evaluation of the use of ADINAT in such analysis.



Steps 3 to 6 are repeatedly used throughout this thesis.

The consolidation studies here are limited by the assumptions stated in the following chapters. The studies can be divided into two parts. The first part deals with linear and nonhomogenous problems, in which Terzaghi's simplifications are generally applied. The second part is concerned with nonlinear analyses, where assumptions such as the validity of Darcy's law and constant load increments are accepted but stress history and secondary effects are neglected.

Applications of ADINAT are illustrated through sample analyses selected from various areas of the program applicability. Predictions of pore pressures are the primary concern in these examples. The guidelines given below are followed in the development of the examples:

- a. Simplicity and brevity;
- b. Usefulness, in the sense that each example represents a certain type of formulation;
- c. Correctness in modelling, by presenting the required derivation and a comparison with published formulations and/or experimental results;
- d. Correctness in computation, by comparison with closed-form or published results where available;





- e. Completeness; the main reasons for those not illustrated are: similarity, duplication, inadequate data, no practical solution exists or the analyses are outside the scope of the assumptions.

A brief discussion of the governing equation and several features in ADINAT are given in Chapter 2. Chapter 3 deals with the application of ADINAT to linear problems. Conventional analysis, as well as homogenous layered and continuous nonhomogeneous problems, are considered.

In view of the variety of soil stratigraphy, loading and boundary conditions that exist in nature, uses of ADINAT in different physical situations will also be illustrated and given with the examples.

A brief review of theories and discussions on the use of ADINAT in nonlinear analyses are presented in Chapter 4. A collection of equations are given for a general treatment of these analyses. These equations are tabulated for convenience in use and reference. Chapter 5 and 6 are devoted to examples of various published nonlinear theories.



## 2. BRIEF DESCRIPTION OF ADINAT

The governing equation in ADINAT will be discussed below. Mass matrix or compressibility matrix analysis and the time integration scheme of the program will be briefly discussed. Stability and accuracy requirements in using the program are also given.

### 2.1 Governing Field Equation

The general field equation solved by ADINAT, for a problem domain idealized as an assemblage of finite elements, can be interpreted physically as a mathematical statement of the energy equilibrium at the nodes of the system at any time, that is (Wilson, Bathe and Peterson, 1974; Holman, 1976)

$$Q^{in} + Q^g = Q^{out} + Q^s \quad (2.1)$$

in which

$Q^{in}$  : rate of energy flow into the elements adjacent to the node and into the node from an external source,



- $Q^g$  : rate of energy generated within the elements adjacent to the node,  
 $Q^{out}$  : rate of energy flow out of the elements adjacent to the node, .  
 $Q^s$  : rate of energy stored within the elements adjacent to the node.

Let  $\theta$  denote the unknown function or the state variable, which can be fluid head or potential, stream function, velocity potential, temperature, warping, stress function, electrical potential and so on. The energy flux,  $q_z dz$ , is generally related to the gradient of  $\theta$ . In considering the one dimensional condition, a very general relationship in terms of the proportionality constant, termed conductivity  $k_z$ , is:

$$q_z dz = v_z \rho_z = -k_z \frac{\partial \theta}{\partial z} \quad (2.2)$$

where  $v_z$  is the flow velocity and  $\rho_z$  is the density. This can also be expressed by similar equations as those shown below using the appropriate physical law for a given problem. For instance, from Darcy's law





$$q_z dz = v_z \gamma_w = -k_z \frac{\partial u}{\partial z} \quad (2.3)$$

or from Fourier's law of heat conduction

$$q_z dz = -k_z \frac{\partial T}{\partial z} \quad (2.4)$$

where  $u$  is the fluid pressure,  $T$  is the temperature,  $k_z$  is the hydraulic conductivity in Eq. 2.3 and the thermal conductivity in Eq. 2.4. Hence,  $Q^{in}$  can be determined by

$$Q^{in} = -k_z A_z \frac{\partial \theta}{\partial z} \quad (2.5)$$

Similarly,

$$Q^{out} = -k_z A_z \frac{\partial \theta}{\partial z} \bigg|_{z+dz}$$



$$= -A_z \left[ k_z \frac{\partial \theta}{\partial z} + \frac{\partial}{\partial z} \left( k_z \frac{\partial \theta}{\partial z} \right) dz \right] \quad (2.6)$$

and

$$Q^g = q^B \cdot A_z \cdot dz \quad (2.7)$$

$$Q^s = q^C \cdot A_z \cdot dz \quad (2.8)$$

$q^B$  is the rate of energy generation per unit volume while  $q^C$  is the rate of energy stored per unit volume. Combining the above relations gives,

$$\frac{\partial}{\partial z} \left( k_z \frac{\partial \theta}{\partial z} \right) = q^C - q^B \quad (2.9)$$

Generalizing this equation in three dimensions yields:

$$\frac{\partial}{\partial x} \left( k_x \frac{\partial \theta}{\partial x} \right) + \frac{\partial}{\partial y} \left( k_y \frac{\partial \theta}{\partial y} \right) + \frac{\partial}{\partial z} \left( k_z \frac{\partial \theta}{\partial z} \right) = q^C - q^B \quad (2.10)$$



where  $k_x, k_y, k_z$  are the conductivities corresponding to the  $x, y, z$  directions.

The associated boundary conditions are

$$\theta = \theta_1 \quad (2.11)$$

on part of the boundary  $S_1$  of the domain, while

$$k_n \frac{\partial \theta}{\partial n} = q^s \quad (2.12)$$

on part of the boundary  $S_2$  of the problem domain, and

$$\theta|_{t=0} = \theta_0 \quad (2.13)$$

$n$  denotes the outward normal to the surface.  $\theta_1, q^s$  are the prescribed values of the unknown and energy flow input to the boundary surface respectively.  $\theta_0$  is the initial value of the unknown  $\theta$ .





## 2.2 Variational Method

### 2.2.1 Variational Principle

A *variational principle* specifies a scalar quantity expressed in an integral form (Desai and Abel, 1972; Zienkiewicz, 1977)

$$\Pi = \int_{\Omega} F(x, y, z, \theta, \theta', \theta'', \dots) d\Omega + \int_S G(\theta, \theta', \dots) dS \quad (2.14)$$

in which  $\theta$  is the unknown variable or function, the superscript prime denotes differentiation with respect to the cartesian coordinates, the number of primes is the order of differentiation.  $\Omega$  is the bounded problem domain and  $S$  is its piecewise smooth boundary,  $F$  and  $G$  are specified operators.  $\Pi$  is called a functional, because it is a function of other functions.

The boundary conditions can be sorted into three classes: the essential, natural and mixed boundary conditions. For quadratic functionals, where the operators contain at most second order derivatives, essential boundary conditions are the prescribed values of the unknown function  $\theta$ . The natural boundary conditions correspond to the first order derivatives of the unknown function,  $\partial\theta/\partial n$ .  $n$  denotes



one of the coordinate direction  $x, y, z$ , or the direction of the outward normal to the boundary surface. The mixed boundary condition is a combination of the above two.

In the variational procedure, tentative solutions, or trial functions,  $\bar{\theta}$ , satisfying the essential boundary conditions are tried for a given problem,

$$\theta \simeq \bar{\theta} = \sum N_i a_i \quad (2.15)$$

where  $N_i$  are base functions prescribed in terms of independent variables such as the coordinates  $x, y, z$ . Parameters  $a_i$  are unknown constants.

From all such possible solutions, the solution to the problem will be the one which makes the functional  $\Pi$  stationary, that is

$$\delta \Pi = 0 \quad (2.16)$$

with respect to small changes  $\delta \theta$ , where



$$\bar{\theta} = \theta + \delta\theta \quad (2.17)$$

$\delta$  is the variation notation, termed the 'variation in'.

From calculus of variations,

$$\delta\Pi = \frac{\partial\Pi}{\partial a_1} \delta a_1 + \frac{\partial\Pi}{\partial a_2} \delta a_2 + \dots = \frac{\partial\Pi}{\partial \mathbf{a}} \delta \mathbf{a} \quad (2.18)$$

Since the stationary condition of  $\Pi$  in Eq. 2.16 is true for any arbitrary  $\delta \mathbf{a}$ , then

$$\frac{\partial\Pi}{\partial \mathbf{a}} = 0 \quad (2.19)$$

which yields a set of equations for the parameters  $\mathbf{a}$ . Note that if the functional  $\Pi$  is quadratic, or the derivatives of the unknown function  $\theta$  do not exceed second order, Eq. 2.19 reduces to a set of linear equilibrium equations of the form (Zienkiewicz, 1977)





$$K\mathbf{a} + \mathbf{f} = 0 \quad (2.20)$$

$\mathbf{f}$  is the action vector, and the matrix  $K$  will always be symmetric. Hence the parameters  $a_i$  are easily computed, and an approximate solution for the problem can be found.

The essence of the variational method is, therefore, to establish the functional  $\Pi$  of the problem, choose a trial function  $\bar{\theta}$  and find the stationary value of  $\Pi$ , that is  $\delta\Pi = 0$ .

This approximation process using trial functions and stationary conditions is known as the Rayleigh-Ritz method. The finite element form of this method employs several simple trial functions  $\bar{\theta}^e$ , each defined only over an element of a discretized domain, instead of using a single trial function which spans the entire domain as in the classical method. This piecewise fit is the basic feature of finite element analysis.

The appropriate functional for a given problem can be established in a number of ways. It can be derived vigorously using some theorems as guidelines, or obtained backward from the governing differential equations and natural boundary conditions with mathematical manipulation. Generally, there does not exist a unique functional for a



particular problem. A specific functional, however, can be applicable to a certain class of problems. (Bathe, 1982)

### 2.2.2 General Formulation

A linear system of differential equations can be written as

$$L\bar{\theta} + \bar{g} = 0 \quad (2.21)$$

where  $\bar{\theta}$  is the vector of the unknown variable, and  $\bar{g}$  is the known action function. If the differential operator  $L$  is linear and self-adjoint, it can be shown that the variation principle is given as (Zienkiewicz, 1977)

$$\Pi = \int_{\Omega} [(1/2)\bar{\theta}^T L\bar{\theta} + \bar{\theta}^T \bar{g}]d\Omega + \text{Boundary Terms} \quad (2.22)$$

the superscript  $T$  stands for the transpose.



Applying this to Eq. 2.10 and using integration by parts, the functional of the field equation is then (Bathe, 1982)

$$\begin{aligned} \Pi = & \int_{\Omega} \frac{1}{2} \left[ k_x \left( \frac{\partial \theta}{\partial x} \right)^2 + k_y \left( \frac{\partial \theta}{\partial y} \right)^2 + k_z \left( \frac{\partial \theta}{\partial z} \right)^2 \right] d\Omega \\ & - \int_{\Omega} \theta (q^B - q^C) d\Omega - \int_{S_2} \theta^s q^s dS - \sum_i \theta^i Q^i \end{aligned} \quad (2.23)$$

where  $\theta^s$ ,  $\theta^i$  are the unknown functions at the boundary surface.  $Q^i$  is the concentrated energy flow input.

The functionals in engineering problems usually possess physical meaning. (Zienkiewicz *et al*, 1966) The first integral of the above functional, may be interpreted as the rate at which energy is being dissipated over the entire domain, where,

$$-\left[ k_x \left( \frac{\partial \theta}{\partial x} \right)^2 + k_y \left( \frac{\partial \theta}{\partial y} \right)^2 + k_z \left( \frac{\partial \theta}{\partial z} \right)^2 \right] = v_x \frac{\partial \theta}{\partial x} + v_y \frac{\partial \theta}{\partial y} + v_z \frac{\partial \theta}{\partial z} \quad (2.24)$$

is proportional to the energy dissipation rate in a unit volume. The fact that the unknown  $\theta$  system gives the





minimum value of this integral which corresponds to the exact solution, is consistent with the principle of minimum potential energy.

From Eq. 2.16 and Eq. 2.23, and since  $\theta$  is the only variable, it can be shown that (Bathe, 1982)

$$\begin{aligned} \int_{\Omega} \delta(\bar{\theta}')^T \mathbf{k} \bar{\theta}' d\Omega \\ = \int_{\Omega} \delta\theta (q^B - q^C) d\Omega + \int_{S_2} \delta\theta^s q^s dS + \sum_i \delta\theta^i Q^i \end{aligned} \quad (2.25)$$

in which

$$(\bar{\theta}')^T = \left[ \frac{\partial \theta}{\partial x} \quad \frac{\partial \theta}{\partial y} \quad \frac{\partial \theta}{\partial z} \right] \quad (2.26)$$

$$\mathbf{k} = \begin{bmatrix} k_x & 0 & 0 \\ 0 & k_y & 0 \\ 0 & 0 & k_z \end{bmatrix} \quad (2.27)$$

Consider the whole problem domain, for any element



$$\theta = \mathbf{A} \bar{\theta} \quad (2.28)$$

$$\bar{\theta}' = \mathbf{B} \bar{\theta} \quad (2.29)$$

where  $\mathbf{A}$  and  $\mathbf{B}$  are the element  $\theta$  and  $\theta$ -gradient interpolation matrices respectively. Assume the rate of energy storage is given by

$$\dot{q}^c = c \dot{\theta} \quad (2.30)$$

the proportional constant  $c$  is called energy capacity (or volume compressibility). A superposed dot indicates time differentiation. The variation in  $\theta$ ,  $\delta\theta$ , in Eq. 2.25 can be considered as a virtual quantity and the equation is analogous to the virtual work equation in structural stress analysis. Hence the system governing energy equilibrium equations can be written as (Bathe, 1982)

$$\mathbf{C} \dot{\bar{\theta}} + \mathbf{K} \bar{\theta} = \mathbf{Q} \quad (2.31)$$

in which  $\mathbf{C}$  is the finite element system capacity matrix,



$$C = \sum_m \int_{\Omega} \mathbf{A}^T c \mathbf{A} d\Omega \quad (2.32)$$

$m$  is the number of elements in the discretized problem domain.  $K$  is the conductivity matrix of the system,

$$K = \sum_m \int_{\Omega} \mathbf{B}^T k \mathbf{B} d\Omega \quad (2.33)$$

$Q$  is the total nodal point energy input,

$$Q = Q_B - Q_c + Q_s + Q_E \quad (2.34)$$

where

$$Q_B = \sum_m \int_{\Omega} \mathbf{A}^T b \mathbf{A} d\Omega \quad (2.35)$$

$$Q_c = \sum_m \int_{\Omega} \mathbf{A}^T c \mathbf{A} d\Omega \quad (2.36)$$

$$Q_s = \sum_m \int_{S_2} \mathbf{A}^T q^s dS \quad (2.37)$$

$Q_E$  is a vector of externally applied nodal energy flows.  $C$  and  $K$  are constant matrices in linear steady analysis.



For nonlinear or transient problems,  $C$  and  $K$  may be either  $\theta$  or time dependent or both, as

$$C(\theta, t) \cdot \dot{\bar{\theta}} + K(\theta, t) \cdot \bar{\theta} = Q(t) \quad (2.38)$$

where  $t$  denotes time.

## 2.3 A Note on The Lumped and Consistent Matrix Analyses

Two types of matrix analysis options, the lumped and consistent matrix method, are available in ADINAT. Some of their basic features are briefly described in this section.

### 2.3.1 Consistent Matrix

Application of the standard finite element method to the energy equilibrium equation produces the capacity matrix  $C$  in the form of Eq. 2.32. This matrix is called *consistent* because it is derived from a procedure according to the variational formulation. (Desai, 1979) The resulting matrix is non-diagonal. The term *consistent* is thought to be





unnecessary since it is the natural outcome of the discretization process. (Zienkiewicz, 1977)

The principle advantage of the consistent matrix method is to provide a more rational analysis. However, its use does not always lead to improved accuracy but always involves additional computational work compared to the lumping method. (Desai and Abel, 1972; Zienkiewicz, 1977)

In the frequency analysis using the consistent matrix approach, more accurate mode shapes and upper-bound frequencies are obtained. (Desai and Abel, 1972)

### 2.3.2 Matrix Lumping Approach

An alternative way of establishing an approximate matrix  $C$  is by physical reasoning. As such, the matrix is assembled in an arbitrary manner by assuming that all the  $C$  values tributary to a node is lumped at that node, thus a diagonal matrix always results although no actual concentrated values exist. This matrix is often referred to as lumped matrix. (Zienkiewicz, 1977)

The resulting diagonal matrix in the conventional lumped matrix approach leads to a simple technique of formulation and solution. In addition, the lumped matrix



utilizes less storage space. This method is hence significantly more convenient and economical for many computational processes compared to the consistent matrix approach. (Zienkiewicz, 1977)

But if a coarse finite element mesh is used the result obtained may be very inaccurate. Surprising solutions may also result in some cases when high-order elements are employed. (Bathe, 1982)

The lumping scheme, nevertheless, can often be constructed so that there will be no accuracy loss. Researchers have found some matrix lumping techniques that improve accuracy. (Fried and Malkus, 1975; Hinton, Rock and Zienkiewicz, 1976; Zienkiewicz, 1977)

The consistent and lumped matrix option can be found in ADINAT users manual input section II.2.

## 2.4 Remark on the Time Integration Scheme

### 2.4.1 General

For transient analysis, a one parameter ( $\alpha$ ) family of one-step numerical time integration method is employed in



ADINAT. This method, as in the general incremental procedure, assumes that the unknown functions at time  $t_n$  have been determined and that they are to be computed at time  $t_n + \Delta t$ , where  $t_n$  is the time at time  $n$ ,  $n \in \{0, 1, \dots, N-1\}$ ,  $N$  is the total number of solution time steps and  $\Delta t$  is the time step interval.

To evaluate the unknown functions at time  $t_n + \Delta t$ , the energy equilibrium of the problem domain is considered at time  $t_n + \alpha \Delta t$ . From Eq. 2.38, (Bathe and Khoshgoftaar, 1979; Bathe, 1981)

$$C(\theta_{t_1}, t_1) \dot{\bar{\theta}}_{t_1} + K(\theta_{t_1}, t_1) \bar{\theta}_{t_1} = Q(t_1) \quad (2.39)$$

$$\dot{\bar{\theta}}_{t_1} = (\bar{\theta}_{n+1} - \bar{\theta}_n) / \Delta t \quad (2.40)$$

$$\bar{\theta}_{t_1} = (1 - \alpha) \bar{\theta}_n + \alpha \bar{\theta}_{n+1} \quad (2.41)$$

$$\bar{\theta}|_{t=0} = \bar{\theta}_0 \quad (2.42)$$

In the above equations  $t_1$  denotes  $t_n + \alpha \Delta t$ , or  $(n + \alpha) \cdot \Delta t$ . Subscripts  $n$  and  $n+1$  denote the time of  $t_n$  and  $t_n + \Delta t$  respectively.  $\bar{\theta}_0$  is the prescribed value at the initial state.





Note that  $\alpha$  is a constant that ranges from 0 to 1. The value of  $\alpha$  for a given problem is chosen to obtain optimum stability and accuracy.

Some known integration procedures such as Euler's forward method, the trapezoidal rule, and Euler's backward method correspond to  $\alpha$  values of zero, 1/2, and 1 respectively. These options are shown in ADINAT users manual input section II.5.

#### 2.4.2 Stability

The properties of the integration method depend on the value of  $\alpha$  that is used. They can be analysed by considering a typical single degree-of-freedom equilibrium equation, (Hughes, 1977; Bathe and Khoshgoftaar, 1979; Bathe, 1982)

$$\ddot{\theta} + \lambda(\theta, t) \cdot \theta = 0 \quad (2.43)$$

where  $\lambda$  can be considered as a ratio related to  $k$  and  $c$ . The characteristic equation of the recursion relation is obtained by using similar procedures from Eq. 2.39 to



Eq. 2.42 and solving for the unknown function  $\theta$  at time  $t_n + \Delta t$ ,

$$\begin{aligned}\theta_{n+1} &= \mathbb{S} \cdot \theta_n \\ &= \left[ \frac{1 - (1-\alpha)\Delta t \lambda_{t+1}}{1 + \alpha\Delta t \lambda_{t+1}} \right] \theta_n\end{aligned}\quad (2.44)$$

where for stability, it is required

$$|\mathbb{S}| \leq 1 \quad (2.45)$$

From this, the following results can be observed:

First, for  $1/2 \leq \alpha \leq 1$ , the above relationship is always satisfied for any  $\Delta t$ , which means that the integration method is unconditionally stable.

Second, for  $0 \leq \alpha < 1/2$ , the relation in Eq. 2.45 is satisfied provided

$$\Delta t \leq 2 / \{ (1-2\alpha) \cdot \lambda_{t+1} \} \quad (2.46)$$



that is, the method is conditionally stable.

Therefore, the conditionally stable methods, particularly the explicit scheme ( $\alpha=0$ ), will always result in instability as long as the time interval  $\Delta t$  exceeds that specified by the highest eigenvalue in the problem. The time saving computation in the explicit scheme may, however, compensate the requirement of using many small time steps.

### 2.4.3 Accuracy

The accuracy of the Euler forward method ( $\alpha=0$ ), the trapezoidal rule ( $\alpha=1/2$ ), the Euler backward method ( $\alpha=1$ ) can be assessed using the Taylor series expansion. (Hughes, 1977; Bathe and Khoshgoftaar, 1979; Snyder and Bathe, 1981)

It can be shown that, both of the Euler methods are first-order accurate whereas the trapezoidal rule is second-order accurate. The remaining  $\alpha$  schemes can also be shown to possess first-order accuracy. Hence, the trapezoidal rule is the only second-order accurate method in the  $\alpha$ -family of integration scheme.



The  $\alpha$  values with the largest error would be  $\alpha=0$  and  $\alpha=1$ . In the Euler backward scheme ( $\alpha=1$ ), the truncated terms in the time derivative equation obtained from Taylor's expansion are alternate in sign, which may result in significantly better solution accuracy.

#### 2.4.4 Oscillation Effect

The following important property should also be considered in the evaluation of the different integration schemes in the  $\alpha$ -family method. Refer to Eq. 2.44 in the previous section, oscillatory solutions will result as the characteristic value  $\$$  becomes negative (Zienkiewicz, 1977). Assume that  $[C^{-1}][K]$  of Eq. 2.31 has a complete set of orthogonal eigenvectors with corresponding eigenvalues  $\lambda_i$ , there will be no oscillation effect provided that (Wood and Lewis, 1975)

$$\Delta t < 1/\{(1-\alpha)\lambda_i\} \quad (2.47)$$

for all  $\lambda_i$ , or





$$\Delta t < \Delta t_{\text{critical}} = 1 / \max \lambda_i \quad (2.48)$$

In each integration scheme for which  $\alpha < 1$ , there exists a time interval beyond which  $\xi$  becomes negative or violates Eq. 2.47. Thus the Euler backward method ( $\alpha=1$ ) is the only one not showing this oscillation phenomena. The predicted solutions using this method converge monotonically to the exact solution as the time interval  $\Delta t$  decreases.

Experience shows that, the time interval required for an accurate solution in the Euler forward method is considerably smaller than the value for stability. Even if the trapezoidal rule shows a higher order accuracy mathematically, it can give less accurate results than Euler backward method for larger time intervals. (Bathe and Khoshgoftaar, 1979)

The critical time interval in Eq. 2.48 is related to the number of degrees-of-freedom. As the number of degrees-of-freedom increases due to the increases in the complexity of the problem or the base functions, the critical time interval becomes proportionately smaller. Therefore, it is impractical to eliminate oscillations by reducing time intervals. (Wood and Lewis, 1975)



The procedure with  $\alpha=2/3$ , named the Galerkin scheme, has been shown to be practical as a compromise between accuracy and oscillation effects. Another optimal scheme, with  $\alpha=0.878 \approx 7/8$  was also suggested by Lambert (1973), using different reasoning. (Zienkiewicz, 1977)

Wood and Lewis (1975) have illustrated that the Crank-Nicholson scheme with a simple averaging process is effective in dealing with oscillations. They also observed that parabolic and cubic finite elements give better performance for accuracy and damping oscillations than linear elements, particularly with smaller time intervals.

## 2.5 Discussion

Eq. 2.10 is the basic governing equation in ADINAT. Various field problems can be analysed using the program by manipulating the variables  $k$ ,  $\theta$ ,  $q^B$  and  $q^C$ . These may be consolidation, seepage, and problems in oil sands and cold region engineering, specifically in the area of geotechnical use.

In the present work, illustrations will be given for the case of one dimensional consolidation.  $q^B$  is then set equal zero and Eq. 2.30 is employed. There is no applied



mass flow loadings and hence they will be ignored in the program input.

Notice that the stationary value of Eq. 2.23 not only includes the governing equation but also the natural boundary condition. Hence, for the ADINAT input, only the essential boundary conditions are required to be explicitly specified, since the natural boundary conditions are automatically incorporated in the formulation when using the variational approach. This is particularly effective in the analysis of the problems with complicated boundaries.

Regarding the essential boundary conditions, nodes at impervious boundaries are deleted as degrees of freedom in the input stage. This simply means that the equilibrium equations Eq. 2.31 will not be considered for these nodes. (see ADINAT users manual input section IV)

The neglecting of natural boundary conditions input and the mechanical generation of system equilibrium equations, as shown in Section 2.2, are the primary advantages of the variational approach as used by the program.

The value of  $\alpha$  used in the following examples is 1 (implicit scheme), mainly for reasons of stability. Other choices of  $\alpha$  values may provide better results, depending on the given problems. Examples showing the use of  $\alpha$  values





other than 1 can be found in Part B of the program manual.

The equilibrium iteration scheme used by ADINAT is the modified Newton-Raphson method. The number of time steps between equilibrium iterations, maximum number of iteration cycles and allowable tolerance for convergence are described in the program manual and input in control card II.4.



### 3. LINEAR APPLICATION OF ADINAT IN ONE DIMENSIONAL CONSOLIDATION

The application of ADINAT's linear capability in one dimensional consolidation will be illustrated in this chapter. Examples will include Terzaghi's classical theory, homogeneous layered and continuous nonhomogeneous problems.

#### 3.1 Terzaghi Classical Theory

##### 3.1.1 General

The classical theory of one dimensional consolidation of clay deposits was developed by Karl Terzaghi in 1925. It is founded on the following basic assumptions:

1. The soil is homogeneous and is saturated.
2. The soil grains and pore fluid are incompressible in comparison with the soil skeleton, and volume change is the result of changes in void ratio.
3. The average fluid flow and compression are one dimensional.
4. The surface applied load causes a uniform increment of vertical total stress throughout the soil layer.
5. Darcy's law of permeability is valid at all hydraulic gradients and in the form of the average



fluid flow velocity related to the excess pore pressure gradient.

6. The stress history of the soil stratum is ignored.
7. The structural viscosity effects, which control secondary compression or creep, are ignored.
8. All initial total stresses are uniformly distributed throughout the soil stratum.
9. The void ratio is a unique linear function of the effective stress and is independent of time, that is, the coefficient of volume change is constant.
10. The coefficient of permeability remains constant throughout the consolidation process.
11. Strains, velocities and stress increments are small.
12. The soil layer is comparatively thin such that the stresses arising from the self-weight of the solids and pore fluid are negligible compared with those applied.

By equating the volume of pore water expulsion with the reduction in void volume of a soil element, Terzaghi established the following governing equation for one dimensional consolidation:



$$C_v \frac{\partial^2 u}{\partial z^2} (z, t) = \frac{\partial u}{\partial t} (z, t) \quad (3.1)$$

where  $C_v$  is the coefficient of consolidation,  $u$  is the excess pore pressure,  $z$  is the vertical distance from the nearest drainage boundary to the element under consideration.

Eq. 3.1 defines the immediate magnitude of the excess pore pressure  $u$  in an element at a specific distance  $z$  in the soil mass at time  $t$ , after an instantaneous application of a time-independent load  $\Delta q$ . Since the load causes a uniform stress increment,  $\Delta \sigma_v$ , the initial excess pore pressure is

$$u(z, 0) = \Delta \sigma_v = \Delta q \quad (3.2)$$

Hence the excess pore pressure  $u$  for a clay layer subjected to double drainage can be expressed as (Taylor, 1948)

$$u(z, t) = \sum_{m=0}^{\infty} \frac{2u_0}{M} \sin \frac{Mz}{H} e^{-M^2 T_v} \quad (3.3)$$





where  $M=\pi(2m+1)/2$  and  $m$  is any integer from 0 to  $\infty$ . The boundary conditions are, for double drainage,

$$u(0,t) = 0 \quad (3.4)$$

$$u(h,t) = 0 \quad (3.5)$$

for single drainage, Eq. 3.5 becomes

$$\frac{\partial u}{\partial z}(H,t) = 0 \quad (3.6)$$

where  $H$  is the minimum drainage path. Note that  $C_v$  is generally expressed as,

$$C_v = \frac{k}{\gamma_w m_v} \quad (3.7)$$

$k$  is the coefficient of permeability and  $m_v$  is the coefficient of volume compressibility, which is defined by



$$\frac{m}{v} = - \frac{1}{1+e} \cdot \frac{de}{d\sigma'} \quad (3.8)$$

The term  $(1+e)$  in the above equation is assumed to be equal to  $(1+e_0)$  in this chapter. The average consolidation  $U$  for the entire stratum in linear theory is determined as

$$U = 1 - \frac{\int u(z,t)dz}{\int u(z,0)dz} \quad (3.9)$$

The above integrals are evaluated for the entire compressible layer. Substituting Eq. 3.2 and Eq. 3.3 into Eq. 3.9 gives

$$U = 1 - \sum_{m=0}^{\infty} \frac{2}{M^2} e^{-M^2 Tv} \quad (3.10)$$

The analogies between the heat transfer equation in ADINAT and the consolidation equation can be seen below,



$$\text{Heat Transfer:} \quad k \frac{\partial^2 T}{H \partial z^2} = \rho c \frac{\partial T}{\partial t} \quad (3.11)$$

$$\text{Consolidation:} \quad k \frac{\partial^2 u}{\partial z^2} = \gamma \frac{m}{w v} \frac{\partial u}{\partial t} \quad (3.12)$$

$k$  is the thermal conductivity.  
H

The input and output parameters of ADINAT can be deduced from the above comparison. It is worthwhile to mention that in the program, the product of  $\rho$  and  $c$  forms a single variable  $c''$ , which is the specific heat per unit volume.



### 3.1.2 Sample Analysis

#### Objective:

To illustrate the application of ADINAT to Terzaghi's one dimensional consolidation theory and predict the pore pressure distribution in a simple soil layer.

#### Problem Description:

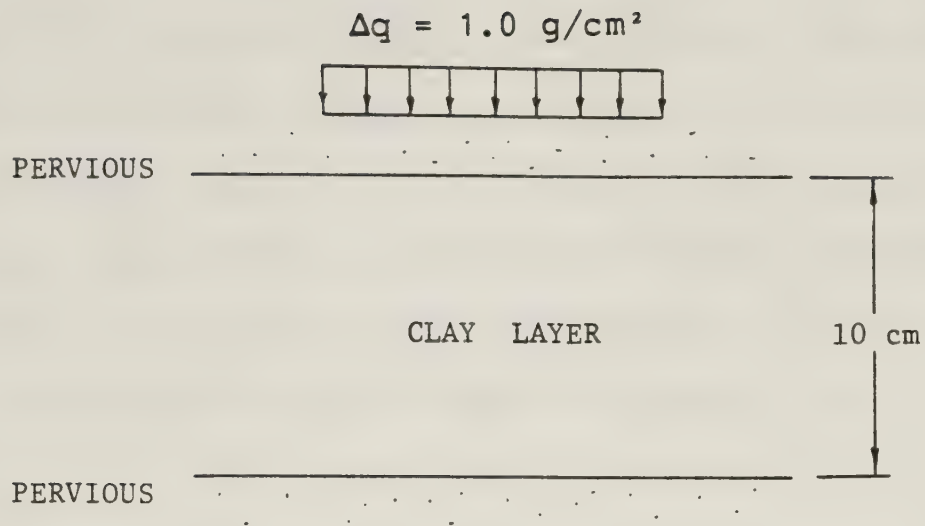
A simple single layer with drainage at both top and bottom is shown in Fig. 3.1. The applied loading and the material properties used in the analysis are given in the figure.

#### Finite Element Model:

For the finite element analysis, ten equal length one dimensional linear elements have been used. The consistent compressibility matrix option and the Euler backward method for time integration have been chosen for computation.







MATERIAL PROPERTIES:      ONE DIMENSIONAL FINITE ELEMENT  
MODEL OF THE CLAY LAYER

$$k = 1.0 \text{ cm/sec}$$

$$m_v = 1.0 \text{ cm}^2/\text{g}$$

$$C_v = 1.0 \text{ cm}^2/\text{sec}$$

$$\gamma_w = 1.0 \text{ g/cm}^3$$

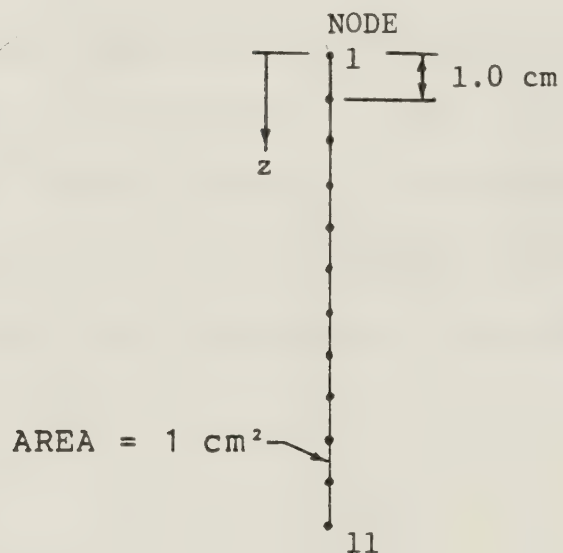


Figure 3.1 Analysis of one dimensional consolidation with double drainage



### Solution Results:

Fig. 3.2 shows the excess pore pressure isochrones predicted by ADINAT at time factors 0.1, 0.4 and 0.9. Only one half of the layer is shown because of the symmetry condition. Closed-form solution of Terzaghi's analysis is also given in the figure and good agreement is shown. The ADINAT input datafile is enclosed in Appendix A.1.

### Remarks:

The computer program used to compute the closed-form solutions is given in Appendix B.1.

A listing of ADINAT input datafile for the case of the impervious base is included in Appendix A.1. The physical problem, loading and finite element model is the same as above, except the thickness now is 5 cm. The pore pressure results, as would be expected, are the same as the upper half layer of the previous case.

### User Hints:

- a. For the case of double drainage, the initial conditions are applied to all but the top and bottom nodes, at which the deleted (zero) pressure condition



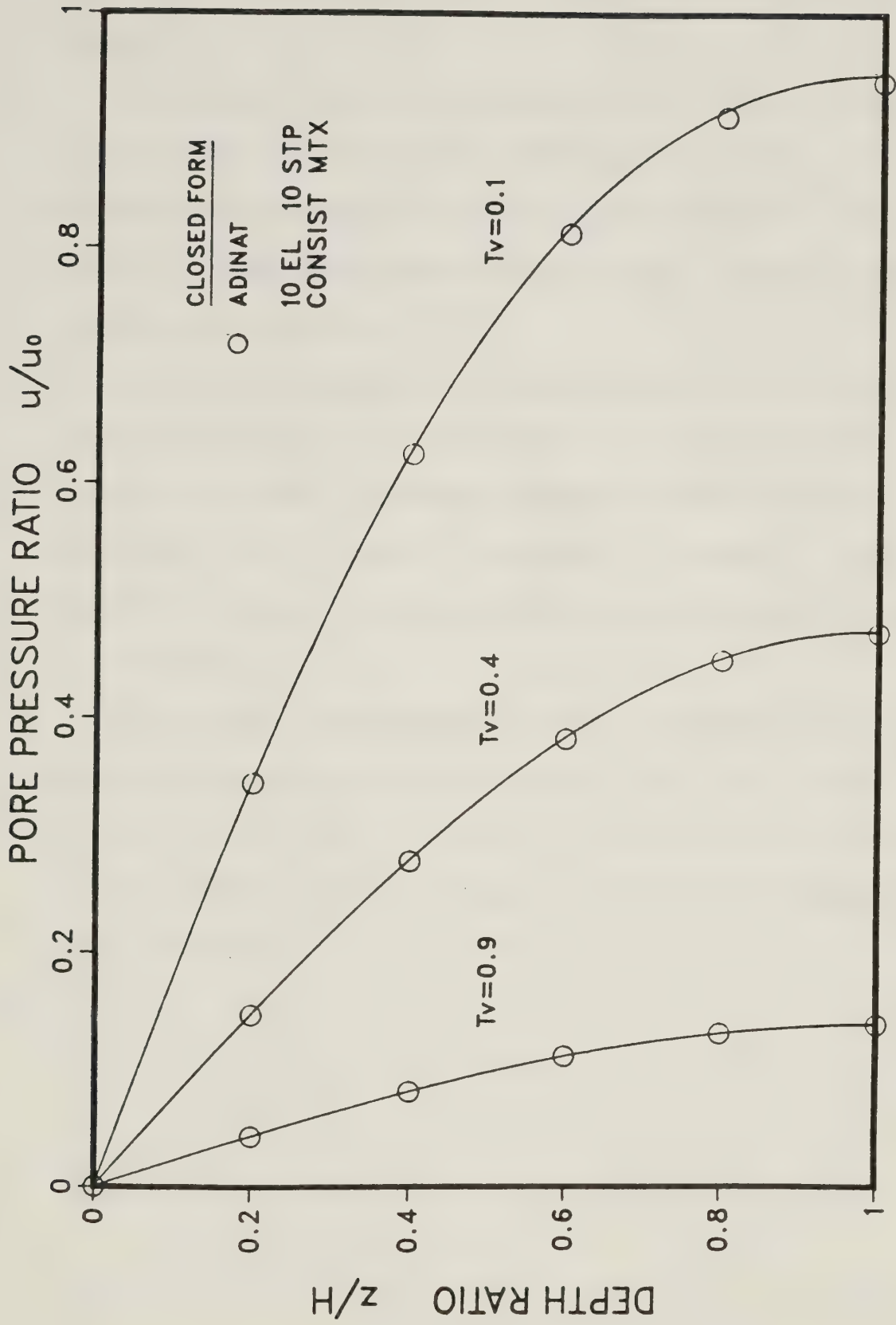


Figure 3.2 Excess pore pressure isochrones for half clay layer



code are assigned in the input. Only the pressure condition of the top node is ignored in singly drained layer.

- b. The output values of excess pore pressure can be obtained for some selected time factors  $T_v$  by deliberately manipulating the input values of the total number of time steps and time interval, which product is the solution time  $t$ .

There are generally three approaches to obtain the required results. One of these is simply to solve for one  $T_v$  at a time. Several computer runs are then necessary for a range of required  $T_v$  values. This method will give a uniform degree of accuracy for all the results.

The other two methods are to solve for all the required time factors in a single run, with either only one or various time intervals. The degree of accuracy for the solutions, obviously, will not be consistent for these methods.





## 3.2 Homogeneous Layered Problems

### 3.2.1 General

In the previous section, the consolidation of a homogeneous soil stratum is considered. Stratified soil systems, however, are a common occurrence. Post glacial lake clays usually show laminated or layered structures, due to various sediments sizes during seasonal deposition. Problems of unlike soil layers are often encountered in earth dams, embankment buildings and landfills.

In this section, the consolidation of layered soils is discussed and examples are given. Since the material properties are assumed constant in each layer, the problems are not considered as *nonlinear*.

A closed-form solution for two layers is given by Gray (1945) and for multiple layers by Barden and Younan (1969), Schiffman and Stein (1970). Numerical solutions are suggested by Abbott (1960), Luscher (1965), Christie (1966), Davis and Lee (1969). Experimental studies are provided by Raymond (1966) and Barden and Younan (1969).

The governing equation, Eq 3.1, can be rewritten as



$$C_v \frac{\partial^2 u^1}{\partial z^2}(z, t) = \frac{\partial u^1}{\partial t}(z, t) \quad (3.13)$$

where the superscript 1 refers to the 1<sup>th</sup> layer. The initial and boundary conditions are the same as in the preceding section. In addition to excess pore pressures, the continuity of flow between adjacent clay layers requires that flow velocities be equal at the common layer interfaces. These conditions are formulated as

$$u^1(z_1, t) = u^{1+1}(z_1, t) \quad (3.14)$$

and

$$\left[ k \frac{\partial u}{\partial z}(z_1, t) \right]_1 = \left[ k \frac{\partial u}{\partial z}(z_1, t) \right]_{1+1} \quad (3.15)$$

where  $z_1$  is the distance from the surface to the interface separating the 1<sup>th</sup> and (1+1)<sup>st</sup> layer. The void ratio - effective stress relationship differs for each layer, as shown in Fig. 3.3.



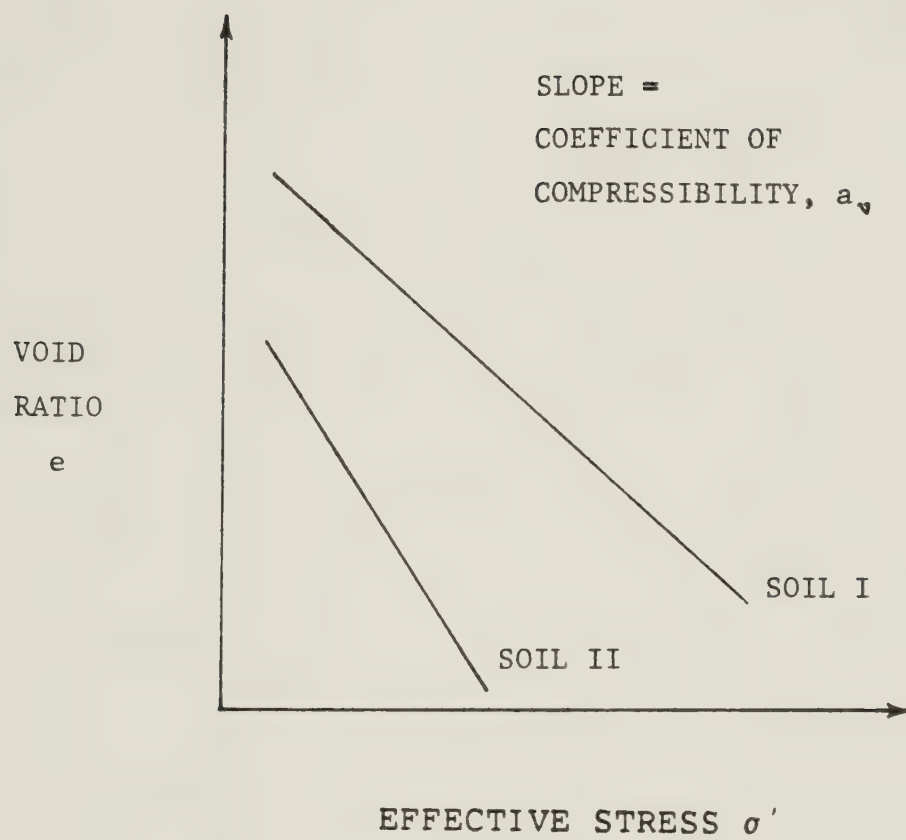


Figure 3.3 Void ratio - effective stress relationship in layered soils



Examples, consisting of two, three and four contiguous soil layers with singly or doubly drained boundaries, are analysed below. Conditions of uniform initial excess pore pressure and time-independent loadings are considered here. Closed-form, experimental or numerical results will also be included to compare with the ADINAT solutions.





### 3.2.2 Two Layers

#### Objective:

To verify ADINAT application to Gray's (1945) two layer consolidation problem.

#### Brief:

The first closed-form solution for a layered problem was published by Gray for two adjacent compressible strata. The impeded drainage condition was also considered in his paper.

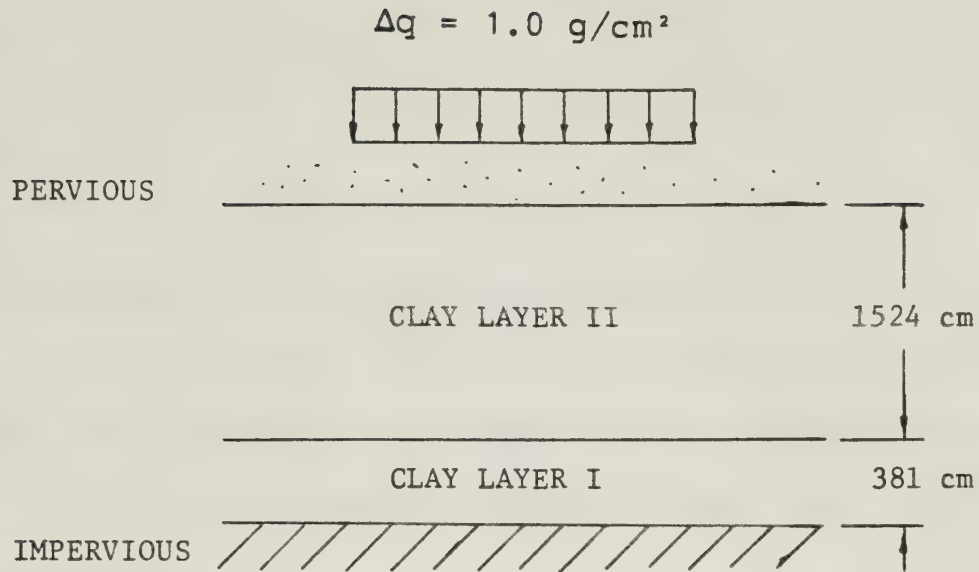
#### Physical Problem:

Two adjoining layers of unlike compressible soils under a constant applied load are shown in Fig. 3.4. Boundary conditions and material properties are given in the figure.

#### Finite Element Model:

In the finite element model 20 equal length elements are applied for the clay stratum. Similar to the previous example, a consistent compressibility matrix option and





#### MATERIAL PROPERTIES:

#### FINITE ELEMENT MODEL FOR THE TWO CLAY LAYERS

##### LAYER I

$$k = 2.5 \times 10^{-7} \text{ cm/sec}$$

$$m_v = 8.0 \times 10^{-5} \text{ cm}^2/\text{g}$$

$$C_v = 3.125 \times 10^{-3} \text{ cm}^2/\text{sec}$$

##### LAYER II

$$k = 2.5 \times 10^{-8} \text{ cm/sec}$$

$$m_v = 2.0 \times 10^{-4} \text{ cm}^2/\text{g}$$

$$C_v = 1.25 \times 10^{-4} \text{ cm}^2/\text{sec}$$

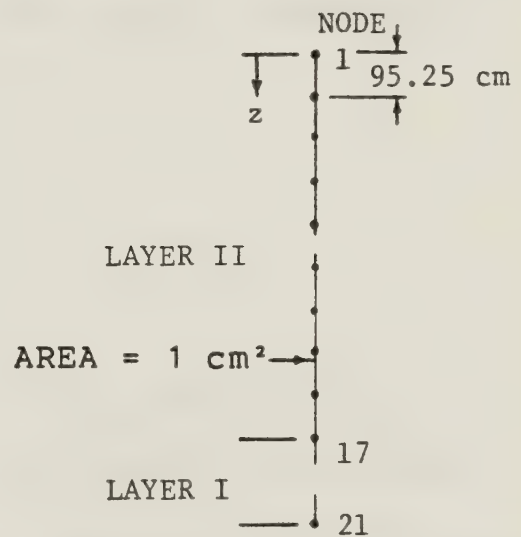


Figure 3.4 Consolidation analysis of two-layered problem



Euler backward method are chosen in the analysis.

### Results:

The ADINAT results of excess pore pressure are plotted in Fig. 3.5 with Gray's solutions and close agreement is observed. The parameter  $T_b$  shown in the figure is the time factor for the bottom layer. ADINAT input data is given in Appendix A.2.

### Others:

A Fortran program is used to compute the pore pressure according to Gray's formulation. The source code can be found in Appendix B.2.

### User Hints:

Two approaches can be used for the data input in layered problems. The material properties for each layer can be input in ADINAT, either as different linear element groups in control card I, or as different sets of material properties in a single element group in control card X. The same results will be obtained for these two methods. Both input methods are used in the sample analyses and they can be seen from the datafiles in Appendix A.



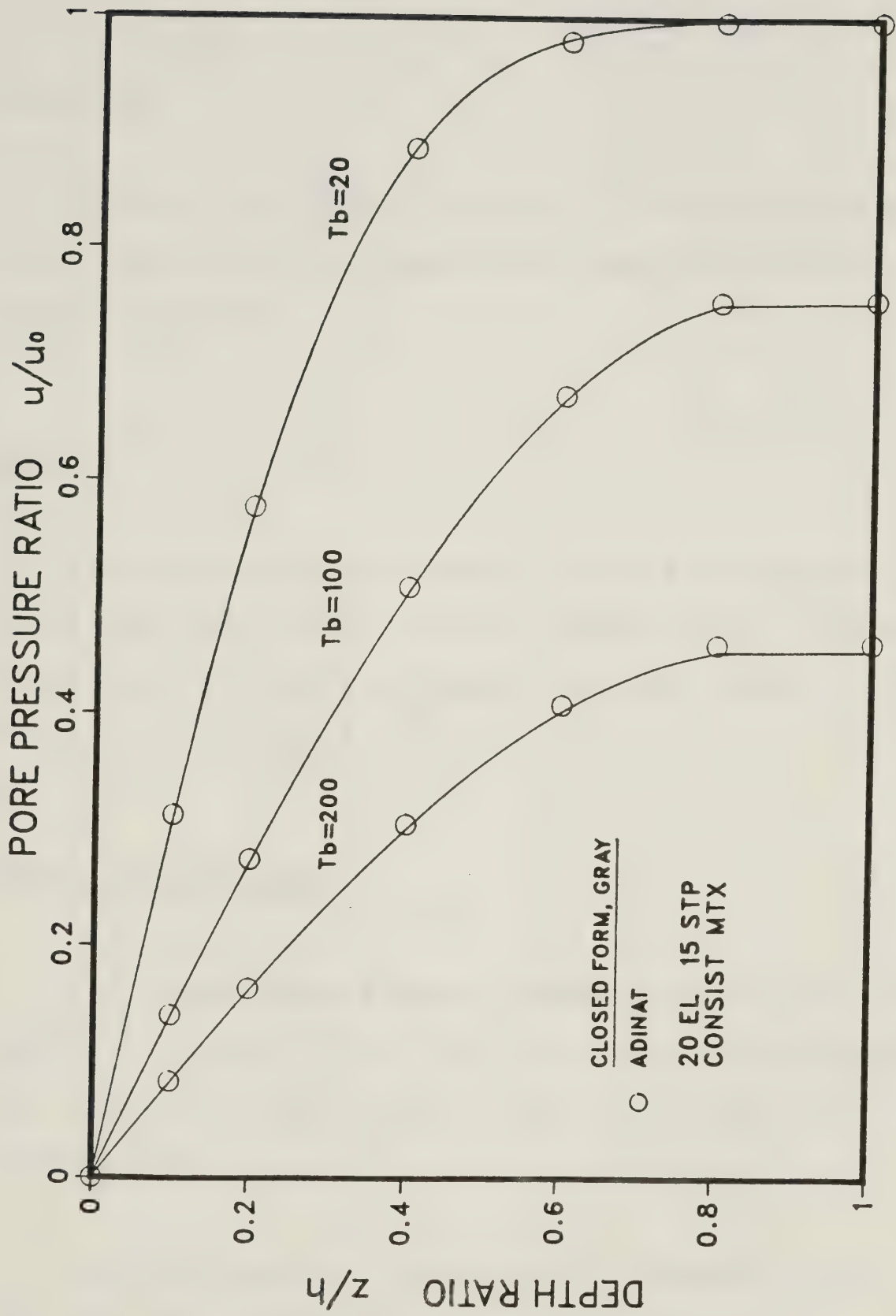


Figure 3.5 Pore pressure distribution for two-layered clay





### 3.2.3 Multiple-Layer System

#### Objective:

To verify the ADINAT analysis of the Barden and Younan (1969) solution of one dimensional consolidation for a three-layered soil.

#### Brief:

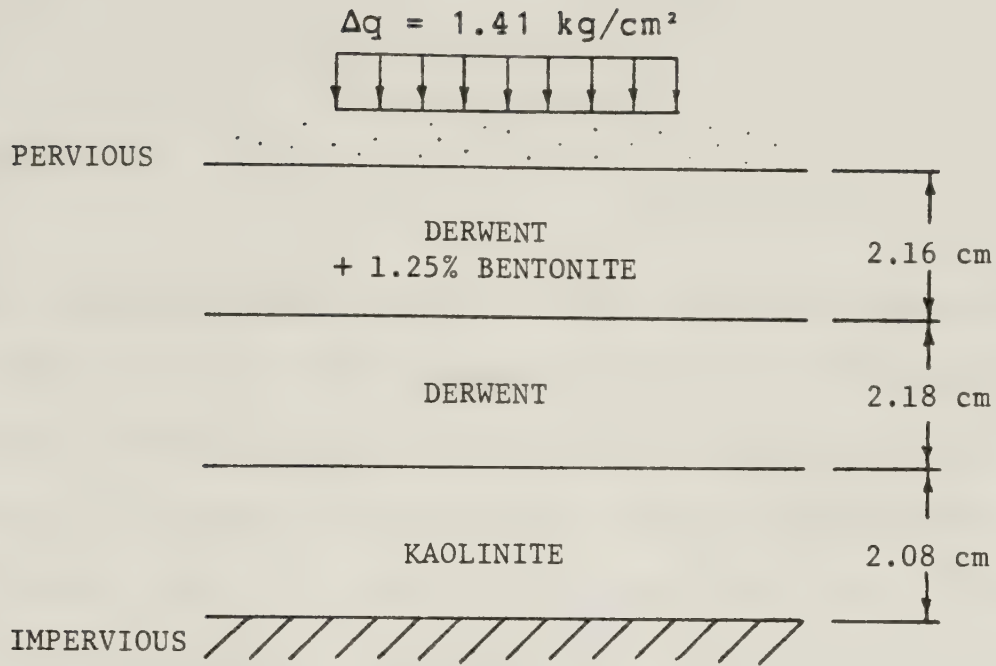
Closed-form and experimental results are given by these authors for artificially prepared layered soils. The same problem here is also analysed by Desai and Johnson (1972) using a finite element scheme.

#### General Descriptions:

The singly drained three-layered soils and the applied loading are shown in Fig. 3.6. The material properties of each layer are given in the figure. Input data is in Appendix A.3.

The finite element model is also included in the Fig. 3.6. In order to be comparable with Desai and Johnson's (1972) results, eight elements are used for each





#### MATERIAL PROPERTIES:

LAYER	$C_v$ ( $10^{-6} \text{ cm}^2/\text{yr}$ )	$k$ ( $\text{cm}/\text{yr}$ )
DERWENT + 1.25% BENTONITE	.716	.375
DERWENT	1.16	.546
KAOLINITE	7.01	3.14

#### FINITE ELEMENT MODEL FOR THE CLAY LAYER

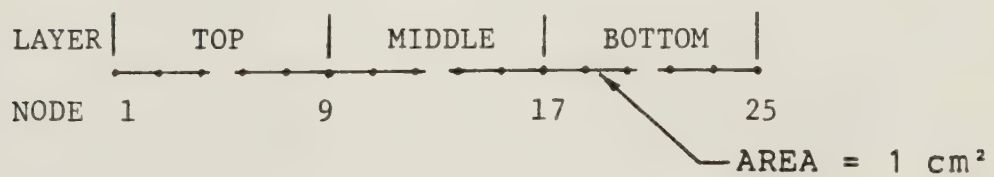


Figure 3.6 One dimensional consolidation analysis of a three-layered clay system



layer.

Results:

Pore pressure dissipation at the interface between Derwent and Kaolinite is plotted in Fig. 3.7. The ADINAT, closed-form, finite element (Desai and Johnson, 1972) and experimental results are compared. Good correspondence can be seen among the first three solutions over the entire range of consolidation, and with the experimental results above a time factor of about 20.



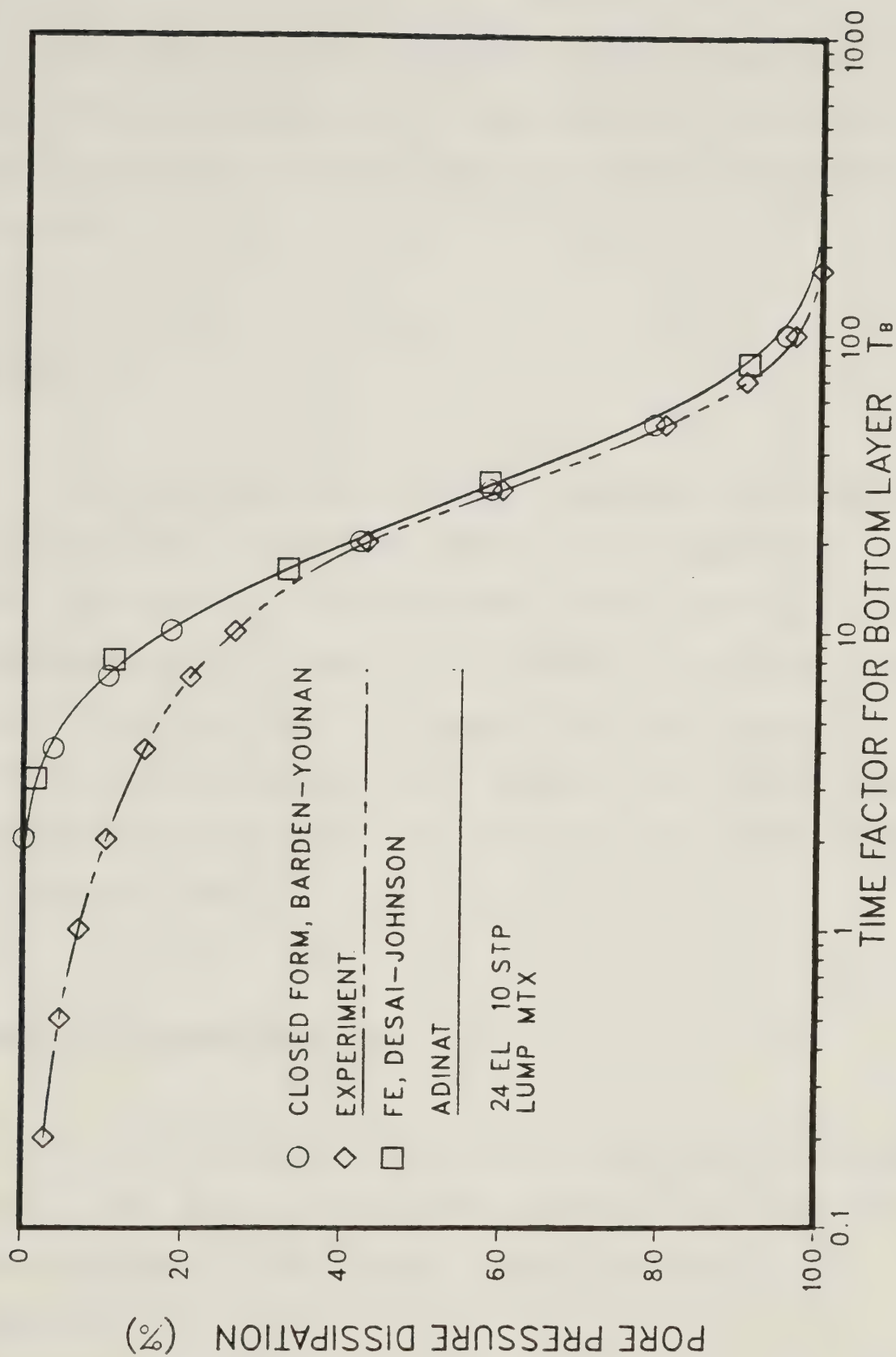


Figure 3.7 Pore pressure dissipation in three layers consolidation





### Objective:

To verify ADINAT predictions in one dimensional consolidation analysis of a four contiguous layer soil system with the Schiffman and Stein (1970) closed-form solution.

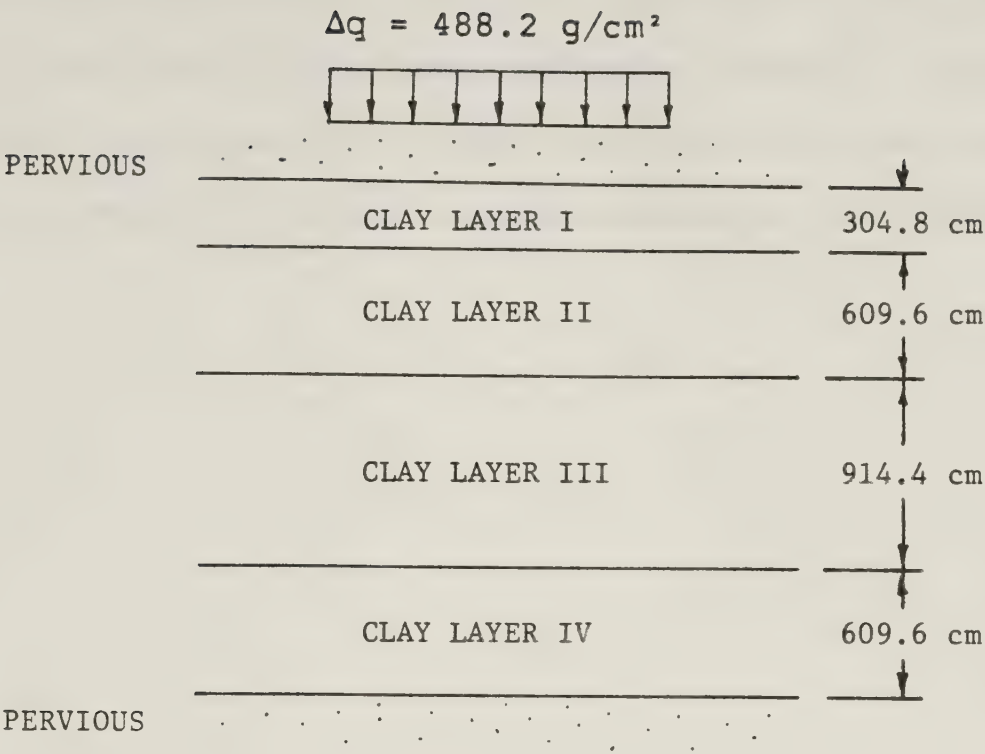
### Brief:

These authors have studied a multiple layer system with a general set of boundary conditions as impervious, impeded and free draining. Unrestricted loading history is also considered and a closed-form solution is provided using Fourier mathematics. An example of a four layer overconsolidated glacial till deposit with doubly drained boundaries and constant load is given. The same example is illustrated here.

### Descriptions and Results:

Fig. 3.8 shows the soil profile and the finite element model. The material properties given by Schiffman and Stein are tabulated in the figure. Input datafile is in Appendix A.4.





MATERIAL PROPERTIES:

LAYER	$C_v$ ( $10^{-4} \text{ cm}^2/\text{sec}$ )	$m_v$ ( $10^{-4} \text{ cm}^2/\text{g}$ )	$k$ ( $10^{-9} \text{ cm/sec}$ )
I	4.42	6.29	2.78
II	20.6	3.99	8.26
III	5.89	1.99	1.17
IV	7.38	3.99	2.95

FINITE ELEMENT MODEL FOR THE WHOLE CLAY LAYER

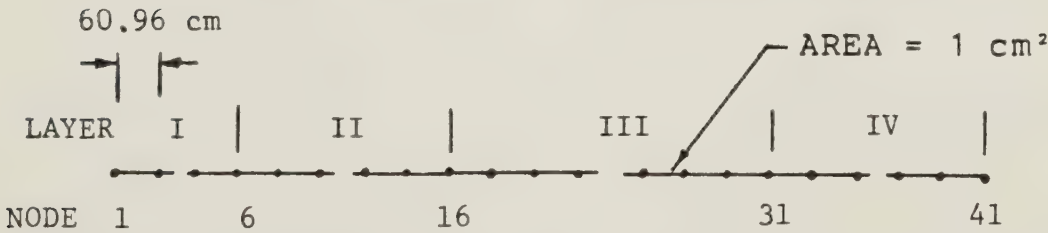


Figure 3.8 Consolidation analysis of a four-layered clay system



A set of excess pore pressure isochrones at 25%, 50% and 75% consolidation are computed using ADINAT. These results and the closed-form solutions from Schiffman and Stein are compared in Fig. 3.9, good agreement is observed.



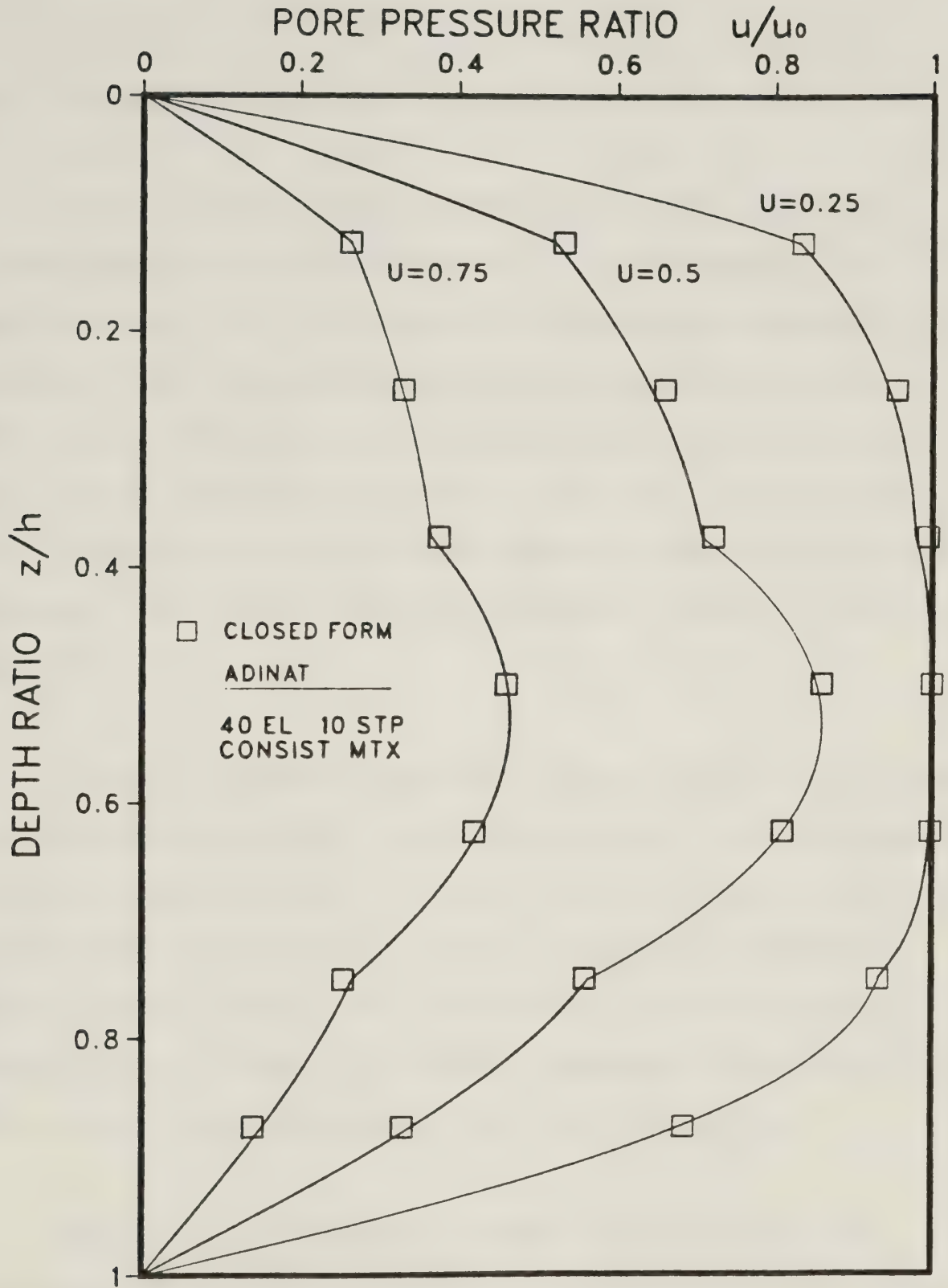


Figure 3.9 Pore pressure isochrones of a four-layered system





### 3.3 Continuous Spatial Variation of Material Properties

#### 3.3.1 General

Another type of spatial variation, where the soil properties vary with depth in a continuous manner, is examined here. In this case, the strains are assumed to be so small that the permeability and compressibility are effectively independent of time but are continuous functions only of the space coordinate  $z$ . This type of problem is only a special case of the more general *material nonlinear* problem. The latter will be developed further in the next chapter.

Values of permeability and compressibility at any depth vary greatly due to the diverse soil type, initial stress state, stress increment and past stress history. The range of possible variations in nature is so wide that only special cases of distributions which can be expressed in special form are used in general analysis. Therefore, only these cases are discussed in the given examples.

These known functions, however, can be useful to simulate many real situations. A polynomial variation can be used to approximate the conditions in normally consolidated or overconsolidated clay layers with sufficient generality. For example, in the case of overconsolidated



clay, where great amount of overburden has been eroded away in the geologic past, the upper layer usually is more permeable and compressible than at depth. This is illustrated in Fig. 3.10a.

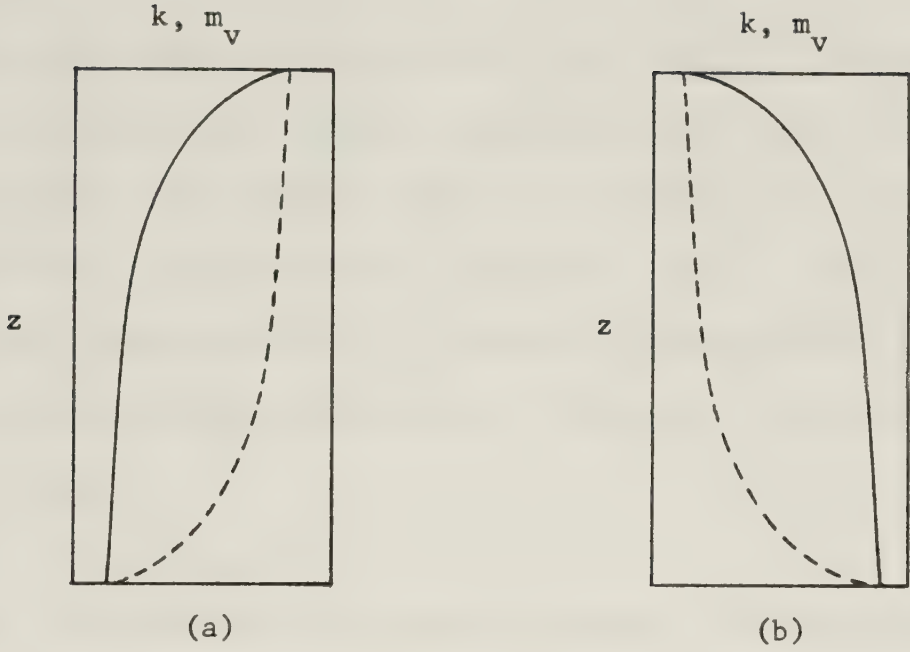
Fig. 3.10b shows the variation in a normally consolidated clay layer. The low permeability and compressibility at top is due to surface drying and the development of a hardened crust. As shown in Fig. 3.10c and Fig. 3.10d, the sinusoidal distribution is close to the condition of partly consolidated clay layers, in which the center part is more or less permeable or compressible than the edges. Examples with these distributions will be analysed below.

The governing equation Eq. 3.1 of the conventional consolidation theory is extended to include spatial variations in the coefficient of permeability and the coefficient of volume compressibility as follows, (Schiffman and Gibson, 1964)

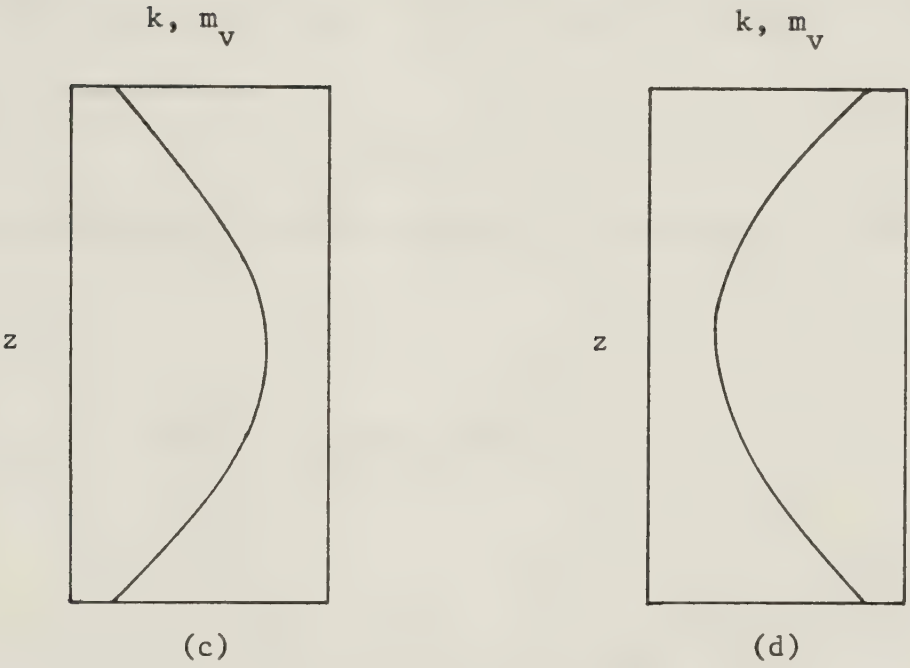
$$\frac{1}{\gamma} \frac{\partial}{\partial z} \left\{ k(z) \frac{\partial u}{\partial z} (z, t) \right\} = m_v(z) \left\{ \frac{\partial u}{\partial t} (z, t) \right\} \quad (3.16)$$

The boundary and loading conditions are the same as in





POLYNOMIAL VARIATION



SINUSOIDAL VARIATION

Figure 3.10 Polynomial and sinusoidal functions for material properties distribution



### Section 3.1.

Note that even if  $k(z)$  and  $m(z)$  vary in the same manner and thus make  $C_v(z)$  a constant (see Eq. 3.7), Eq. 3.16 does not reduce to Eq. 3.1. This only takes place when  $k$  and  $m$  are separately constant. Eq. 3.16, however, is only an approximation to a general analysis of a thick layer with nonlinear properties. This will be discussed in the next chapter.

Note further that values of average consolidation for nonhomogeneous problems are determined from the settlement ratios. The same applies to nonlinear problems in the following chapters.

The polynomial relationship is expressed in the form of

$$KM(z) = KM_0 (1 + \eta z')^\epsilon \quad (3.17)$$

while the sinusoidal variations are given by

$$KM(z) = KM_0 + A \cdot \sin(\pi z') \quad (3.18)$$





in which  $KM$  indicates either the permeability or compressibility,  $z'$  is the depth ratio  $z/h$ . The subscript  $\phi$  denotes values at the top of the soil layer. The parameters  $\epsilon$  and  $\eta$  describe the mode of variation of  $k$  and  $m$ .  $A$  is the amplitude of the sine function.

Published analytical solutions can only be found for a limited number of particular cases. Edelmann (1953) considered the linear variation in the modulus of elasticity and Martins (1965) has analysed the case of a linear variation in permeability. Closed-form solutions are provided for both cases. Schiffman and Gibson (1964) have considered a series of cases for the spatial variations in permeability and compressibility. Known distributions such as polynomial and sinusoidal functions are used in their analysis.



### 3.3.2 Variable Permeability with Constant Compressibility

#### Objective:

To verify ADINAT application to one dimensional consolidation with linear variation in permeability and constant compressibility.

#### Brief:

Martins (1965) studied the special case of linear permeability and constant compressibility. He assumed  $\epsilon=1$  in Eq. 3.17 for compressibility and use a polar coordinate transformation of Eq. 3.16, which becomes

$$\chi \left\{ -\frac{1}{r} \frac{\partial u}{\partial r}(r,t) + \frac{\partial^2 u}{\partial r^2}(r,t) \right\} = \frac{\partial u}{\partial t}(r,t) \quad (3.19)$$

where

$$\chi = k_0 \eta^2 / (4 \gamma_w h^2 m_v) \quad (3.20)$$

and



$$r = (1 + \eta z')^{1/2} \quad (3.21)$$

$k_0$  is the value of the coefficient of permeability at the layer surface. This equation can be considered as the diffusion equation of an infinite cylinder with radial flow and constant diffusivity, where Carslaw and Jaeger's (1959) solution is then applied.

### General Description:

Fig. 3.11 shows the soil profile and the boundary conditions of the problem. The material properties and finite element model are given in the figure. Consistent compressibility matrix computation and Euler backward method are used. Datafile is included in Appendix A.5.

### Results:

The pore pressure isochrone predicted by ADINAT at 50% consolidation is given in Fig. 3.12, and compared with Martins' closed-form solution. Good agreement is shown. The maximum discrepancy is about 4% at the bottom part of the curve.



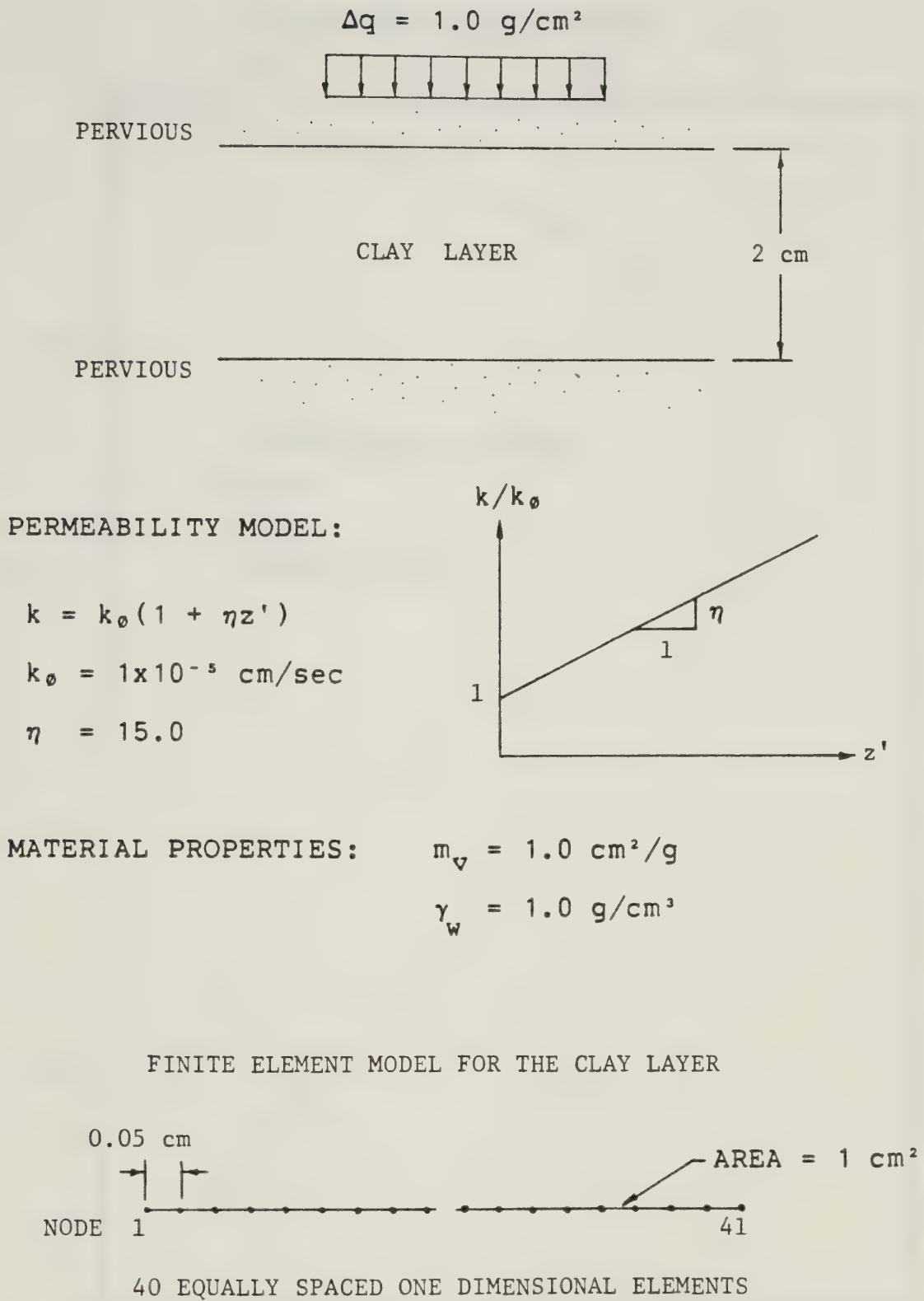


Figure 3.11 One dimensional consolidation analysis with linear variation in permeability





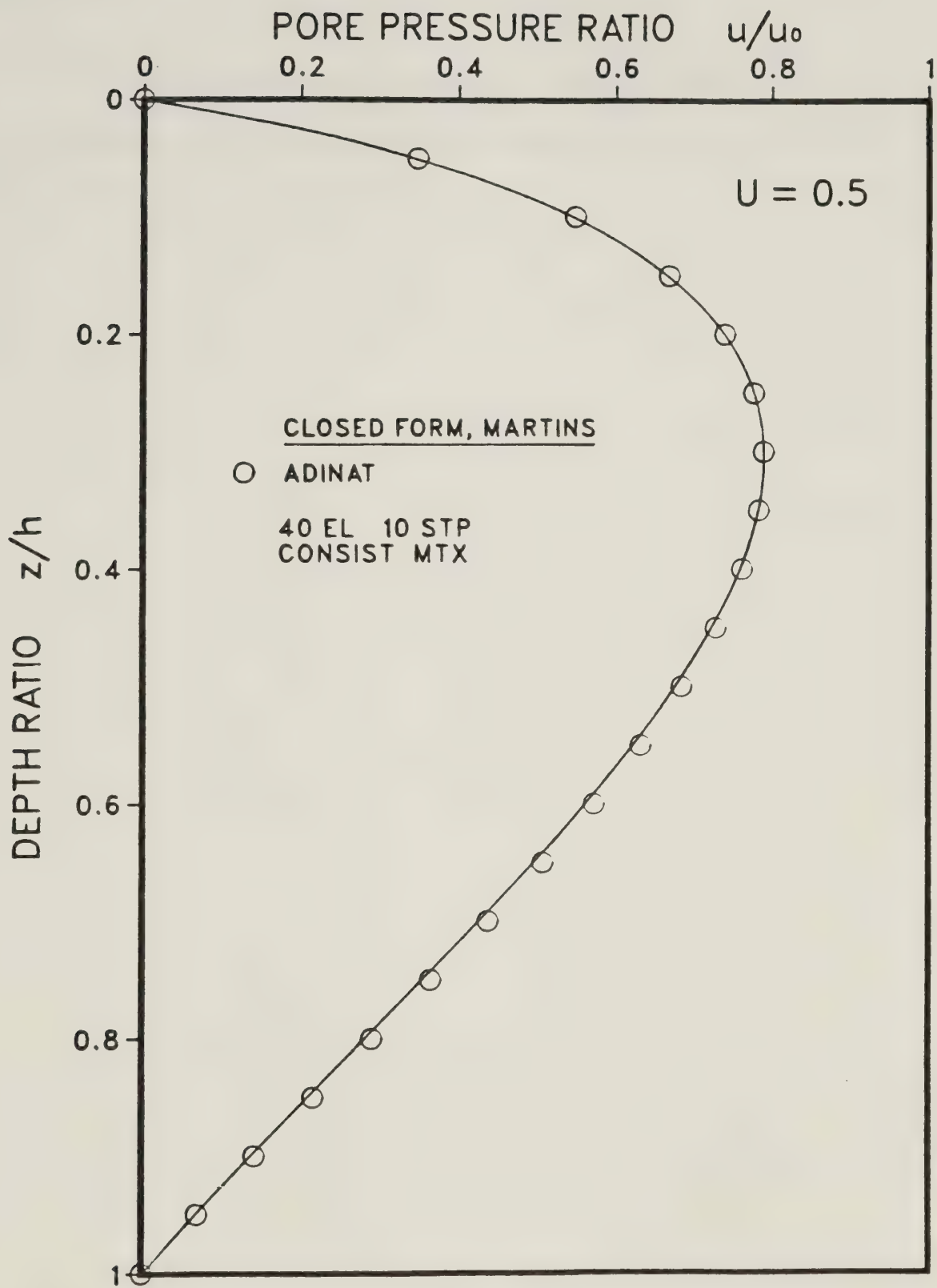


Figure 3.12 Pore pressure isochrones: linear variation in permeability



Remarks:

A Fortran program is used to compute Martins' formulation. A listing of the source code is included in Appendix B.3.



### Objective:

To show ADINAT analysis of one dimensional consolidation with polynomial distribution in permeability and constant compressibility.

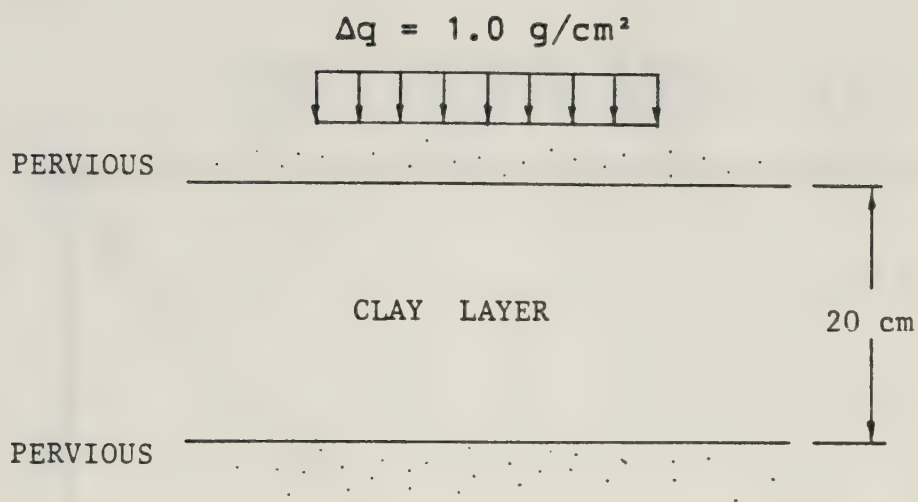
### General Description:

Eq. 3.16 and 3.17 are applied in the analysis. The parameters  $\epsilon$  and  $\eta$  estimated from Schiffman and Gibson (1964) are used. Fig. 3.13 shows the soil profile and the finite element model. Forty equally spaced one dimensional elements will be used for the clay layer. Euler backward method and lumped compressibility matrix are utilized in the computation. The ADINAT input data is given in Appendix A.6

### Results:

The ADINAT output of excess pore pressure at 50% consolidation is compared with the solution of Schiffman and Gibson and presented in Fig. 3.14. Good correspondence is observed.





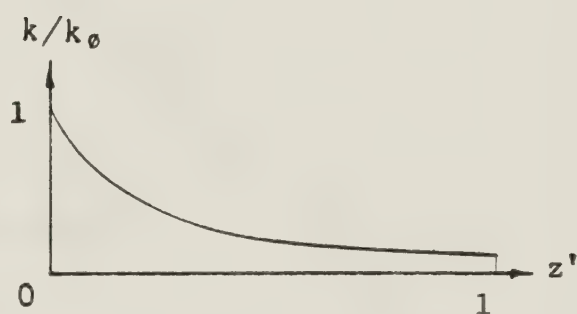
PERMEABILITY MODEL:

$$k = k_0 (1 + \eta z')^\epsilon$$

$$k_0 = 1.0 \text{ cm/sec}$$

$$\eta = 8.0$$

$$\epsilon \approx -1.05$$



MATERIAL PROPERTIES:

$$m_v = 1.0 \text{ cm}^2/\text{g}$$

$$\gamma_w = 1.0 \text{ g/cm}^3$$

FINITE ELEMENT MODEL FOR THE CLAY LAYER

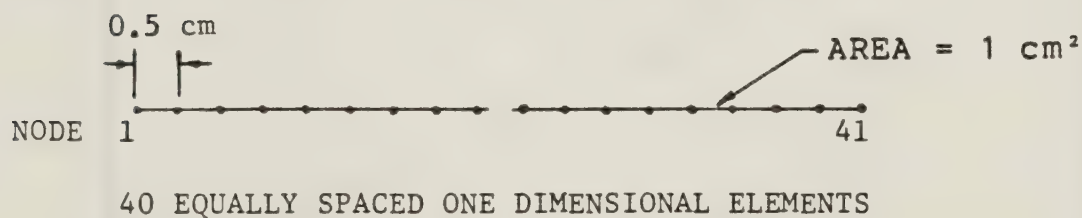


Figure 3.13 One dimensional consolidation analysis with polynomial variation in permeability





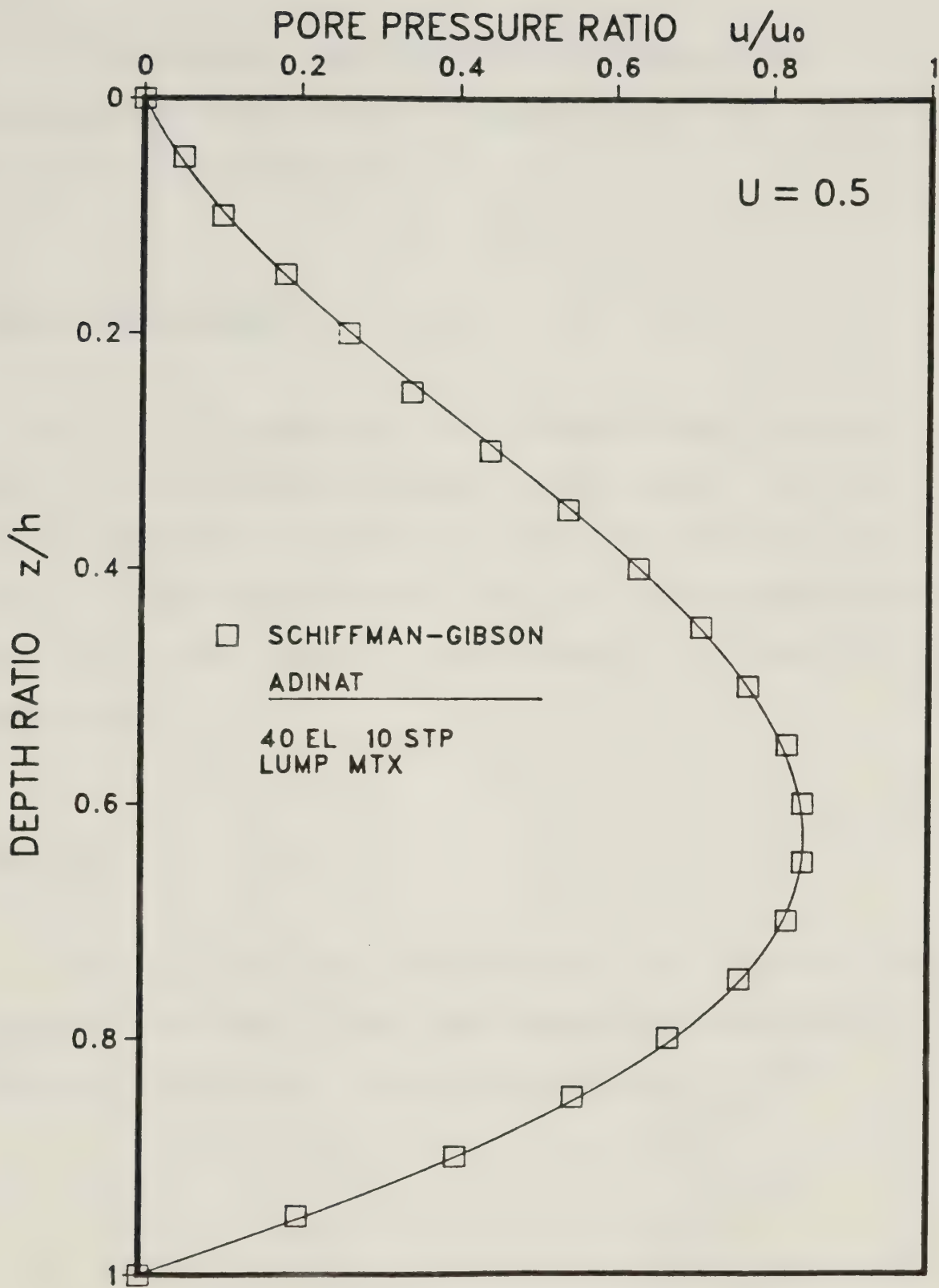


Figure 3.14 Pore pressure isochrones: polynomial variation in permeability



### Objective:

To show ADINAT application to one dimensional consolidation with constant compressibility and sinusoidal variation in permeability.

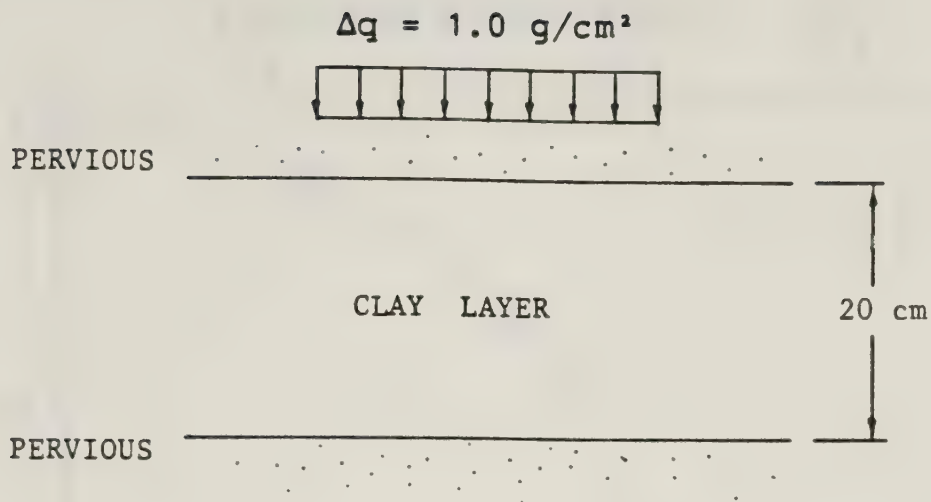
### General Description:

Eq. 3.18 is employed with the amplitude  $A$  obtained from Fig. 13 in Schiffman and Gibson (1964) paper. The soil profile and material properties are given in Fig. 3.15. The finite element model and the time integration method are the same as the last example. ADINAT datafile is shown in Appendix A.7.

### Results:

Fig. 3.16 shows the comparison of the resulting pore pressure isochrone at 50% consolidation and the solution by Schiffman and Gibson. Good correspondence is noted.





PERMEABILITY MODEL:

$$k = k_0 + A \cdot \sin(\pi z')$$

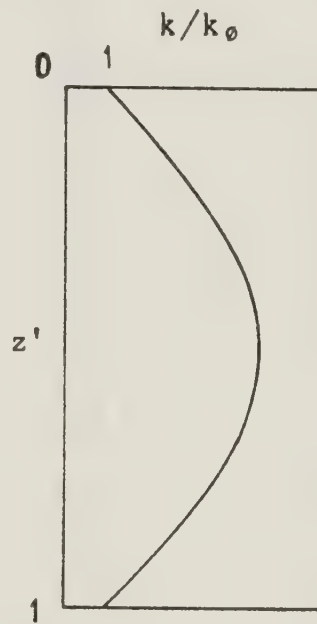
$$k_0 = 1.0 \text{ cm/sec}$$

$$A = 9.0 \text{ cm/sec}$$

MATERIAL PROPERTIES:

$$m_v = 1.0 \text{ cm}^2/\text{g}$$

$$\gamma_w = 1.0 \text{ g/cm}^3$$



FINITE ELEMENT MODEL FOR THE CLAY LAYER



40 EQUALLY SPACED ONE DIMENSIONAL ELEMENTS

Figure 3.15 One dimensional consolidation analysis with sinusoidal variation in permeability



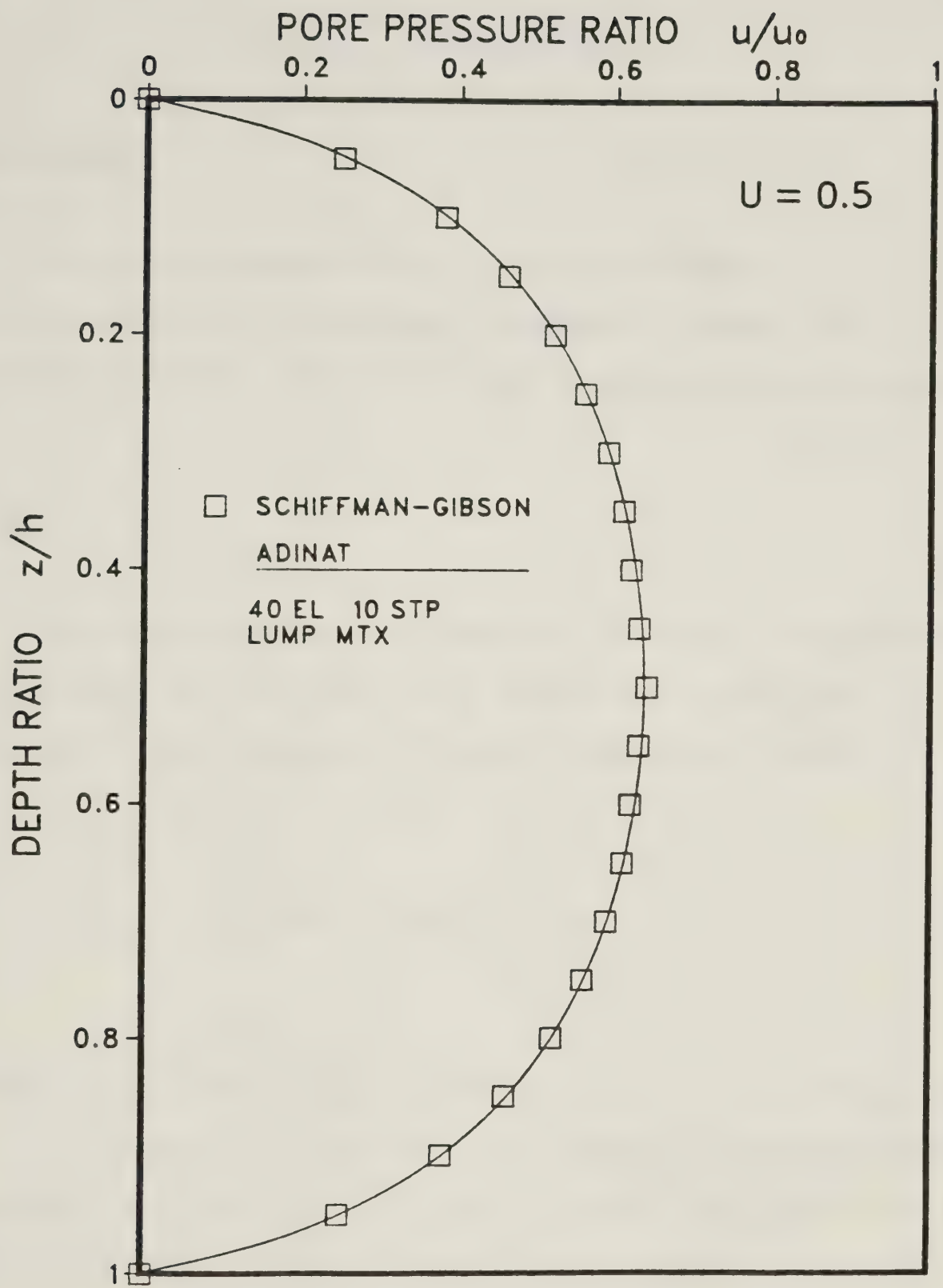


Figure 3.16 Pore pressure isochrones: sinusoidal variation in permeability





### 3.3.3 Variable Compressibility with Constant Permeability

#### Objective:

To verify ADINAT analysis of the variable compressibility and constant permeability model with Edelman's (1953) solution in one dimensional consolidation.

#### Brief:

Edelman assumed the linear variation of the modulus of elasticity,  $E$ , in a soil layer whence the coefficient of compressibility becomes a rational function of depth,

$$m_v(z) = 1/(c'\gamma_z + \sigma'_0) \quad (3.22)$$

where  $\sigma'_0$  is a small initial effective stress at surface and  $c'$  is a constant characterize the manner of variation of  $E$ . Putting this into Eq. 3.16, which can be simplified into a different form



$$\frac{\partial^2 u}{\partial z^2}(z, t) = \frac{1}{z + \eta} \cdot \frac{\partial u}{\partial T_1}(z, t) \quad (3.23)$$

where

$$\eta = \sigma'_0 / (c' + \gamma_t) \quad (3.24)$$

and

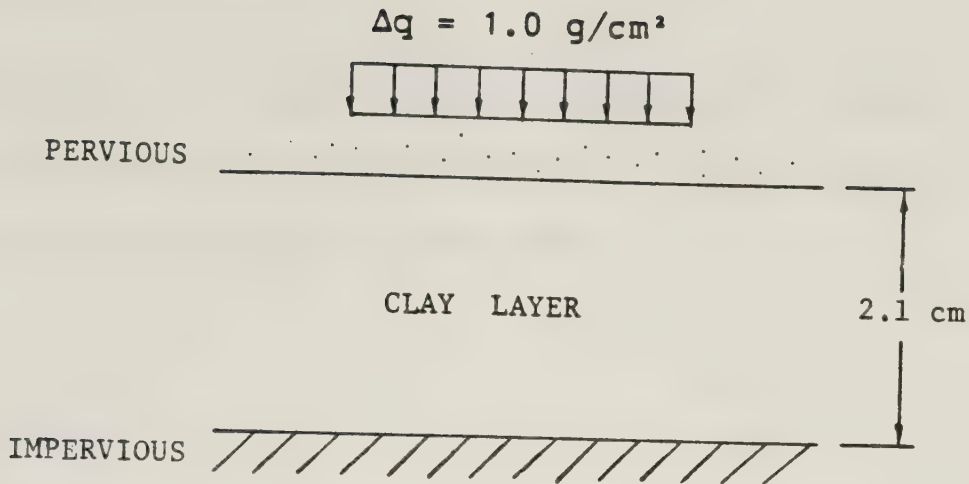
$$T_1 = (c' + k + \gamma_t + t) / \gamma_w \quad (3.25)$$

Edelmann employed Fourier mathematics to obtain a closed-form solution for the case of a single layer with only top drainage.

### Description of Problem:

The singly drained soil profile and loading are presented in Fig. 3.17. The compressibility model and material properties are given in the figure. The ADINAT datafile is given in Appendix A.8.



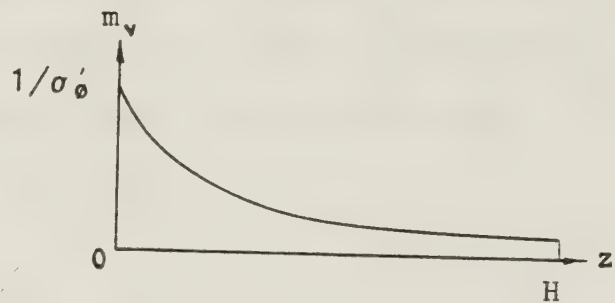


COMPRESSIBILITY MODEL:

$$m_v = \frac{1}{c' \gamma_t z + \sigma'_0}$$

$$c' = 15$$

$$\sigma'_0 = 6 \text{ g/cm}^2$$



MATERIAL PROPERTIES:

$$k = 1 \times 10^{-5} \text{ cm/sec}$$

$$\gamma_t = 1.0 \text{ g/cm}^3$$

$$\gamma_w = 1.0 \text{ g/cm}^3$$

FINITE ELEMENT MODEL FOR THE CLAY LAYER



40 EQUALLY SPACED ONE DIMENSIONAL ELEMENTS

Figure 3.17 One dimensional consolidation analysis with variable compressibility



### Finite Element Model:

The finite element model consists of 40 equally spaced elements. Consistent compressibility matrix option and Euler backward method is employed.

### Results:

Edelmann's closed-form solution at 50% consolidation is compared with the ADINAT pore pressure results in Fig. 3.18 and good agreement is shown. Maximum error of the pore pressure ratio occurs at the top part of the plot and is in the order of  $10^{-2}$ .

### Remarks:

A Fortran IV program is prepared for Edelmann's solution. The roots for the Bessel functions in the calculation can be found from Abramowitz and Stegun (1972). The required time for the analysis is computed by the program. A listing of the source code is included in Appendix B.4.





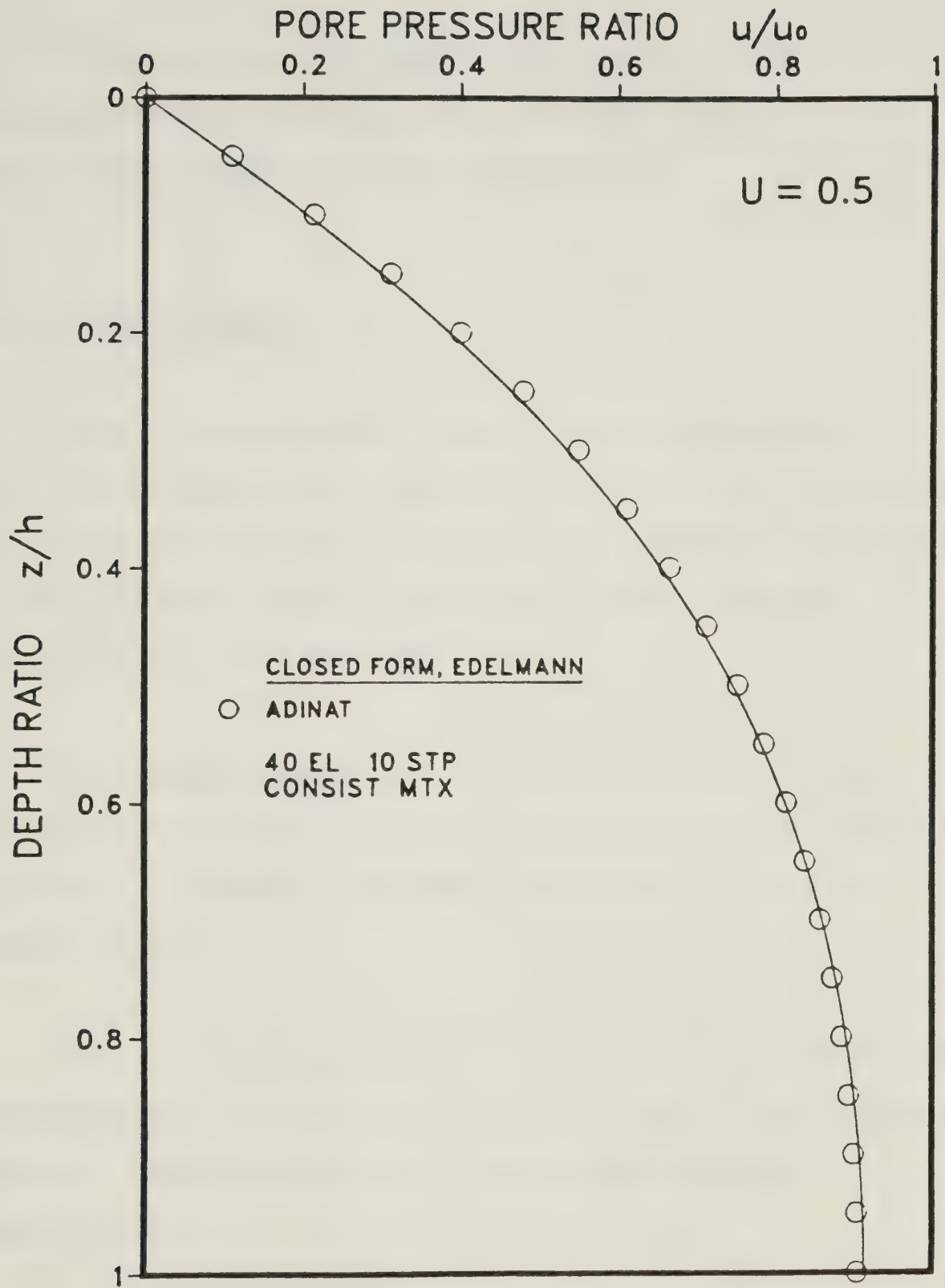


Figure 3.18 Pore pressure isochrones: variable compressibility



### Objective:

To demonstrate the ADINAT application in one dimensional consolidation with sinusoidal variation in compressibility and constant permeability.

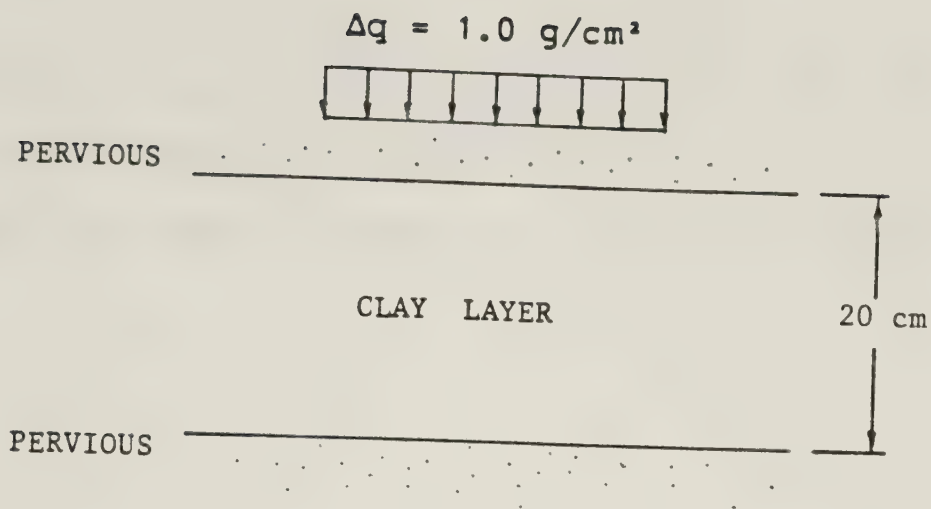
### General Description:

The soil profile and loading for the problem is described in Fig. 3.19. With the values  $\epsilon$  and  $\eta$  estimated from Schiffman and Gibson (1964), the compressibility model is obtained and given in the figure. Other material properties are also included.

The required time in the analysis is estimated by Simpson's rule in the numerical integration of the average settlement. A sample of ADINAT input data is given in Appendix A.9.

The finite element model, time integration method and compressibility matrix option are the same as the previous example. These can also be found in the datafile in Appendix A.9.





COMPRESSIBILITY MODEL:

$$m_v = m_{v0} - A \cdot \sin(\pi z')$$

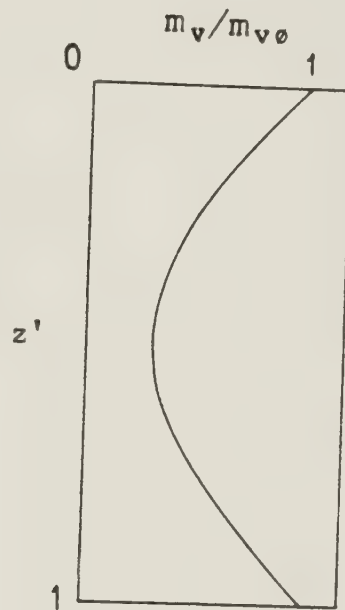
$$m_{v0} = 1.0 \text{ cm}^2/\text{g}$$

$$A = 0.9 \text{ cm}^2/\text{g}$$

MATERIAL PROPERTIES:

$$k = 1.0 \text{ cm}^2/\text{sec}$$

$$\gamma_w = 1.0 \text{ g/cm}^3$$



FINITE ELEMENT MODEL FOR THE CLAY LAYER



40 EQUALLY SPACED ONE DIMENSIONAL ELEMENTS

Figure 3.19 One dimensional consolidation analysis with sinusoidal variation in compressibility



Results:

The excess pore pressure isochrone at 50% consolidation predicted by ADINAT is compared with the solution by Schiffman and Gibson in Fig. 3.20, and close agreement is found.





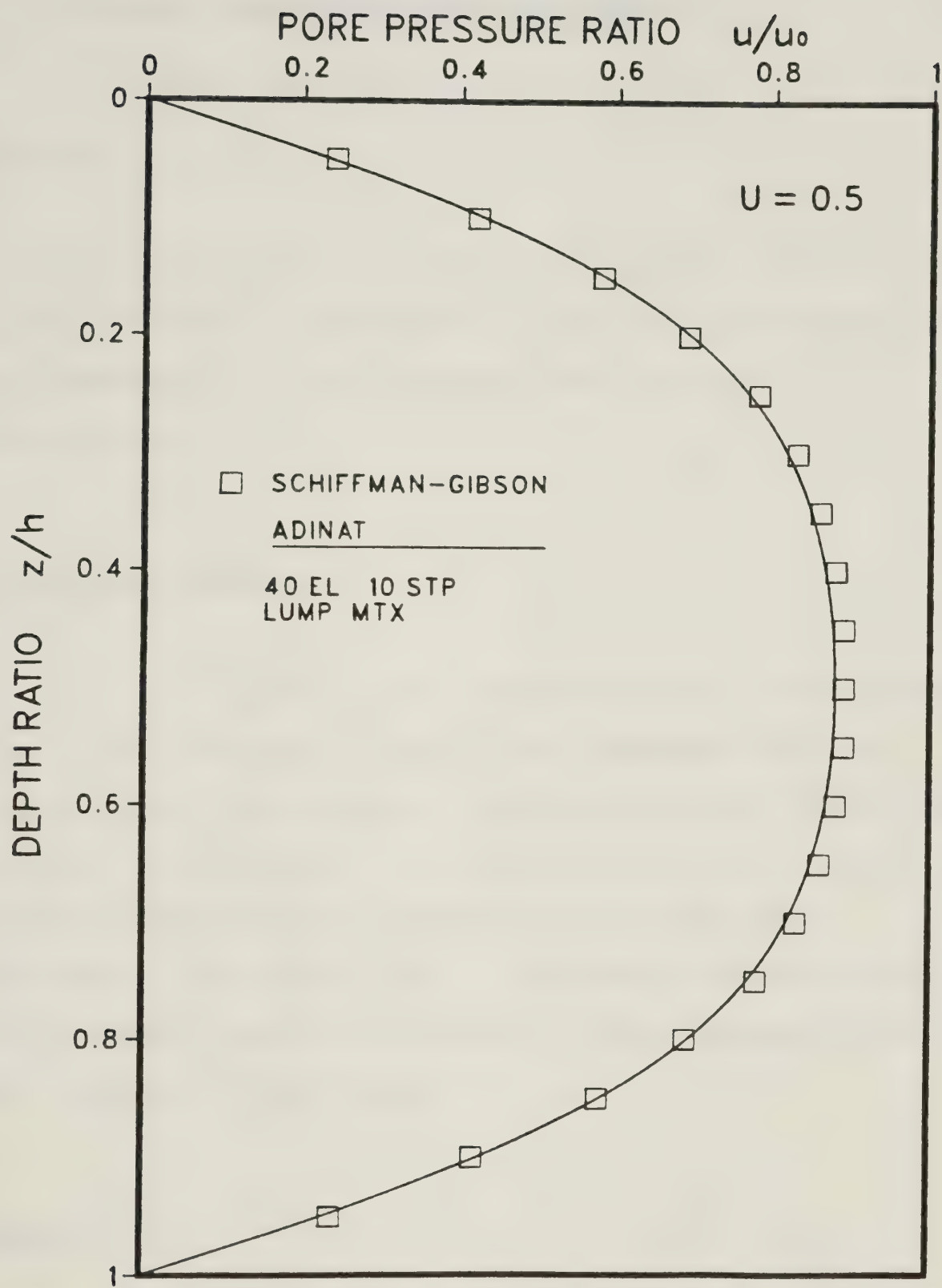


Figure 3.20 Pore pressure isochrones: sinusoidal variation in compressibility



### 3.3.4 Geometric Variation of Material Properties with Constant Coefficient of Consolidation

#### Objective:

To illustrate the application of ADINAT in the analysis of one dimensional consolidation with variable permeability and compressibility but constant coefficient of consolidation.

#### General Description:

Fig. 3.21 gives the soil profile and the finite element model. The estimated values of the parameters for the material model are obtained from Schiffman and Gibson (1964) and given in the figure. The estimated time is found from Simpson's rule for the integration of the average settlement. The ADINAT data is included in Appendix A.10. Finite element model is the same as previous example and is also included in the figure.

#### Results:

The pore pressure predicted by ADINAT is presented in Fig. 3.22. This same problem is also analysed by Schiffman and Gibson (1964), close approximation is shown.



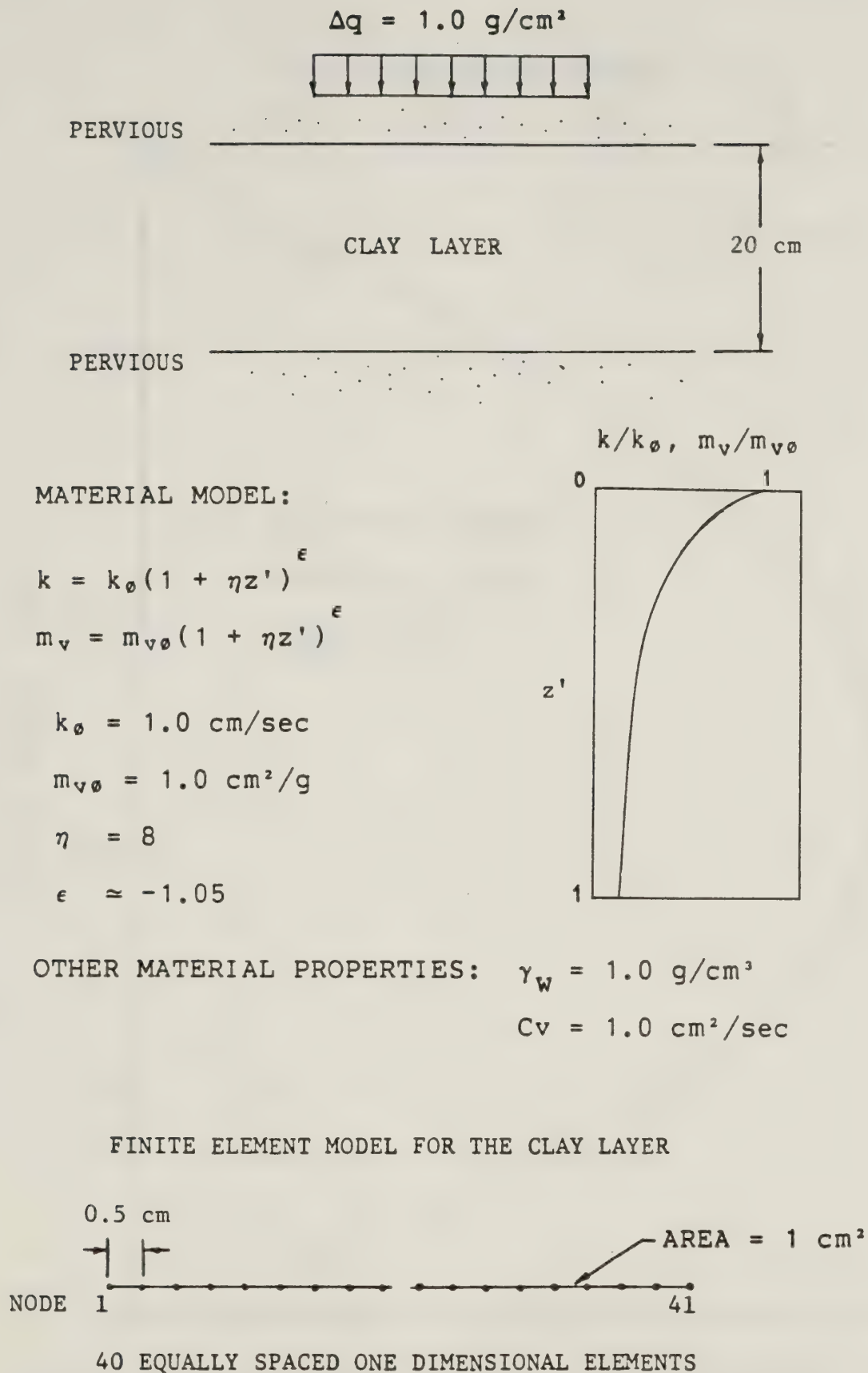


Figure 3.21 One dimensional consolidation analysis with polynomial distribution in material properties



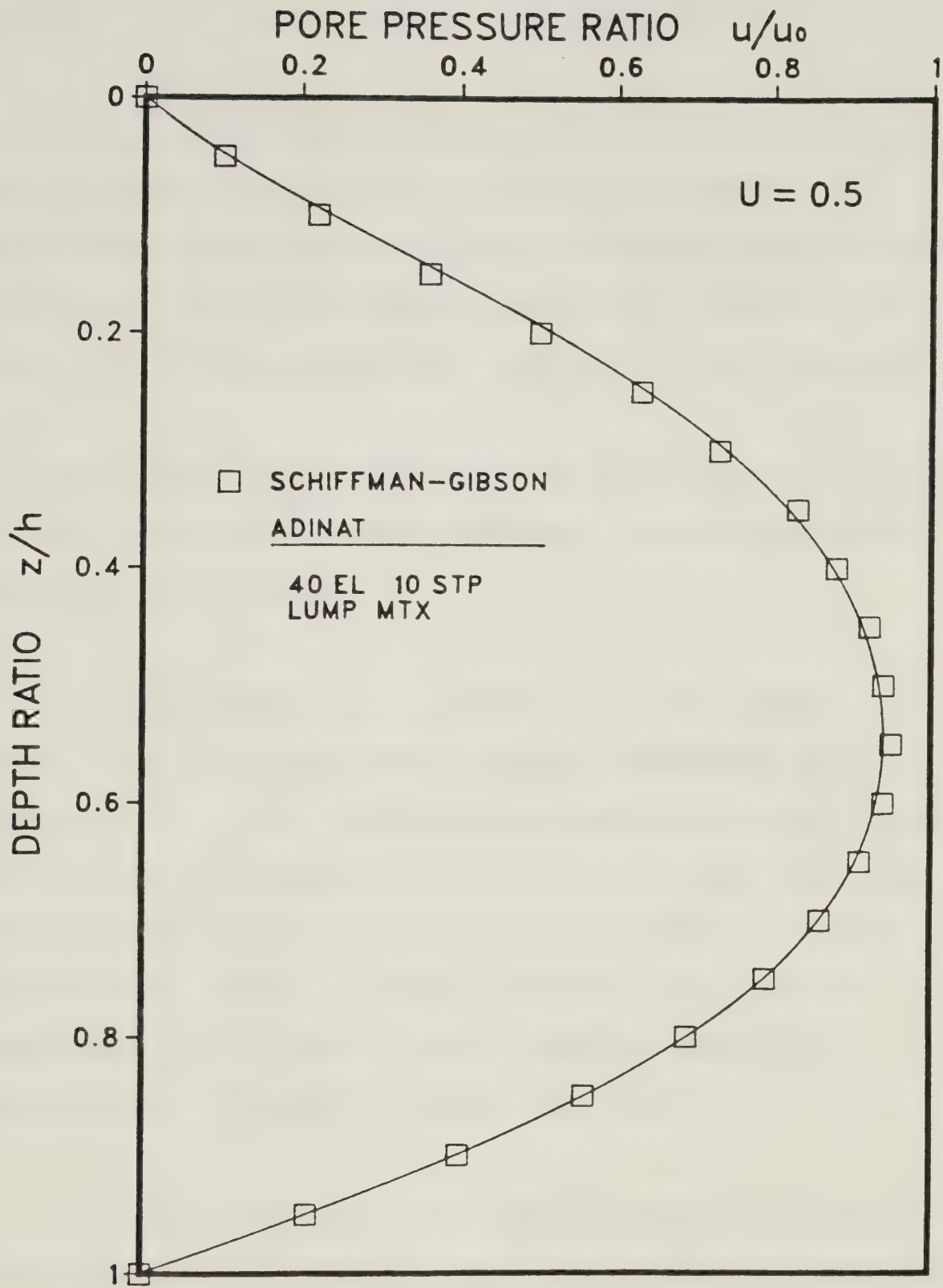


Figure 3.22 Pore pressure isochrones: variable permeability & compressibility with constant coefficient of consolidation





### 3.4 Discussion

In most of the above examples, closed-form solutions are given for a qualitative evaluation of ADINAT predictions. The high accuracy of the ADINAT results found in the examples in this chapter serves as a check on the reliability of the program when applied to such analyses.

The average cost of each run is about thirty to fifty cents for the size of above problems. This indicates the economy in using ADINAT.

Since the excess pore pressure  $u$  is the primary unknown, the finite element procedure guarantees the continuity of  $u$ , but not the pore pressure gradient,  $\partial u / \partial z$ . The continuity requirement in layered problems, Eq. 3.15, is therefore not strictly satisfied for the case of linear finite element model in ADINAT. However, the accurate predictions in the above layered examples assure the applicability of ADINAT in these analyses.

For layered problems with severe changes in material properties, however, higher-order finite element model may become necessary. A further discussion on this subject can be found in Desai and Johnson (1972).



For the analyses in this chapter, pore pressure results are obtained directly from the program output. Subroutines or supplemental programs may be then required for the computation of other quantities such as average pressure dissipation or settlement. The settlement ratio becomes a more important measure for comparisons in nonlinear problems.



#### 4. GENERAL STUDY OF ADINAT APPLICATION IN NONLINEAR CONSOLIDATION

Nonlinear consolidation theories have been developed for decades and numerous studies can be found. The scope of discussions in this and subsequent chapters are in accordance with the first seven assumptions given in Section 3.1.1.

Published nonlinear theories are reviewed first. The general governing equation is then derived, in order to facilitate the discussions that follows. It will be shown to embrace most nonlinear theories previously cited.

From this equation, different formulations with the common major simplifications will be generated. They will be presented in a form suitable for use with ADINAT. These formulations form the basis for the numerical treatment of various nonlinear theories using the program. Applications of these *ADINAT equations* to existing theories will be illustrated in subsequent chapters.



#### 4.1 Review of Nonlinear Consolidation Theories

For many years, attempts have been made by researchers to modify the assumptions implicit in Terzaghi's theory to agree more closely with the physical conditions and material properties in real soils. The changes of permeability and compressibility in a mass of real soil are two significant departures from Terzaghi's classical theory, which result in a nonlinear theory.

The first attempt to deal with these nonlinearities was by Schiffman and Gibson (1964), as illustrated in Section 3.3, where they considered the material properties as spatially dependent. Since the effective stress also varies with depth, the nonlinear behavior is then treated as a geometric nonhomogeneous problem. This, however, is only an approximate approach to the nonlinear stress - strain and stress - permeability behavior.

Permeability is known to vary with effective stress in a complicated manner. A widely used law is the linear relationship between the void ratio and logarithm of permeability. Raymond (1966 & 1969), Poskitt (1969), Poskitt and Birdsall (1971), Berry and Poskitt (1972) have used this relation in their theories. Other relationships are also proposed by Schmid (1957), Barden and Berry (1965),





Hwang (1966), Samarasinghe *et al* (1982), Juarez-Badillo (1983). Schiffman (1958) and Koppula (1970) are special cases of above.

For the stress range normally encountered in practice, the usually accepted compressibility law is the linear relationship between void ratio and logarithm of vertical effective stress. Barden and Berry (1965), Davis and Raymond (1965), Mikasa (1965), Raymond (1966), Poskitt (1969), Wilson and Hwang (1969), all used this relation in developing their analyses.

It is found in some soils that the coefficient of consolidation varies much less than either the permeability or the compressibility. Researchers such as Janbu (1965), Davis and Raymond (1965), Mikasa (1965), Davis (1971) assume this coefficient to be a constant in their formulations. The assumptions of constant coefficient of consolidation and the logarithmic effective stress - void ratio relation have also been shown from laboratory tests for normally consolidated clay, and where the secondary effect is ideally absent.

Many of these theories, such as Barden and Berry (1965), Davis and Raymond (1965), Raymond (1966), have introduced nonlinear behaviour in thin layered soils although the restraint of small strain still persists.



Small strain analysis is based on the assumption that the deformation is small compared to the thickness of the consolidating layer. Therefore, it is not appropriate for an extremely compressible soil and only applicable to oedometer tests or soil samples sufficiently thin that self-weight effects can be neglected.

The usual coordinate system employed in small strain theory is the Eulerian system where all measurements are referenced in space to a fixed datum. Settlement computation based on small strain theory, therefore, may probably be in error since the movement of boundaries is assumed in a fixed coordinate system.

Finite strain analyses for thin layers have been developed by McNabb (1960), Mikasa (1965), Gibson *et al* (1967), Poskitt (1969), Wilson and Hwang (1969), Mesri and Rockhsar (1974), Lee and Sills (1979). The first presents the fundamental nonlinear equations without providing any solution, while the last uses a moving boundary approach to account for large strains. Mikasa assumed a constant coefficient of consolidation and used the compression strain as the primary unknown in his governing equation. Settlement analysis is given by Wilson and Hwang for different load ratios and permeability parameters. Mesri and Rockhsar have carried out an extensive study of nonlinear behavior, but failed to account for the influence



of soil depth.

The general theory formulated by Gibson *et al* (1967) reduced many limitations in the conventional analysis. It represents a basis for various future researchers. Studies using thin layer approximation have been made by these authors and other researchers. Their original equation with logarithmic material laws and further assumptions is solved analytically by Poskitt (1969) using a perturbation method. A closed-form solution is found by Simons and Beng (1969) with other approximations. Comparison with experimental results is presented by Burland and Roscoe (1969).

Experimental investigation and numerical solution of Gibson *et al*'s (1967) general formulation is suggested by Monte and Krizek (1976). In their analysis, the liquid limit state is taken as the datum for strain measurement. Linear relationship between constrained modulus and strain is employed, and the permeability is expressed as a second degree function of void ratio.

In engineering projects such as large embankments, reclamation filling and groundwater lowering, very deep and soft clay layers may be involved. Large settlements may occur and effects of self-weight cannot be ignored.





Raymond (1969) and Davis (1971) have considered effects of layer depth in the context of small strain. Numerical solutions for Davis (1971) is suggested by Viggiani (1973). Variable load increment is also considered by Raymond (1969) and Viggiani (1973).

Studies accounting for both the finite strain and self-weight effect are given by Janbu (1965) and Gibson *et al* (1981). The former gives a general formulation similar to Gibson *et al* (1967), and particular cases are considered. The latter uses the same general formulation as in an earlier paper (Gibson *et al*, 1967). Numerical solution for Gibson *et al* (1981) is provided by DeSimone and Viggiani (1976) using assumptions of constant coefficient of consolidation and logarithmic compressibility law.

Despite the fact that Gibson *et al*'s (1967, 1981) governing equation allows the general consideration of many factors influencing consolidation, other approximations are incorporated in order to arrive at a solution. This is a serious drawback to most extensive nonlinear theories, where mathematical difficulties usually arise for a general solution.

The advantages of ADINAT in such a situation is apparent since its finite element procedure, accompanied with the nonlinear solution strategies, provide a powerful





tool for solving any form of expression for material laws. Since different soils are subject to different forms of material laws, ADINAT allows the choice of the most proper constitutive relations based on experimental evidence rather than being restricted to any particular form due to analytical convenience.

#### 4.2 Derivation of the General Governing Equation

The one dimensional consolidation of a simple clay layer with drainage at both top and bottom is used for the illustration here. (Gibson *et al*, 1967) The soil profile and coordinate systems employed are given in Fig. 4.1. Fig. 4.1(a) and (b) show the Lagrangian system and refers all events to the initial (zero time) state. The convective coordinate, on the other hand, always moves with the actual material configuration. This is described in Fig. 4.1(c) and (d).

To further clarify the following derivation, a pictorial description of the stress distribution in the soil layer in the convective coordinate system is shown in Fig. 4.2. In the following derivation, the first seven assumptions stated in Section 3.1 are still valid.



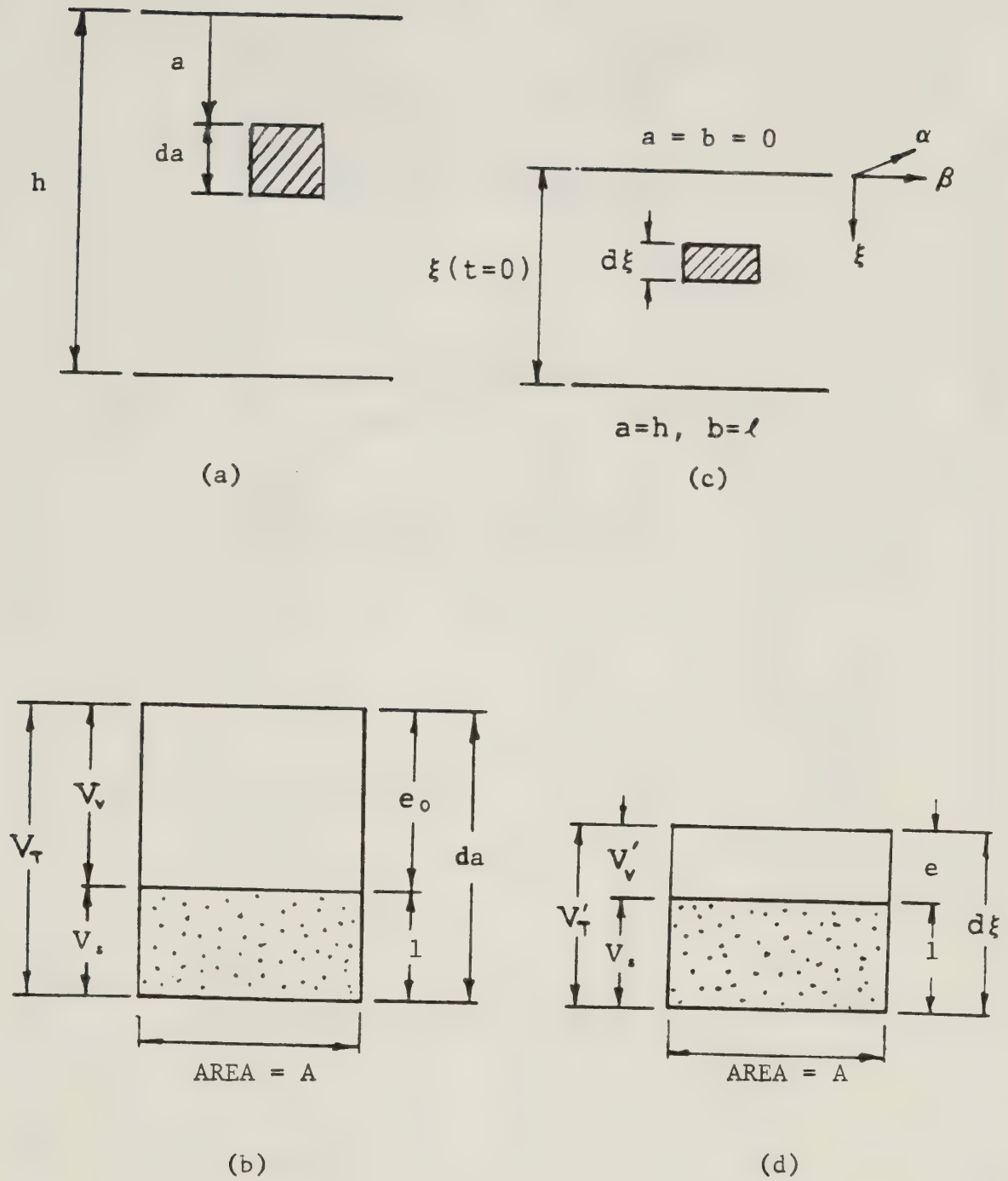


Figure 4.1 Coordinate system: (a),(b) Lagrangian: initial state, (c),(d) Convective: consolidating state



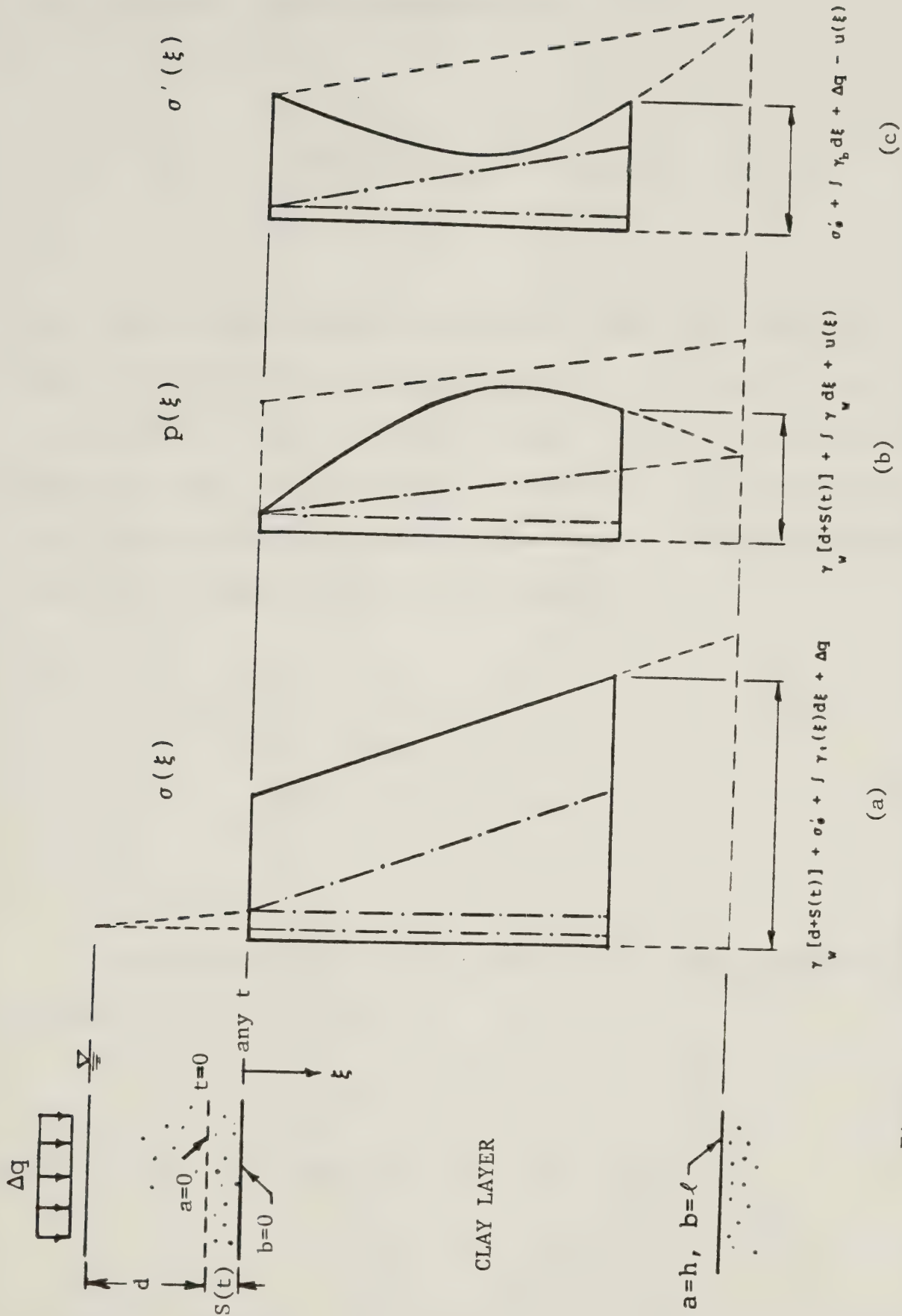


Figure 4.2 Stress and pore pressure distribution in convective coordinate



Referring to the convective system  $\xi$  in Fig. 4.2, the total stress  $\sigma$  can be expressed as

$$\sigma(\xi) = \gamma_w [d+S(t)] + \sigma'_0 + \int \gamma_t(\xi) d\xi + \Delta q \quad (4.1)$$

$d$  is the distance between the water table and the soil layer,  $S(t)$  is the settlement at time  $t$ ,  $\gamma_t$  and  $\gamma_w$  are the unit weight of soil and water respectively. Differentiating partially with respect to  $\xi$ ,  $\xi = \xi(a,t)$ , and transforming into the Lagrangian system by the chain rule, a vertical equilibrium equation can be obtained

$$\frac{\partial \sigma}{\partial a} - \frac{\gamma_s + e\gamma_w}{1+e} \cdot \frac{\partial \xi}{\partial a} = 0 \quad (4.2)$$

The expression for pore pressure in the convective system from Fig. 4.2 is

$$p(\xi) = \gamma_w [d+S(t)] + \int \gamma_w d\xi + u(\xi) \quad (4.3)$$

Similarly as above, equilibrium of the pore fluid can be





determined,

$$\frac{\partial u}{\partial a} \cdot \frac{\partial a}{\partial \xi} = \frac{\partial p}{\partial a} \cdot \frac{\partial a}{\partial \xi} - \gamma_w \quad (4.4)$$

The net inflow rate of fluid weight is

$$q_{\xi} - (q_{\xi} + dq_{\xi}) = - \frac{\partial}{\partial \xi} \left[ (v - v_s) \frac{e\gamma_w}{1+e} \right] \cdot d\xi d\alpha d\beta \quad (4.5)$$

and the rate of change of fluid weight can be given as

$$\frac{\partial}{\partial t} (W_w) = \frac{\partial}{\partial t} \left[ \frac{e\gamma_w}{1+e} \cdot d\xi d\alpha d\beta \right] \quad (4.6)$$

By continuity, equating the above two equations and putting

$$\frac{\partial \xi}{\partial a} = \frac{1+e(a,t)}{1+e(a,0)} \quad (4.7)$$



gives

$$\frac{\partial}{\partial a} \left[ (v_w - v_s) \frac{e\gamma w}{1+e} \right] = - \frac{\partial}{\partial t} \left[ \frac{e\gamma w}{1+e} \cdot \frac{\partial \xi}{\partial a} \right] \quad (4.8)$$

By Darcy's law, the fluid flux is proportional to the excess pressure gradient as

$$(v_w - v_s) \frac{e\gamma w}{1+e} = -k \frac{\partial u}{\partial a} \frac{\partial a}{\partial \xi} \quad (4.9)$$

Substituting Eq. 4.4 and Eq. 4.8 into the above equation, one obtains

$$\frac{\partial}{\partial a} \left[ k \left( \gamma_w - \frac{\partial p}{\partial a} \cdot \frac{\partial a}{\partial \xi} \right) \right] = - \frac{\partial}{\partial t} \left[ \frac{e\gamma w}{1+e} \cdot \frac{\partial \xi}{\partial a} \right] \quad (4.10)$$

From Eq. 4.2, 4.7, 4.10 and the effective stress principle,



$$\sigma = \sigma' + p \quad (4.11)$$

gives

$$\frac{\partial}{\partial a} \left\{ \frac{k}{1+e} \left[ (1+e_0) \frac{\partial \sigma'}{\partial a} - (\gamma_s - \gamma_w) \right] \right\} = - \frac{\gamma_w}{1+e_0} \cdot \frac{\partial e}{\partial t} \quad (4.12)$$

Note that  $e_0$  is not a constant but varies throughout the soil depth. Eq. 4.12 is a general governing equation with effective stress and void ratio as the primary unknowns.

### 4.3 Overview of ADINAT Application

#### 4.3.1 General

One dimensional nonlinear theories can be classified into four main categories: small strain - thin layer, small strain - thick layer, finite strain - thin layer and finite strain - thick layer analyses. Applications of ADINAT to each category will be illustrated.

Mathematically, small strain requires the volume at any time  $t$ ,  $(1+e)$ , is equal to the initial volume  $(1+e_0)$ . Thin



layer analysis does not allow for the total stress variation with depth due to the self-weight of the soil mass, which means  $\gamma_s = \gamma_w$ . This first and second limitation individually is not imposed on the finite strain and thick layer analysis respectively.

General formulations used in ADINAT for each of these categories are derived below. Their uses will be discussed latter. Examples are provided in the following chapters.

#### 4.3.2 Formulations of *ADINAT* Equations

##### (a) Small Strain and Thin Layer Analysis

As discussed previously,  $(1+e)$  equals  $(1+e_0)$  and  $\gamma_s$  equals  $\gamma_w$  are presumed for small strain and thin layer analysis. Further, the Lagrangian coordinate  $a$  can be replaced by the Eulerian coordinate  $z$ , since the initial configuration holds for all time. Using these simplifications in Eq 4.12,

$$\frac{\partial}{\partial z} \left[ k \frac{\partial \sigma'}{\partial z} \right] = - \frac{\gamma_w}{1+e_0} \frac{\partial e}{\partial t} \quad (4.13)$$





and further assuming

$$\sigma' = \sigma(e) \quad (4.14)$$

gives

$$\frac{\partial}{\partial z} \left[ k \frac{\partial \sigma'}{\partial z} \right] = - \frac{\gamma w}{1+e_0} \frac{de}{d\sigma'} \frac{\partial \sigma'}{\partial t} \quad (4.15)$$

or in a different form

$$\frac{\partial}{\partial z} \left[ - \frac{k}{\gamma} \frac{d\sigma'}{de} \frac{\partial e}{\partial z} \right] = \frac{1}{1+e_0} \frac{\partial e}{\partial t} \quad (4.16)$$

<sub>w</sub>

which can be rewritten as

$$\frac{\partial}{\partial z} \left\{ \frac{k}{\gamma} \left[ - (1+e_0) \frac{d\sigma'}{de} \right] \frac{\partial e}{\partial z} \right\} = \frac{\partial e}{\partial t} \quad (4.17)$$

<sub>w</sub>

Since



$$u(z,t) = \sigma'_0 + \Delta q - \sigma'(z,t) \quad (4.18)$$

Differentiating partially with respect to  $z$  and  $t$  gives

$$\frac{\partial u}{\partial z} = - \frac{\partial \sigma'}{\partial z} \quad (4.19)$$

and

$$\frac{\partial u}{\partial t} = - \frac{\partial \sigma'}{\partial t} \quad (4.20)$$

Substituting these equations in Eq. 4.15 and Eq. 4.17, together with the conventional definitions of  $C_v$ , Eq. 3.7, and  $m_v$ , Eq. 3.8, result in

$$\frac{\partial}{\partial z} \left[ k \frac{\partial u}{\partial z} \right] = \gamma_w m_v \frac{\partial u}{\partial t} \quad (4.21)$$

and



$$\frac{\partial}{\partial z} \left[ C_v \frac{\partial e}{\partial z} \right] = \frac{\partial e}{\partial t} \quad (4.22)$$

### (b) Small Strain and Thick Layer Analysis

Assume small strain,  $(1+e) \approx (1+e_0)$  and thus replace  $a$  by the Eulerian coordinate  $z$  in Eq. 4.12, which can be rewritten as

$$\frac{\partial}{\partial z} \left\{ k \frac{\partial \sigma'}{\partial z} - \frac{k}{1+e_0} (\gamma_s - \gamma_w) \right\} = - \frac{\gamma_w}{1+e_0} \frac{\partial e}{\partial t} \quad (4.23)$$

In thick layer analysis using the Eulerian coordinate, self-weight of soil mass is to be accounted for in the effective stress relationship as

$$\sigma'(z,t) = \sigma'_0 + \Delta q + \int_w (\gamma_s - \gamma_w) dz - u(z,t) \quad (4.24)$$

or



$$\sigma'(z,t) = \sigma'_\phi + \Delta q + (\gamma_s - \gamma_w) \cdot \int \frac{dz}{1+e_0} - u(z,t) \quad (4.25)$$

The subscript  $\phi$  refers to the initial value at soil surface.  
Differentiating partially with respect to  $z$  and  $t$  give

$$\frac{\partial \sigma'}{\partial z} = \frac{\gamma_s - \gamma_w}{1+e_0} - \frac{\partial u}{\partial z} \quad (4.26)$$

and

$$\frac{\partial \sigma'}{\partial t} = - \frac{\partial u}{\partial t} \quad (4.27)$$

Substituting Eq. 4.14, 4.26 and 4.27 into Eq. 4.23, gives

$$\frac{\partial}{\partial z} \left[ k \frac{\partial u}{\partial z} \right] = \gamma_w m_v \frac{\partial u}{\partial t} \quad (4.28)$$





## (c) Finite Strain and Thin Layer Analysis

For finite strain and thin layer formulation, setting  $\gamma_s \approx \gamma_w$  in Eq. 4.12 obtains

$$\frac{\partial}{\partial a} \left[ k \frac{1+e_0}{1+e} \frac{\partial \sigma'}{\partial a} \right] = - \frac{\gamma_w}{1+e_0} \frac{\partial e}{\partial t} \quad (4.29)$$

Assuming Eq. 4.14 and using the definitions of Eq. 3.7 and Eq. 3.8, Eq. 4.29 can be expressed as

$$\frac{\partial}{\partial a} \left[ k \frac{1+e_0}{1+e} \frac{\partial \sigma'}{\partial a} \right] = \gamma_w \frac{mv(1+e)}{1+e_0} \frac{\partial \sigma'}{\partial t} \quad (4.30)$$

or

$$\frac{\partial}{\partial a} \left[ \frac{Cv(1+e_0)}{(1+e)^2} \frac{\partial e}{\partial a} \right] = \frac{1}{1+e_0} \frac{\partial e}{\partial t} \quad (4.31)$$

Since



$$u(a,t) = \sigma'_0 + \Delta q - \sigma'(a,t) \quad (4.32)$$

in thin layer analysis.  $\sigma'_0$  refers to the initial effective stress in soil. Differentiating partially with respect to  $a$  and  $t$  similar to Section (a), and substituting into Eq. 4.30, gives:

$$\frac{\partial}{\partial a} \left[ k \frac{1+e_0}{1+e} \frac{\partial u}{\partial a} \right] = \gamma_w \frac{mv(1+e)}{1+e_0} \frac{\partial u}{\partial t} \quad (4.33)$$

#### (d) Finite Strain and Thick Layer Analysis

To account for both large deformation and self-weight in finite strain and thick layer formulation, the convective system should be considered. The effective stress relation is

$$\sigma'(\xi) = \sigma'_0 + \Delta q + \int_w (\gamma_t - \gamma_w) d\xi - u(\xi) \quad (4.34)$$

or it can be determined in Lagrangian coordinate through



$$\sigma'(a,t) = \sigma'_0 + \Delta q + (\gamma_s - \gamma_w) \cdot \int \frac{da}{1+e_0} - u(a,t) \quad (4.35)$$

which gives

$$\frac{\partial \sigma'}{\partial a} = \frac{\gamma_s - \gamma_w}{1+e_0} - \frac{\partial u}{\partial a} \quad (4.36)$$

and

$$\frac{\partial \sigma'}{\partial t} = - \frac{\partial u}{\partial t} \quad (4.37)$$

Substituting Eq. 4.36 into Eq. 4.12 gives:

$$\frac{\partial}{\partial a} \left[ - \frac{k(1+e_0)}{1+e} \frac{\partial u}{\partial a} \right] = - \frac{\gamma_w}{1+e_0} \frac{\partial e}{\partial t} \quad (4.38)$$

Substituting Eq. 4.14 and 4.37 into Eq. 4.38 and rearranging to obtain



$$\frac{\partial}{\partial a} \left[ k \frac{1+e_0}{1+e} \frac{\partial u}{\partial a} \right] = \gamma \frac{mv(1+e)}{w(1+e_0)} \frac{\partial u}{\partial t} \quad (4.39)$$

Since  $(1+e_0)$  varies with depth, another reduced coordinate  $b$  may be used. This coordinate may be physically interpreted as the volume of soil solids lying between the datum and the point being considered, that is (McNabb, 1960)

$$b(a) = \int \frac{da}{1+e(a,0)} \quad (4.40)$$

or

$$\frac{db}{da} = \frac{1}{1+e_0} \quad (4.41)$$

The mathematics of the effective stress relation now becomes a little simpler,

$$\sigma'(b,t) = \sigma'_0 + \Delta q + (\gamma_s - \gamma_w) b - u(b,t) \quad (4.42)$$

Substitute Eq. 4.41 into Eq. 4.39 produces,





$$\frac{\partial}{\partial b} \left[ \frac{k}{1+e} \frac{\partial u}{\partial b} \right] = \gamma \frac{m}{w v} (1+e) \frac{\partial u}{\partial t} \quad (4.43)$$

In reduced coordinates, the soil layer thickness is reduced to  $\ell$ , which is obtained from the integral, Eq. 4.40, evaluated throughout the depth.

#### 4.3.3 Remark on the Use of *ADINAT* Equations

A summary table of the above *ADINAT* equations is given in Fig. 4.3. The equations are renumbered from *A1* to *A9* in the figure for the convenience of later reference. Equations *A1* and *A2* are Eq. 4.21 and Eq. 4.22 respectively. *A3* is an obvious outcome from Eq. 4.21 when the parameter  $m$  is assumed constant. *A4* refers to Eq. 4.28 while *A5* is a direct result by assuming  $m$  constant. *A6*, *A7*, *A8*, *A9* are corresponding to Eq. 4.33, 4.31, 4.39, 4.43 respectively.

Notice that even if the forms of equation *A1* and *A4* are similar, their uses are different. For the application of *ADINAT*, *A4* must be evaluated as described in Chapter 6.

The purpose of Fig. 4.3 is to facilitate the users study of nonlinear consolidation using *ADINAT*. The



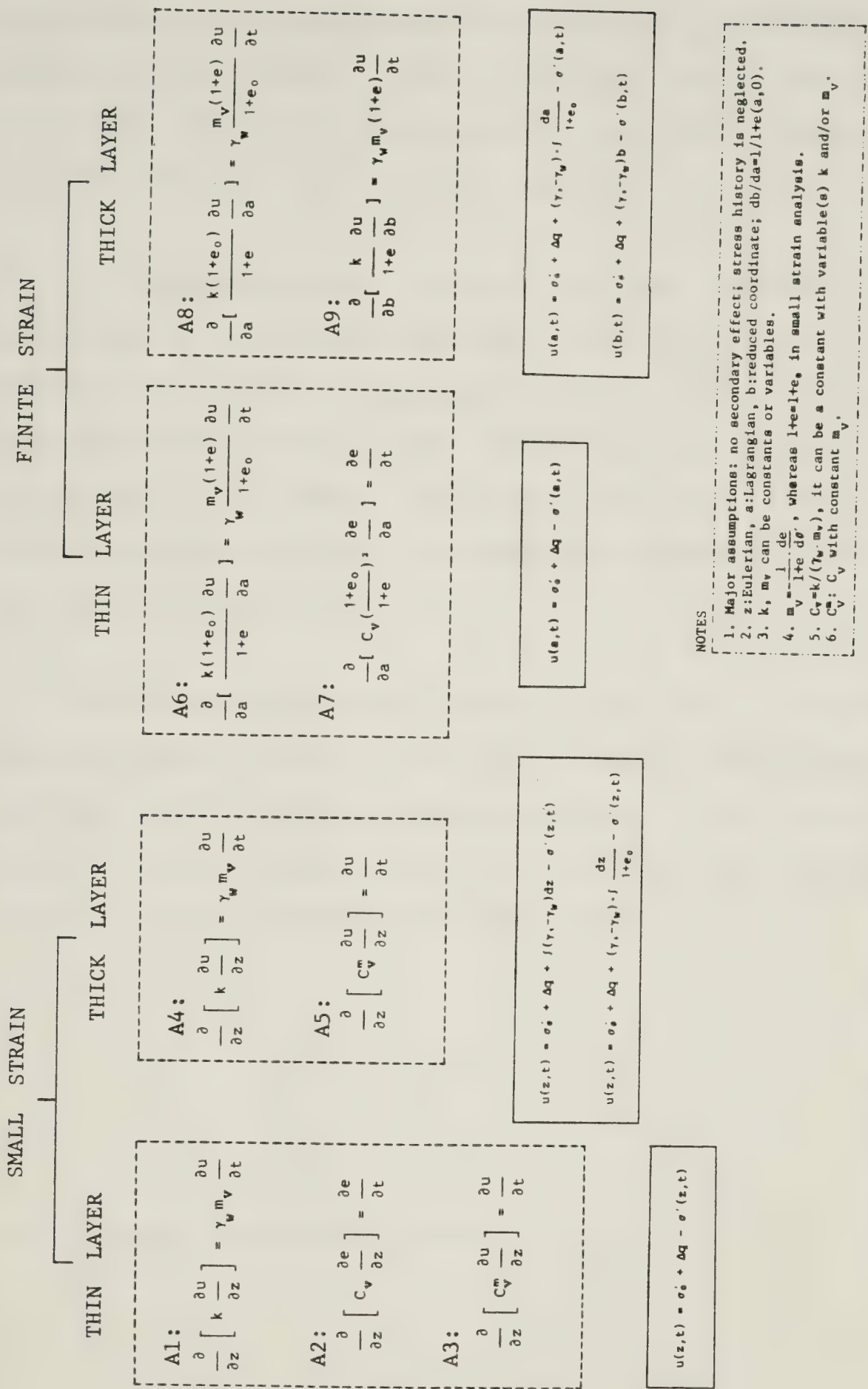


Figure 4.3 General formulations for ADINAT applications in nonlinear consolidation



equations given in each category, although in different forms, are equivalent. This will further assist the use of the table. The table can be used in many ways. Examples will be shown in the succeeding chapters.

In these equations,  $C_v$  and  $m_v$  are just notations presented in the conventional way. They are different from the mean values  $\overline{C_v}$ ,  $\overline{m_v}$  obtained from the conventional tests, and can be substituted by any appropriate expressions according to their basic definitions, Eq. 3.7 and 3.8. Unless otherwise stated, the material parameters can be constants or variables.

Since the logarithmic material relations are so widely accepted that a table of the finite strain - thin layer analysis with this approximation is given in Fig. 4.4. Several letter symbols adopted in the figure are defined below.  $\alpha$  is the commonly used load ratio,

$$\alpha = \sigma'_f / \sigma'_0 \quad (4.44)$$

and  $\beta$  is the permeability ratio,



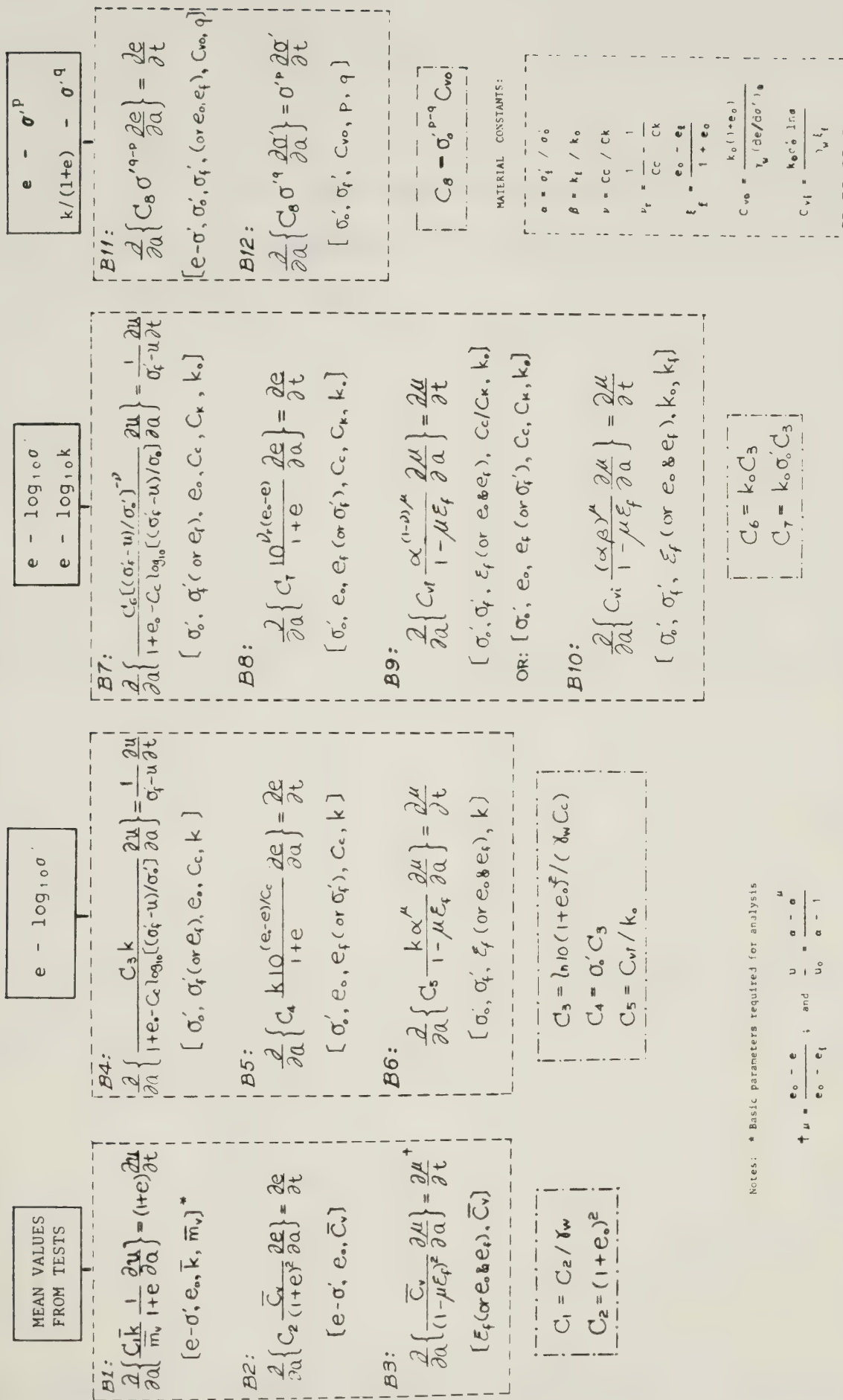


Figure 4.4 Useful equations for ADINAT applications in finite strain, thin layer consolidation





$$\beta = k_f / k_o \quad (4.45)$$

This ratio is always less than 1 in monotonic compression as in the illustrated examples. The constant  $\nu$  is a ratio of the material coefficients

$$\nu = C_c / C_k \quad (4.46)$$

where  $C_c$  is the slope of the  $e\text{-}\log_{10}\sigma'$  line, and  $C_k$  is the slope of the  $e\text{-}\log_{10}k$  line.  $\nu_r$  is defined in a similar way:

$$\nu_r = \frac{1}{C_c} - \frac{1}{C_k} \quad (4.47)$$

$\mu$  is termed local consolidation ratio and with a value ranges from 0 to 1. It is given by

$$\mu = \frac{e_o - e}{e_o - e_f} \quad (4.48)$$

The application of this variable simplifies the equations as



well as their use, as may be seen from Fig. 4.4. If the logarithmic effective stress - void ratio relation is assumed, the excess pore pressure ratio can be determined from  $\mu$  as follows

$$\frac{u}{u_0} = \frac{\alpha - \alpha^\mu}{\alpha - 1} \quad (4.49)$$

The total final strain,  $\xi_f$ , is expressed in the form of

$$\xi_f = \frac{e_0 - e_f}{1 + e_0} \quad (4.50)$$

$C_{vi}$  is the initial consolidation ratio defined as

$$\begin{aligned} C_{vi} &= k_0 \sigma'_0 (1+e_0) \ln 10 / \gamma_w C_c \\ &= k_0 \sigma'_0 \ln \alpha / \gamma_w \xi_f \end{aligned} \quad (4.51)$$

Another approximation utilizing a more general material law used by Janbu (1963) and Koppula (1970) is



$$\frac{de}{d\sigma'} = \left[ \frac{de}{d\sigma'} \right]_1 \left( \frac{\sigma'}{\sigma'_0} \right)^p \quad (4.52)$$

and

$$\frac{k}{1+e} = \left[ \frac{k}{1+e} \right]_1 \left( \frac{\sigma'}{\sigma'_0} \right)^q \quad (4.53)$$

These are also given in Fig. 4.4. The constants  $p$  and  $q$  range between  $-1$  and  $0$ , the subscript  $1$  denotes a reference state. The logarithmic constitutive relations just mentioned are special cases of these equations.

Since average values of material parameters are often obtained from conventional laboratory tests, the formulations using these parameters are also included in Fig. 4.4. The values of  $\bar{k}$ ,  $\bar{m}_v$  and  $\bar{C}_v$  are assumed to be obtained independently from tests. Note that  $\bar{C}_v$  may not be equal to  $\bar{k}/(\gamma_w \cdot \bar{m}_v)$ .

For the same reason as above, sets of different forms but mathematically equivalent equations are given in Fig. 4.4. Derivations of these equations are long but similar to those in Fig. 4.3 and will not be shown. Equations in the figure are numbered from  $B1$  to  $B12$  for



future reference.

All of these *ADINAT equations* in Fig. 4.3 and Fig. 4.4 are presented in a clear manner for ADINAT input, and are given in the following form:

$$\frac{\partial}{\partial x} \left[ K \frac{\partial \theta}{\partial x} \right] = M \frac{\partial \theta}{\partial t} \quad (4.54)$$

where  $x$ ,  $\theta$ ,  $K$ ,  $M$  represents the coordinate, unknown, permeability, compressibility variable respectively. To illustrate, whatever expression is entered in the *position* 'K' will be input as the permeability variable in the program. The same applies to the rest of the parameters.

Equations with an unknown variable  $\sigma'$  are primary for obtaining pore pressure results, while those with an unknown  $e$  are mainly for computing deformations. Both of these, obviously, can be used for either purpose with supplemental programs to reduce the original output.





#### 4.4 Discussion

It can be seen from Fig. 4.3 that pore pressure can still be maintained as the primary unknown. Hence the strain or void ratio is no longer a privileged variable in nonlinear analysis, as is believed by some researchers. (Mikasa, 1965; Janbu, 1965; Gibson *et al*, 1967 & 1981)

The difficulty with the general governing equation, Eq. 4.39, is that the settlement cannot be assessed directly. On the other hand, Gibson *et al*'s (1981) equation may be useful for computing soil settlement, if a solution can be easily obtained.

Regarding the use of the general equation, Eq. 4.39, inconvenience may seem to arise in the computation of soil thickness  $\ell$ , or the evaluation of the integral Eq. 4.40. It is believed that this cannot be avoided in any analyses dealing with both finite strain and self-weight effects. These two quantities, however, are readily and accurately computed by simple computer-based techniques.

The general governing equation A8 or A9 (Fig. 4.3) is more useful than Eq. 4.12 since solutions are readily obtained by ADINAT for the former, while no general solution is yet available for the latter. Therefore, A8 or A9 represents the working governing equation for a general



treatment of one dimensional consolidation. Note again that both of these equations are basically equivalent to Eq. 4.12.

As mentioned before, most theories can be recovered from the proposed ADINAT formulations. Examples are given below.

Terzaghi's classical theory is an obvious outcome of ADINAT A1 when  $k$  and  $m$  are assumed constant. The governing equation found by Davis and Raymond (1965) is equivalent to A2 with constant  $C_v$ . Barden and Berry (1965), Raymond (1966) are particular cases of A1 where logarithmic material relations are used.

Analyses of soils with large depth by Schiffman and Gibson (1964), Raymond (1969), Davis (1971), Viggiani (1973) are special cases of ADINAT small strain - thick layer formulation, A4. In comparison with A8, Schiffman and Gibson's (1964) general equation is still deficient in the context of thick layer formulation.

McNabb (1960) has given an equation similar to A7. The equation given by Poskitt (1969) can be obtained from A6 with the logarithmic material properties and change of variables. It is identical to equation B10 in Fig. 4.4. Gibson *et al* (1967), Simons and Beng (1969), Wilson and



Hwang (1969) are particular cases of A7. Notice that the parameter  $C_v$  mentioned in the paper by Gibson *et al* is actually a mean value  $\overline{C_v}$  in which  $(1+e) \approx (1+e_0)$  is implied.

The general governing equation Eq. 13 in Gibson *et al* (1967) or Eq. 19 in Gibson *et al* (1981) is basically identical to A8 or A9 respectively. Solutions are provided for A8 and A9 using ADINAT but no general solution has yet found for Gibson *et al*'s equation. Numerical solutions are given, however, only for an approximate form of the latter. (DeSimone and Viggiani, 1976; Gibson *et al*, 1981)



## 5. APPLICATION OF ADINAT IN NONLINEAR CONSOLIDATION OF THIN SAMPLE

### 5.1 General

In order to verify the proposed equations in Fig. 4.3 and show their applications, a series of examples, most of which are taken from published papers, are presented. Thin layer formulations, which most of these papers considered, are discussed in this chapter. Applications of thick layer theories will be given in the next chapter.

Another purpose of these examples is to show the correlation among the mathematical equivalent counterparts in the same category. Therefore, as many as possible will be used in the same example. Applications of the equations in Fig. 4.4 will also be illustrated as well.

ADINAT equations in Fig. 4.3 and Fig. 4.4 are referred to as  $A1$ , ... and  $B1$ , ... in the following examples. It is worthwhile to note again that equations in Fig. 4.3 are *general* equations. That is, they can be used for any material laws. Equations in Fig. 4.4 are special cases of the finite strain - thin layer formulations in Fig. 4.3, with the assumed constitutive relations shown at the top of each sub-category.





In these two figures, equations in each set are mathematically equivalent and are presented in a simplified manner. The choice among them is a matter of preference or may depend on the primary output of the unknown variable.

For all examples illustrated here, unless otherwise specified, one dimensional linear finite element model is used, twenty equally spaced elements are assumed for the compressible layer. Ten solution time steps are assumed. Lumped compressibility matrix option and the Euler backward integration method are employed.



## 5.2 Examples of Samarasinghe-Huang-Drnevich Study

### Objective:

To show ADINAT application in Samarasinghe *et al* (1982) studies.

ADINAT Nonlinear Equation: A1, A6, B4, B5, B6

### Brief:

These authors have found that the permeability relation is determined by

$$k = Ck \frac{e^{\sigma}}{1+e} \quad (5.1)$$

where  $Ck$  is either a constant or it may vary with void ratio. They also assumed a common linear void ratio - logarithm effective stress relationship.



### Description and Analysis:

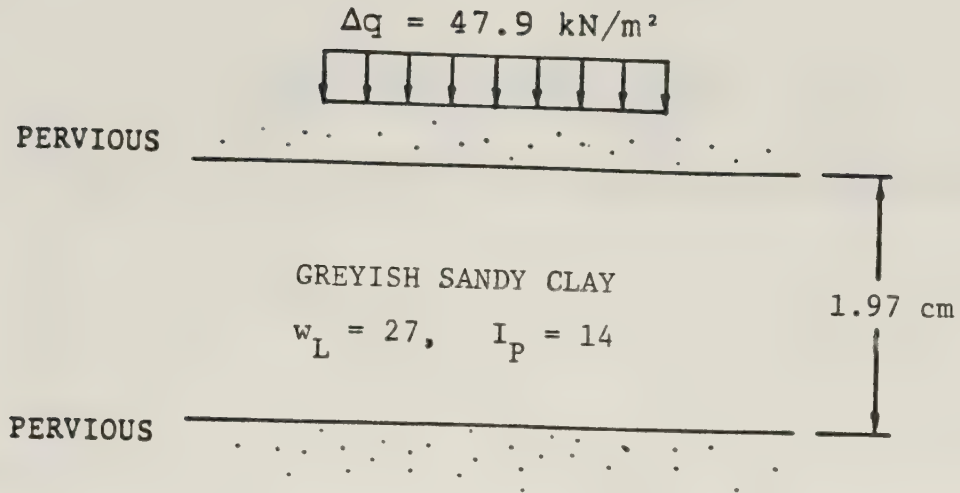
An example of greyish sandy clay with the exponent  $n=5.2$  and  $C_k=6.58 \times 10^{-4} \text{m/day}$  obtained from Samarasinghe *et al* is illustrated here. Fig. 5.1 shows the soil profile, loading and boundary conditions. Material properties and the finite element model are also given in the figure.

Both small and finite strain analyses are performed for comparison. Since a logarithmic compressibility relationship is included, equations in Fig. 4.4 can be employed. Datafiles are given in Appendix A.11.

### Results:

ADINAT prediction of pore pressure distribution at 50% consolidation is plotted in Fig. 5.2. Solutions with *A6* and *B4* are coincident since they are virtually the same equation. The small discrepancy among the results of *B4*, *B5* and *B6* may be due to the difference in numerical input into the program. The finite strain analysis in this example, is only little different from the small strain analysis.





PERMEABILITY MODEL:

$$k = Ck \frac{e^n}{1 + e}$$

$$Ck = 6.58 \times 10^{-4} \text{ m/day}$$

$$n = 5.2$$

COMPRESSIBILITY MODEL:

$$e - \log \sigma'$$

$$Cc = 0.16$$

OTHER PARAMETERS:  $e_o = 0.622$

$$\sigma'_o = \Delta q$$

$$\gamma_w = 9.81 \text{ kN/m}^3$$

FINITE ELEMENT MODEL FOR THE CLAY LAYER



20 EQUALLY SPACED ONE DIMENSIONAL ELEMENTS

Figure 5.1 One dimensional consolidation with nonlinear material properties: Samarasinghe et al, 1982





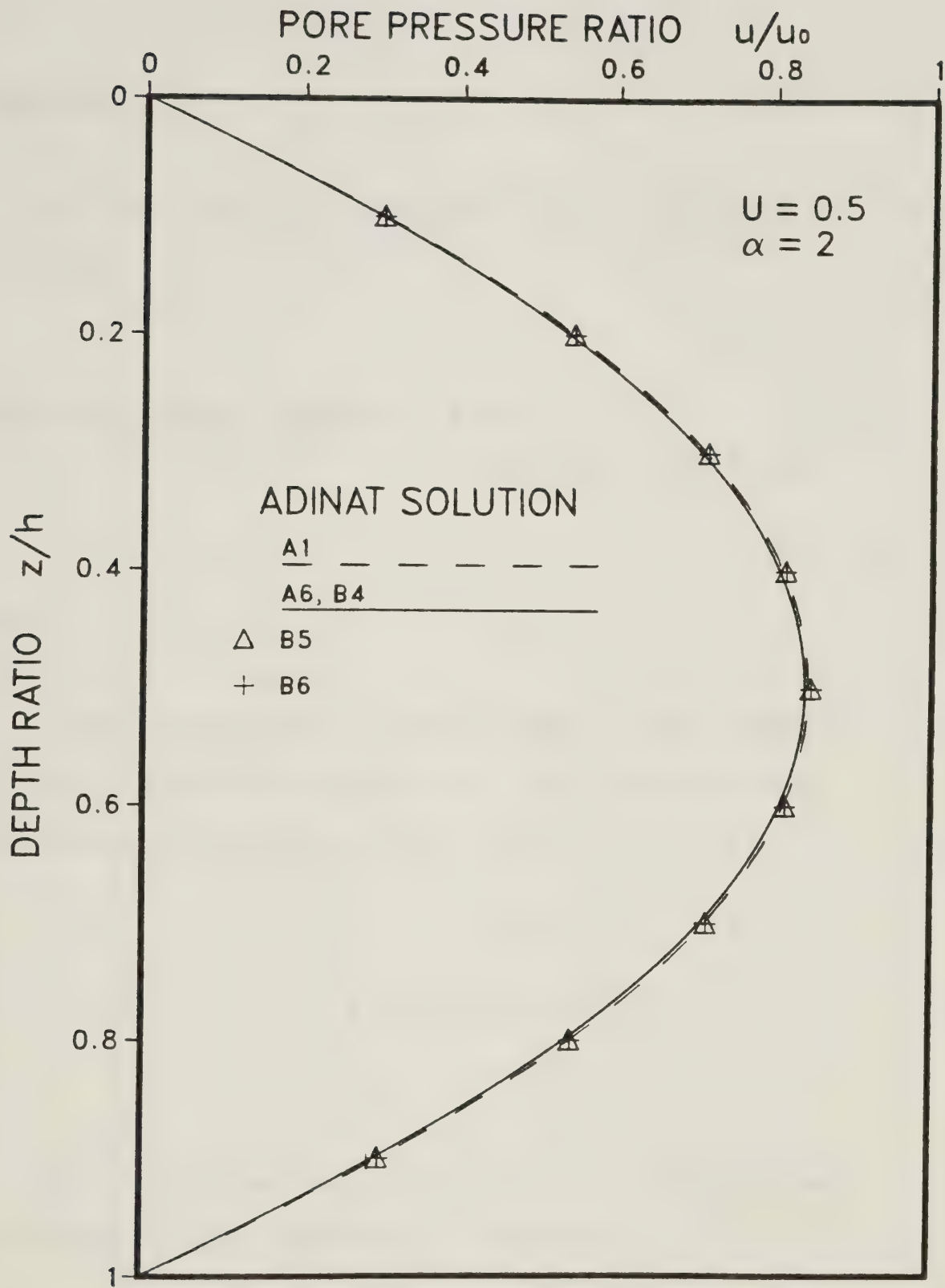


Figure 5.2 Pore pressure isochrones: thin layer analysis



### 5.3 Barden and Berry Analysis

#### Objective:

To show the use of ADINAT in a Barden and Berry (1965) analysis.

ADINAT Nonlinear Equation: (A) A1

(B) A1, B1, B2, B3, B6

#### Brief:

These authors have given a small strain theory utilizing logarithmic material laws. They further approximated the permeability relationship as

$$k = k_f (1 + \eta' u^n) \quad (5.2)$$

in which  $k_f$  is the value of  $k$  when  $u=0$ .  $\eta'$  and  $n$  are constants. Their equation is given as



$$\frac{\partial}{\partial z} \left[ \frac{k_f}{\gamma_w} (1 + bu^n) \frac{\partial u}{\partial z} \right] = \frac{.434Cc}{1 + e_o} \frac{1}{\sigma_f - u} \frac{\partial u}{\partial t} \quad (5.3)$$

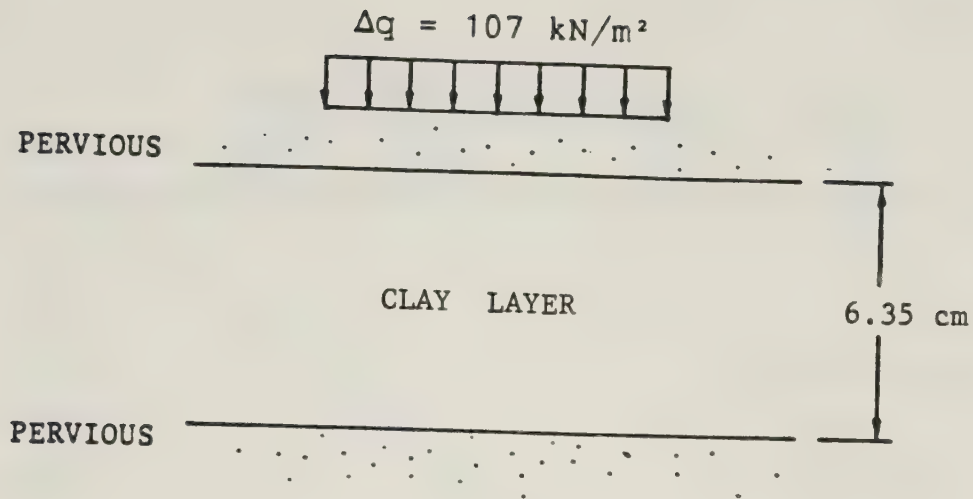
They resorted to finite difference method to solve the nonlinear equation generated.

### Description and Analysis:

(A) An illustration of Barden and Berry's formulation with an assumed soil profile and material properties as given in Fig. 5.3 is presented. Values of 1 and 1/2, as referred to by Barden and Berry, are used for the value of  $n$  in Eq. 5.2.

(B) Results of a laboratory test given by Barden and Berry is analysed by ADINAT. Both small strain and finite strain analyses are illustrated.  $B1$ ,  $B2$ ,  $B3$  and  $B6$  in the finite strain formulations are included for comparison. The loading and material properties are shown in Fig. 5.5.





#### PERMEABILITY MODEL:

$$k = k_f (1 + \eta' u^n)$$

$$\eta' = [(k_o/k_f) - 1]/u_o^n$$

$$k_o = 4.18 \times 10^{-5} \text{ m/day}$$

$$k_f = 2.09 \times 10^{-5} \text{ m/day}$$

$$n = 1/2 \text{ or } 1$$

#### COMPRESSIBILITY MODEL:

$$e - \log \sigma'$$

$$Cc = 0.4$$

#### OTHER PARAMETERS:

$$e_o = 0.959$$

$$\sigma'_o = \Delta q$$

$$\gamma_w = 9.81 \text{ kN/m}^3$$

#### FINITE ELEMENT MODEL FOR THE CLAY LAYER

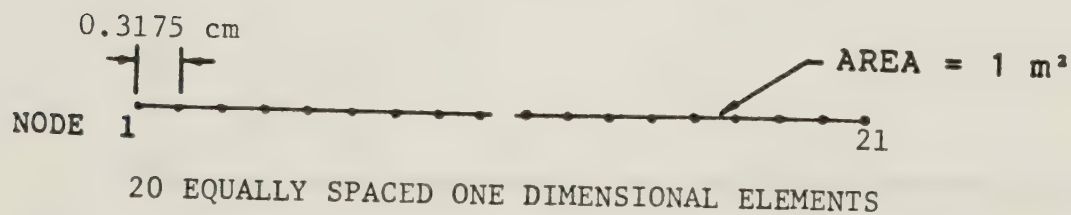


Figure 5.3 One dimensional consolidation with Barden-Berry (1965) formulation





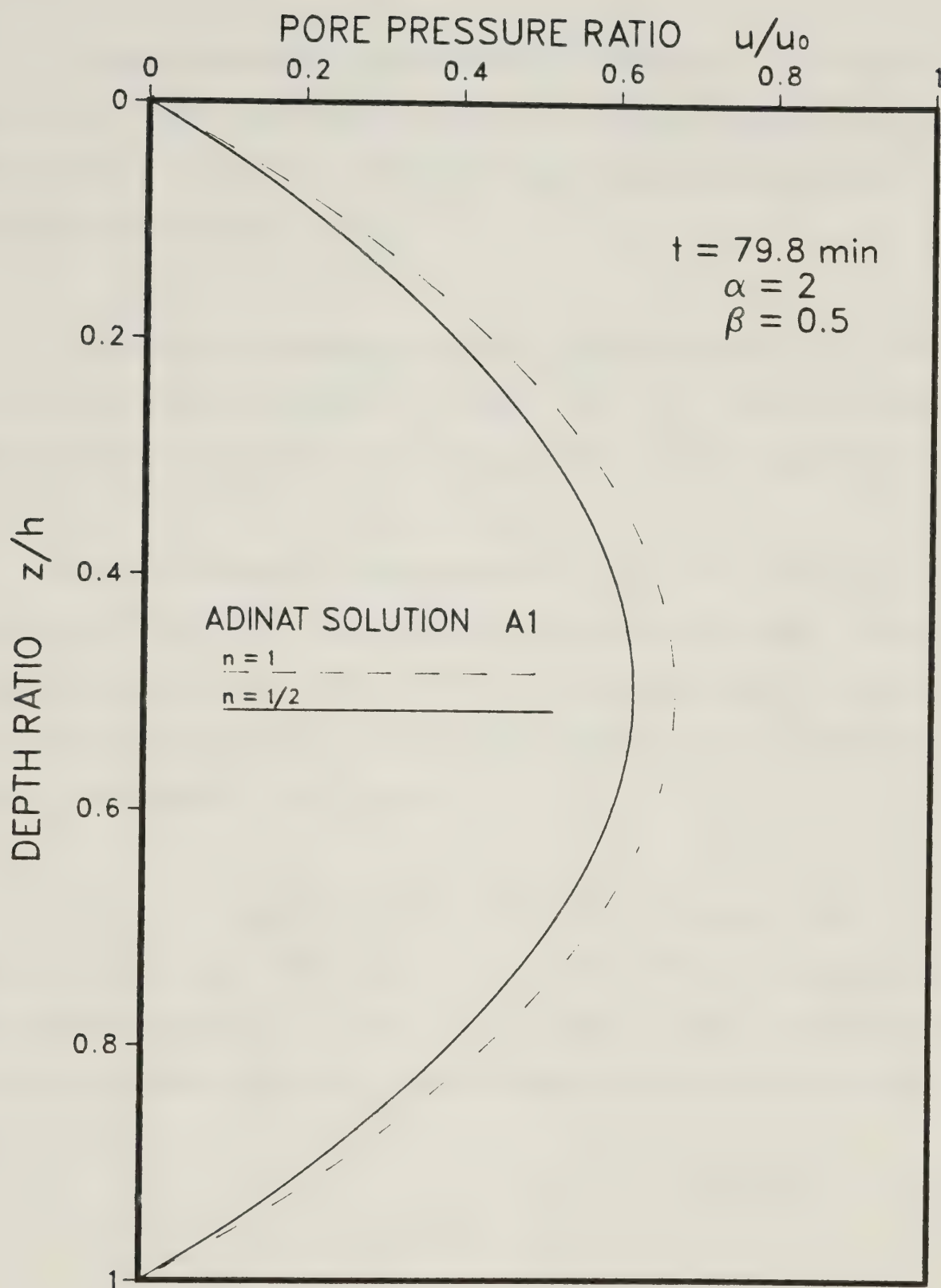


Figure 5.4 Pore pressure results in small strain - thin layer analysis



## Results

(A) Pore pressures predicted by ADINAT are plotted in Fig. 5.4. Not much difference can be seen between the cases of  $n$  equals 1 and  $1/2$ . This has also been observed by Barden and Berry as well.

(B) The pore pressure results given by ADINAT at 50% consolidation are plotted in Fig. 5.6. The finite strain solutions are found close to the small strain analysis.

The discrepancies among the finite strain results are due to the fact that different parameters are used. Note that the numerical value of  $\bar{C}_v$  used in equations B2 and B3 is not exactly equal to  $\bar{k}/(\gamma \cdot \bar{m})$ , whereas the parameters  $\bar{k}$  and  $\bar{m}$  are used in B1 and B6.

Fig. 5.7 compares ADINAT predictions, Barden and Berry's solutions and the test results. The experimental settlement curve is fitted to the ADINAT curve to obtain the required time, which in turn is used to determine the pore pressures.



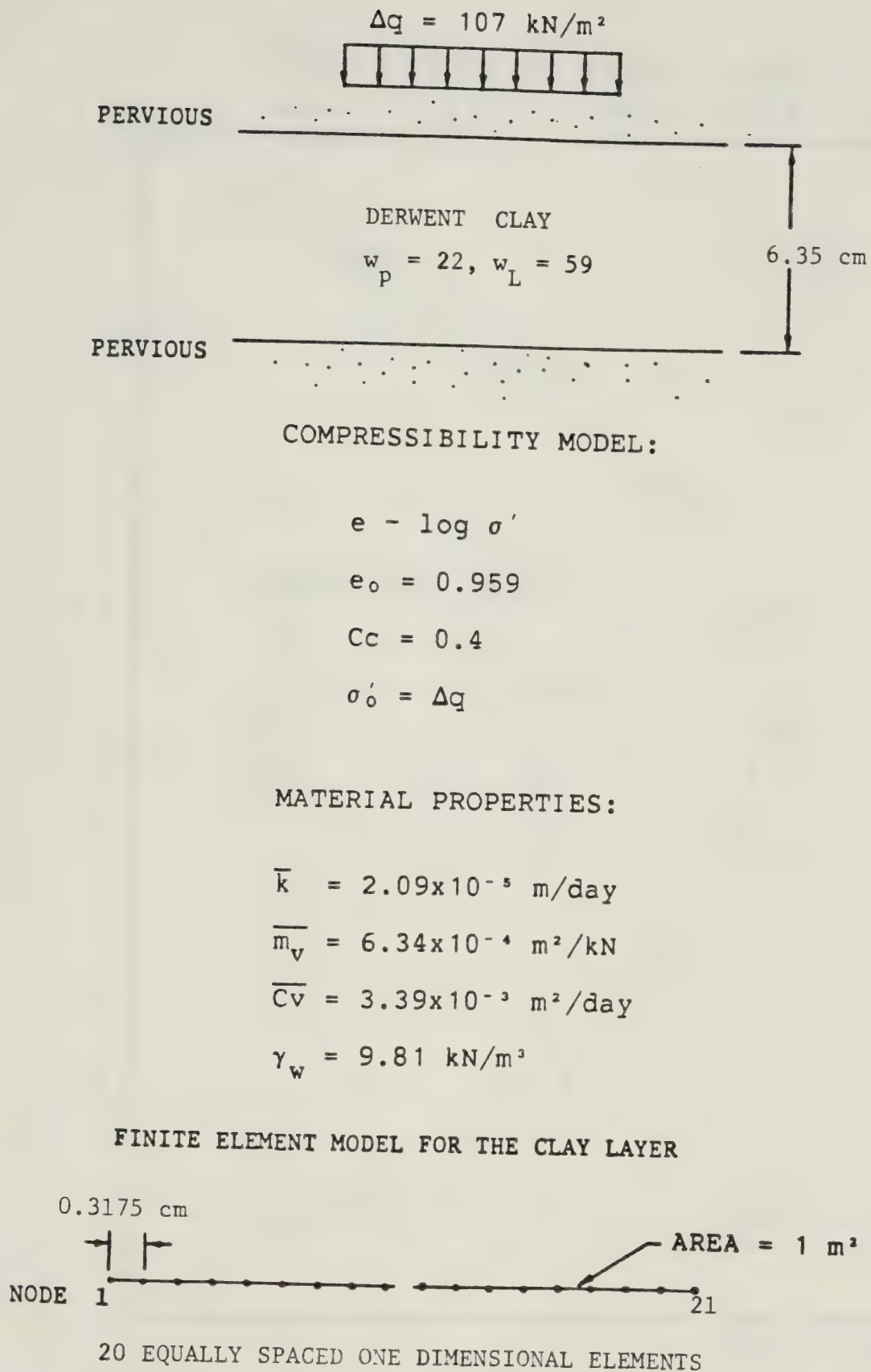


Figure 5.5 Nonlinear consolidation analysis with experimental results



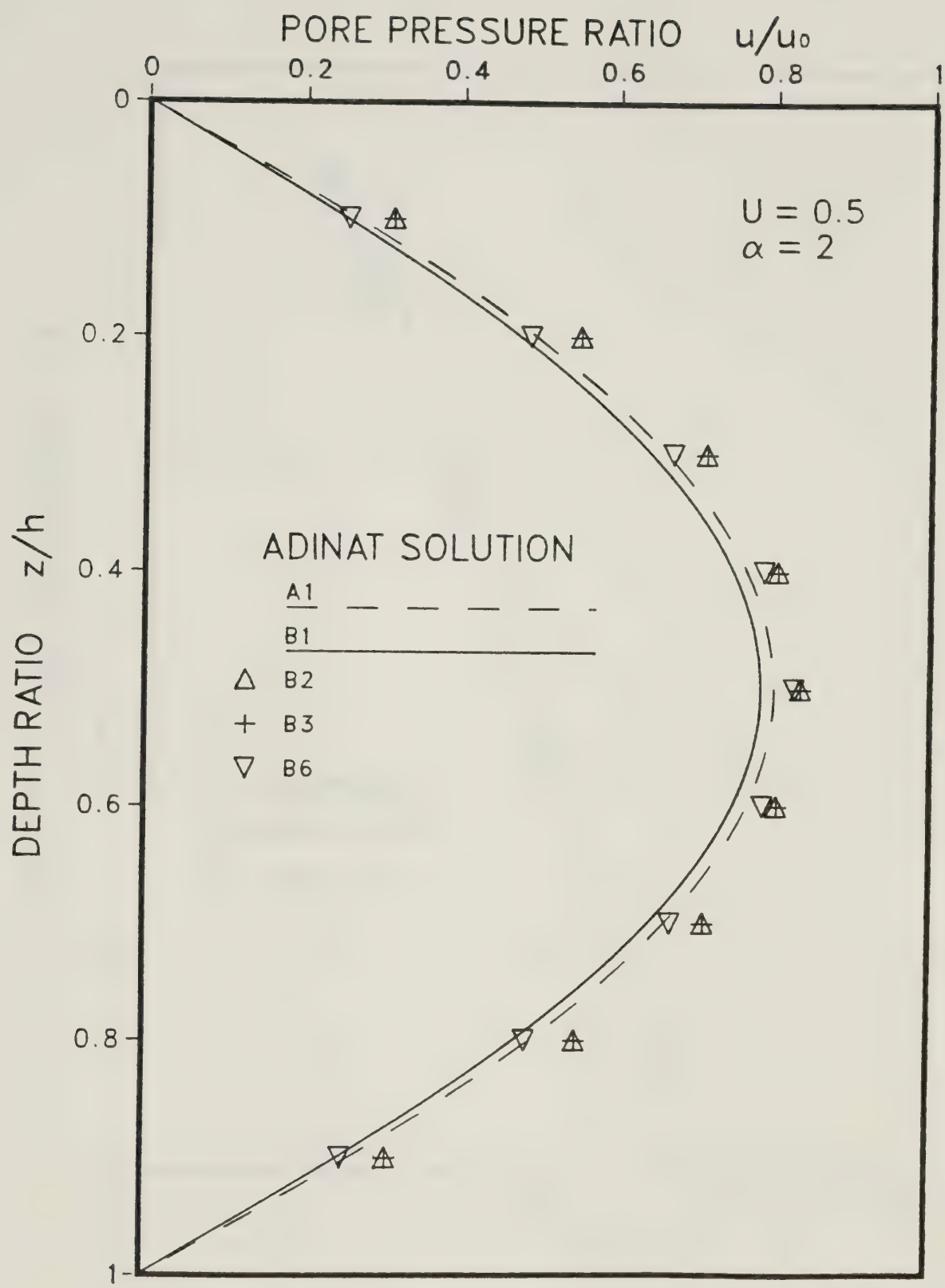


Figure 5.6 Pore pressure distribution in thin layer analysis





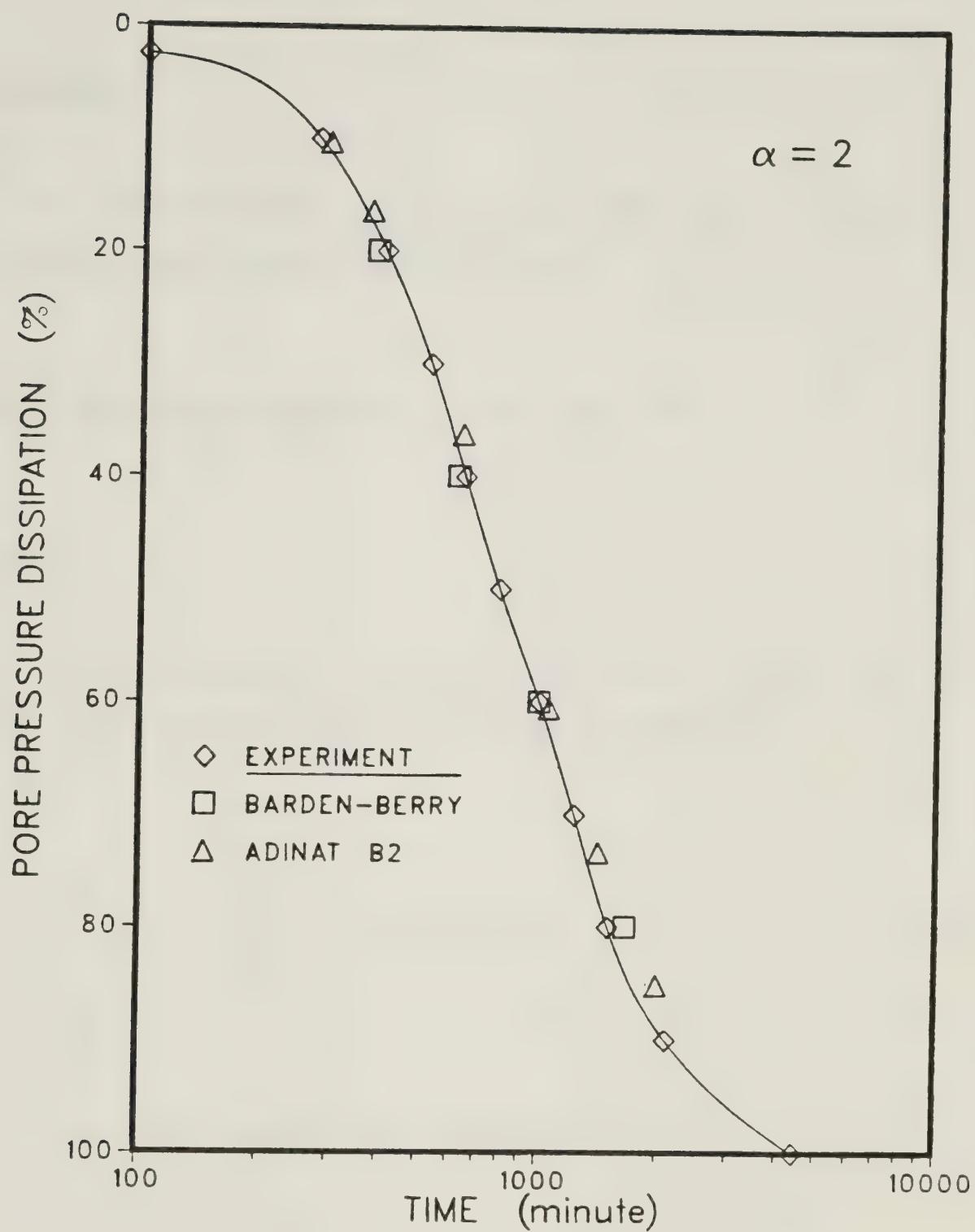


Figure 5.7 Pore pressure dissipation: small strain - thin layer analysis



## 5.4 Davis-Raymond Theory

### Objective:

To verify ADINAT small strain - thin layer formulation with Davis and Raymond theory (1965).

### ADINAT Nonlinear Equation: A1, A2, B2, B3

### Brief:

These authors have given a small strain theory for thin layer and the governing equation is expressed as

$$-\overline{Cv} \left[ \frac{1}{\sigma'} \frac{\partial^2 u}{\partial z^2} + \left( \frac{1}{\sigma'} \right)^2 \left( \frac{\partial u}{\partial z} \right)^2 \right] = \frac{1}{\sigma'} \frac{\partial \sigma'}{\partial t} \quad (5.4)$$

It can be rearranged into a different form

$$\overline{Cv} \frac{\partial^2 \omega}{\partial z^2} = \frac{\partial \omega}{\partial t} \quad (5.5)$$



with

$$\omega = \log_{10} \frac{\sigma'_f}{\sigma'_f} \quad (5.6)$$

The general solution can be determined through

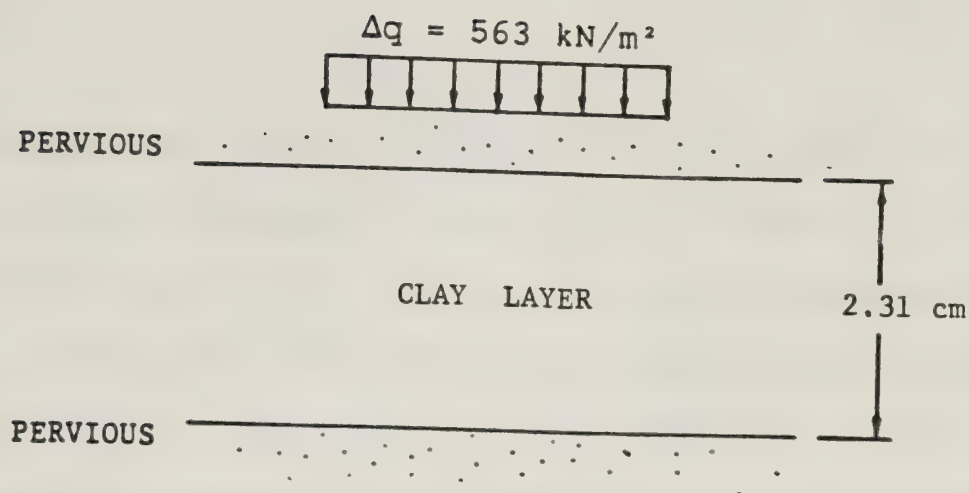
$$u = \sigma'_f \left[ 1 - \left( \sigma'_0 / \sigma'_f \right)^B \right] \quad (5.7)$$

B is Terzaghi's excess pore pressure ratio. Eq. 5.4 or 5.5 is a special case of equation A2, and Davis and Raymond have assumed the logarithmic compressibility law and that the coefficient of consolidation is constant.

### Description and Analysis:

The soil profile, loading and material properties are given in Fig. 5.8. Finite strain analyses are also carried out with equations B2 and B3. Eq. 21 in Davis and Raymond (Eq. 5.5 in above) is directly solved using ADINAT as well. Terzaghi's theory is also analysed by the program for comparisons. Datafiles are attached to Appendix A.13.





MATERIAL MODEL:  $e - \log \sigma'$   
 CONSTANT  $C_v$

MATERIAL PROPERTIES:  $e_o = 0.9$   
 $\sigma'_o = 188 \text{ kN/m}^2$   
 $\gamma_w = 9.81 \text{ kN/m}^3$   
 $C_c = 0.6$   
 $\overline{C_v} = 8.9 \times 10^{-4} \text{ m}^2/\text{day}$

#### FINITE ELEMENT MODEL FOR THE CLAY LAYER



20 EQUALLY SPACED ONE DIMENSIONAL ELEMENTS

Figure 5.8 One dimensional nonlinear analysis with constant coefficient of consolidation





Results:

ADINAT pore pressure results at  $T_v=0.4$  and with a load ratio of 4 are presented in Fig. 5.9. The small strain formulations agree well with the closed-form solutions using Eq. 5.7. Solutions with  $A_2$  and with Davis and Raymond's Eq. 21 are almost identical, as they should be. The finite strain solution shows a lower prediction while Terzaghi's theory gives an even lower result.



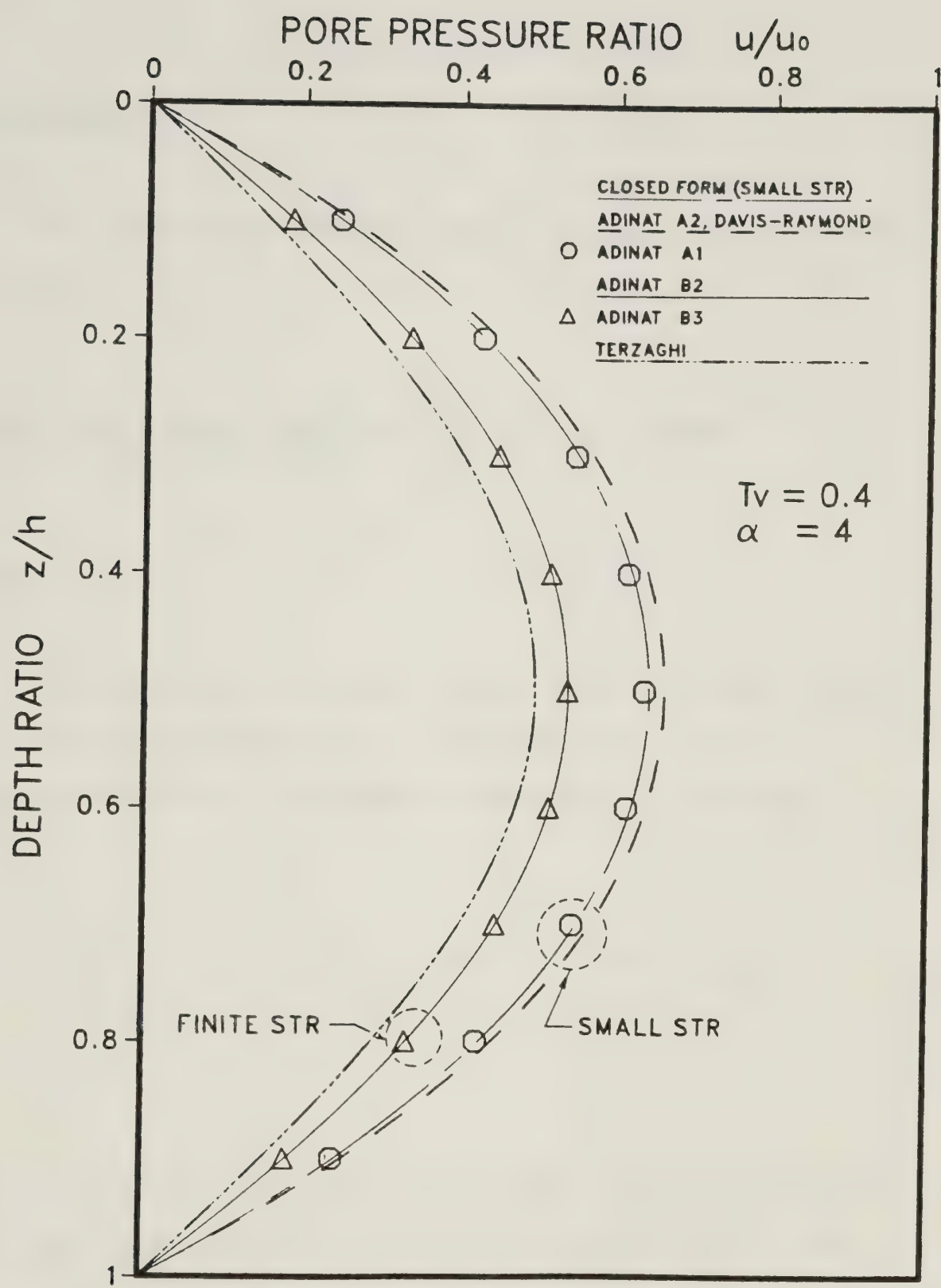


Figure 5.9 Pore pressure isochrones: thin layer analysis



## 5.5 Raymond Theory

### Objective:

To illustrate ADINAT application to Raymond's (1966) theory.

### ADINAT Nonlinear Equation: A1, A6, B7, B8, B9

### Brief:

Raymond has developed a small strain theory which includes both variations in permeability and in compressibility. His general equation is given as

$$\frac{k_n \sigma_n'}{\gamma w} \left( \frac{1}{\sigma'} \right)^{\nu} \left[ \frac{\nu}{\sigma'} \frac{\partial \sigma'}{\partial z} \frac{\partial u}{\partial z} - \frac{\partial^2 u}{\partial z^2} \right] = \frac{0.434 C_c}{\sigma' (1+e_0)} \frac{\partial \sigma'}{\partial t} \quad (5.8)$$

$k_n$ ,  $\sigma_n'$  are constants which may equal to the initial value at the soil surface, or obtained by assuming  $e = 0$ . This equation is the same as equation A1 with the assumed logarithmic material laws.



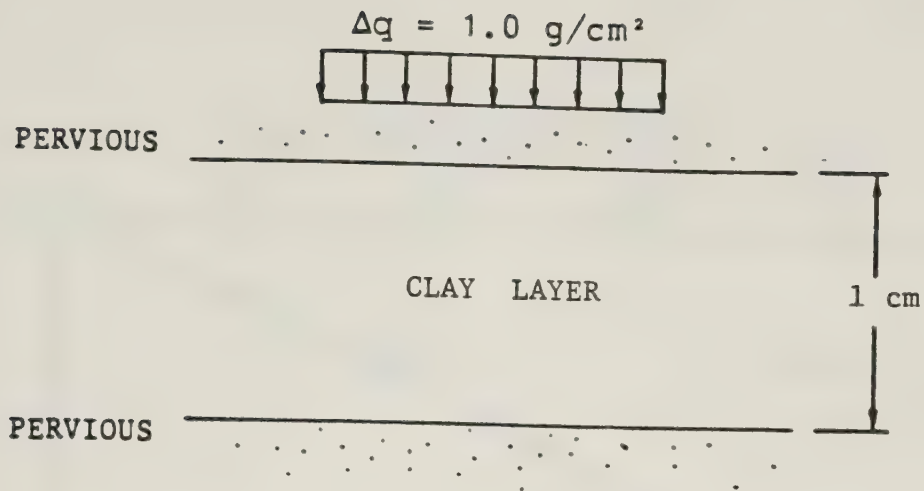
### Description and Results:

Fig. 5.10 shows the soil profile and material properties. The loading and finite element model is given in the figure. Input data is included in Appendix A.14.

Pore pressures obtained by ADINAT using the above cited theory are shown in Fig. 5.11. Finite strain analysis is also performed and compared in the figure. Again, there is little difference between these results. This same problem is also solved by Raymond (1969) with small strain approximation and good coresspondence in the solution is noted.







COMPRESSIBILITY MODEL:

$$e - \log \sigma'$$

$$e_0 = 0.9$$

$$\sigma'_0 = 4.0 \text{ g/cm}^2$$

$$Cc = 0.6$$

PERMEABILITY MODEL:

$$e - \log k$$

$$k_0 = 0.01 \text{ cm/sec}$$

$$Ck = 0.3$$

OTHER PARAMETER:  $\gamma_w = 1.0 \text{ g/cm}^3$

FINITE ELEMENT MODEL FOR THE CLAY LAYER

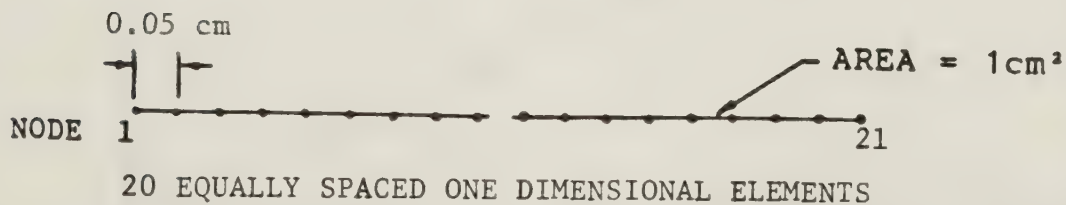


Figure 5.10 One dimensional consolidation with logarithmic material laws



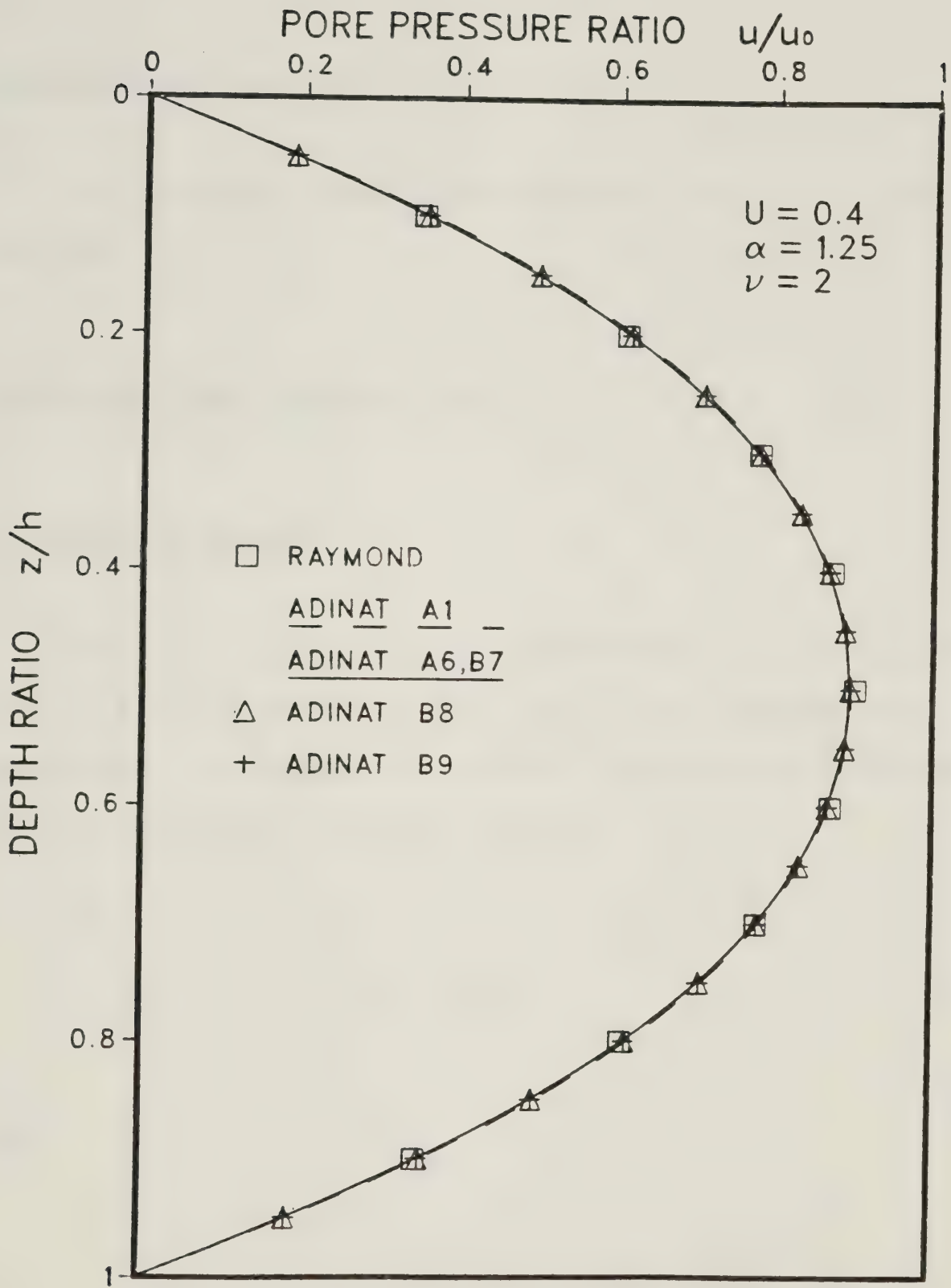


Figure 5.11 Pore pressure distribution in thin layer analysis



## 5.6 Poskitt Analysis

### Objective:

To illustrate ADINAT application to Poskitt's (1969) analysis.

### ADINAT Nonlinear Equation: B10

### Analysis and Results:

Poskitt has given the same formulation as equation B10 in Fig. 4.4. In order to make use of the given material parameters in his paper, changes of variables have been made to equation B10 as follows. Setting

$$c = (\pi a)/h \quad (5.9)$$

and

$$\tau = (4 \cdot Cvi \cdot t)/h^2 \quad (5.10)$$



and substituting these into equation B10, yields:

$$\frac{\partial}{\partial c} \left[ \frac{\pi^2 (\alpha \beta)^\mu}{4(1-\mu \xi_f)} \frac{\partial \mu}{\partial c} \right] = \frac{\partial \mu}{\partial \tau} \quad (5.11)$$

ADINAT analysis is carried out with the parameters  $\alpha$ ,  $\beta$ ,  $\xi_f$ ,  $C_{vi}/h^2$  of 2, 0.37, 0.05,  $1.535 \times 10^{-4}/\text{min}$  respectively. Input data is in Appendix A.15. The output pore pressure at  $\tau=0.5$  is shown in Fig. 5.12.





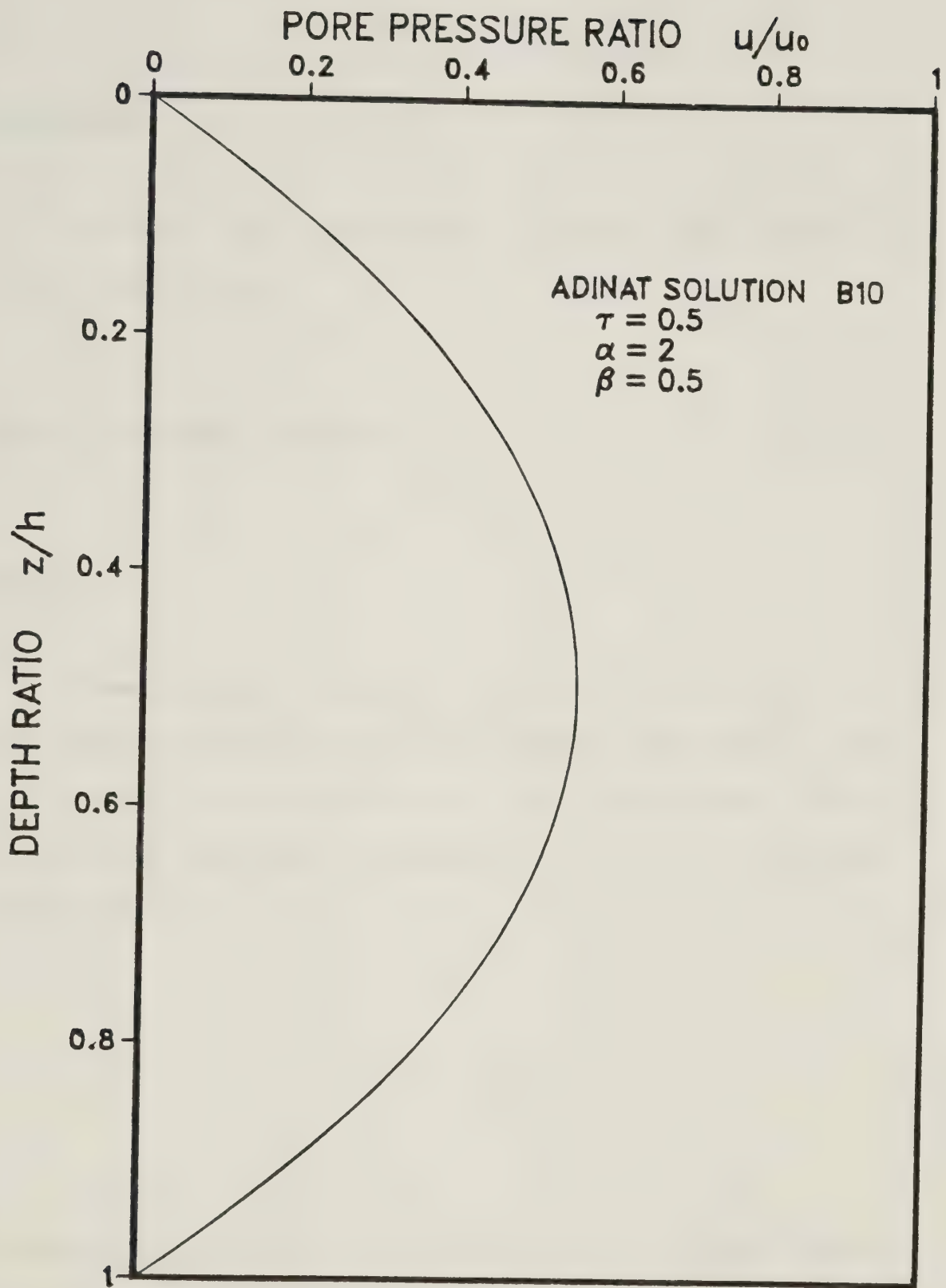


Figure 5.12 Pore pressure isochrones in finite strain - thin layer analysis



## 5.7 Gibson-England-Hussey General Formulation

### Objective:

To verify the application of ADINAT with Gibson *et al* (1967) formulation.

### ADINAT Nonlinear Equation: A7

### Brief:

These authors have given a general treatment of one dimensional consolidation of a single layer soil. They arrive at a governing equation for a thin layer which is essentially identical to equation A7, but the variable coefficient  $C_F$ ,

$$C_F = C_v \left( \frac{1+e_0}{1+e} \right)^2 \quad (5.12)$$

is assumed either as a constant, or as a linear function of the void ratio  $e$ ,



$$C_F = C_0 + \hat{m}(e - e_0) \quad (5.13)$$

$C_0$  and  $\hat{m}$  are constants which constitute the mode of variation of  $C_F$ .

### Description and Analysis

For comparing the results of Gibson *et al* and to use their given parameters, equation A7 is tranformed into

$$\frac{\partial}{\partial a'} \left[ C'_F \frac{\partial e'}{\partial a'} \right] = \frac{\partial e'}{\partial T'} \quad (5.14)$$

by changes of variables, where  $a'$  and  $e'$  are ratios of the orginal values.  $T'$  and  $C'_F$  are defined below,

$$T' = C_0 t / h \quad (5.15)$$

$$C'_F = (\lambda e' + 1) / (\lambda + 1) \quad (5.16)$$

in which



$$\lambda = \hat{m}e_o / (C_o - \hat{m}e_o) \quad (5.17)$$

The case of  $\lambda=0.4$  and  $e_f/e_o=0.2$  with the linear material model, Eq. 5.13, is illustrated here. The datafile is given in Appendix A.16.

### Results:

ADINAT predictions for the particular case taken from Gibson *et al* is presented in Fig. 5.13. Only half of the layer is shown because of symmetry. Gibson *et al*'s results are also reproduced in the figure and good correspondence is found.





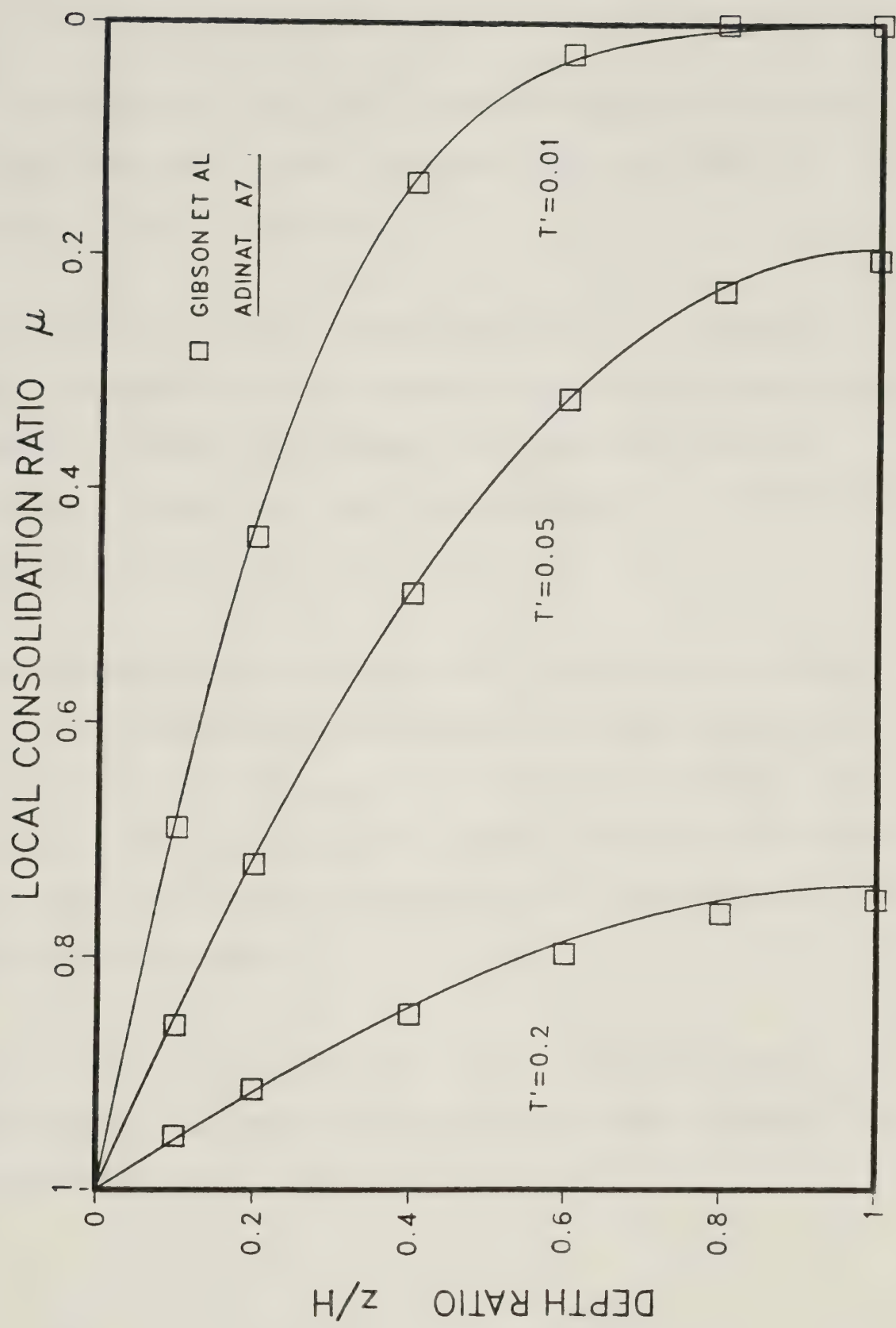


Figure 5.13 Variation of local consolidation ratio  $\mu$



## 5.8 Discussion

It can be seen from the preceding examples that most thin layer theories can be retrieved from the *ADINAT equations* in Fig. 4.3 and Fig. 4.4.

For cases where analytical solutions are obtained, excellent agreement can be observed with ADINAT predictions. For others in which the original work has been further approximated in some way, good correspondence is also shown between the theories and ADINAT results.

The closeness of the results of the finite strain and small strain analyses may be due to the insignificance of strain or void ratio changes in a thin layer problem. An exception to this occurs in the example of Davis and Raymond (1965), where the greater discrepancy is probably due to the larger load ratio used.

The difference between the finite strain and small strain analyses becomes more notable in deeper layers. This will be clearly seen in the next chapter.



## 6. NONLINEAR CONSOLIDATION ANALYSIS OF THICK SOIL LAYER BY ADINAT

### 6.1 General

In the following, examples are given for cases of thick layer analyses. Only a few papers on this subject can be found, probably because of the necessity of dealing with highly nonlinear equations. General descriptions, including the finite element model, given in Section 5.1, are also relevant here.

Notice that  $e_0$  is not a constant in thick layer analysis and the initial effective stress in the Lagrangian coordinate requires

$$\sigma'(a,0) = \sigma'_0 + (\gamma_s - \gamma_w) \cdot \int \frac{da}{1+e_0} \quad (6.1)$$

This expression appears much simpler in the reduced coordinate  $b$ ,

$$\sigma'(b,0) = \sigma'_0 + (\gamma_s - \gamma_w) \cdot b \quad (6.2)$$



The effective stress and excess pore pressure relationship at time  $t$  are given by Eq. 4.35 or Eq. 4.42, depending on the coordinate system used.

By integrating Eq. 4.40 and assuming the logarithmic effective stress - void ratio relationship, the following is obtained:

$$a = (1+e_0)b - \frac{\sigma'_0 Cc}{(\gamma_s - \gamma_w) \ln 10} (1 - \varsigma + \varsigma \ln \varsigma) \quad (6.3)$$

in which

$$\varsigma = 1 + \frac{\gamma_s - \gamma_w}{\sigma'_0} \cdot b \quad (6.4)$$

The above expression shows the relation between the coordinate  $a$  and  $b$ , and it is useful in the thick layer analysis with the logarithmic compressibility law. Notice that the Lagrangian coordinate  $a$  in Eq. 6.3 can be substituted by the Eulerian coordinate  $z$ , in which case the same change also applies to Eq. 4.40 for computing the reduce coordinate  $b$ .

Several other symbols are employed in thick layer analysis. The initial void ratio at surface,  $e_0$ , can be





determined by

$$e_{\emptyset} = \frac{\gamma_s - \gamma_t}{\gamma_t - \gamma_w} \quad (6.5)$$

The surface load ratio is computed from

$$\alpha_s = (\sigma'_{\emptyset} + \Delta q) / \sigma'_{\emptyset} \quad (6.6)$$

and the self-weight ratio  $\xi$  is defined as

$$\xi = [(\gamma_s - \gamma_w) \ell] / \sigma'_{\emptyset} \quad (6.7)$$

where  $\ell$  can be obtained from

$$\ell = \int_0^h \frac{d\theta}{1 + e(\theta, 0)} \quad (6.8)$$

$\theta$  represents either the coordinate  $a$  or  $z$ . A similar ratio using the average submerged weight, termed the average self-weight ratio, is expressed as



$$\bar{\xi} = ( \bar{\gamma}_b \cdot h ) / \sigma'_0 \quad (6.9)$$

in which  $\bar{\gamma}_b$  is the average submerged unit weight of the soil.

In applying thick layer formulations to ADINAT, the compressible stratum is divided into contiguous layers of finite depth. The expressions for material parameters are then evaluated at the mid-point of each layer. Although this is only an approximate method of the application of the thick layer formulation, good results are obtained as is shown in the following examples.



## 6.2 Small Strain Approximation

### 6.2.1 Raymond Analysis

#### Objective:

To illustrate the use of ADINAT in Raymond's (1969) analysis.

#### ADINAT Nonlinear Equation: A4

#### Brief:

Raymond has studied the consolidation of deep deposits in the context of small strain. Variations of permeability and compressibility are taken into account by employing logarithmic relationships. Notice that Raymond's dimensionless equation Eq. 11 used in his calculations is a further approximation of his general equation Eq. 6.

The formulation Raymond has given to solve deep deposit problems may be considered as a pseudo-thick layer analysis. The variable term  $(1+e_0)$  is incorporated in the time factor to which the output time for pore pressures results are also referred. This is of little practical use because the real time may not be easily determined.



### Description and Analysis:

An example is shown with the same variables assumed by Raymond. Setting

$$T = \frac{k'_n \sigma'_n \ln 10 (1+e_o)}{\gamma \frac{C_c}{w} h^2} t \quad (6.10)$$

and with the logarithmic material laws, gives from equation A4,

$$\frac{\partial}{\partial z'} \left[ \sigma'^{-\nu} \frac{\partial u}{\partial z'} \right] = \frac{\sigma'_\theta^{(1-\nu)}}{\sigma'} \frac{\partial u}{\partial T} \quad (6.11)$$

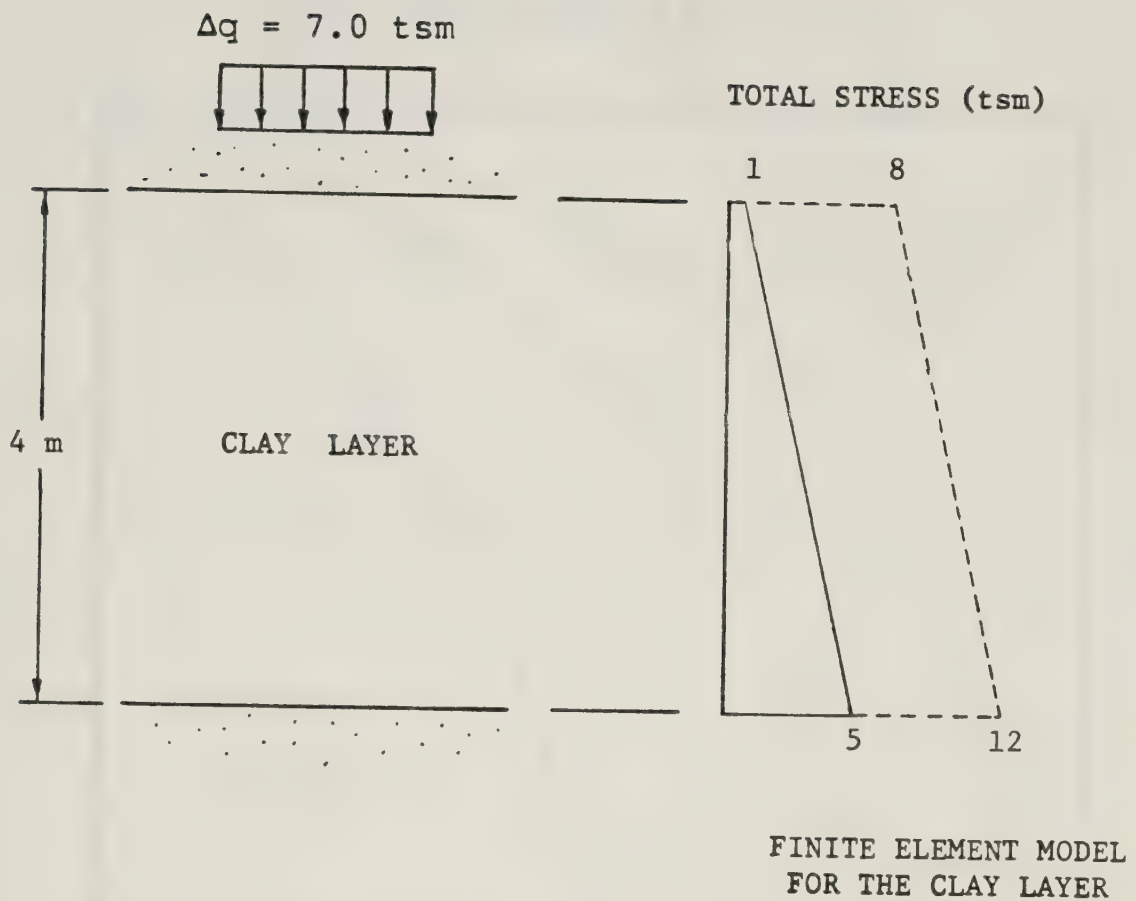
where  $z'$  is the depth ratio  $z/h$ . The soil profile and material parameters for ADINAT analysis are given in Fig. 6.1. The datafile is contained in Appendix A.17.

### Results:

Pore pressure distribution at 40%, 60% and 80% consolidation computed by ADINAT are compared with Raymond's solutions in Fig. 6.2. Slightly larger differences can be seen at later consolidation stages.







MATERIAL MODEL:

$$e - \log \sigma'$$

$$e - \log k$$

MATERIAL PROPERTIES:

$$\nu = 2$$

$$\frac{\gamma}{b} = 1.0 \text{ t/m}^3$$

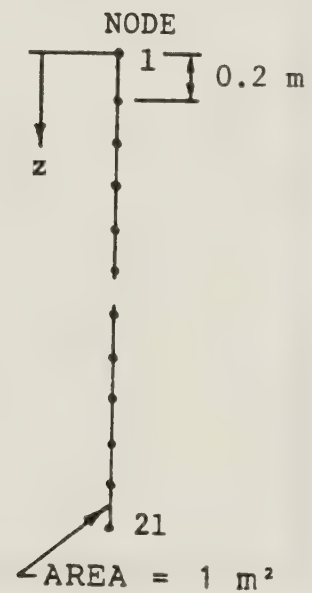


Figure 6.1 Nonlinear consolidation analysis in Raymond (1969) example



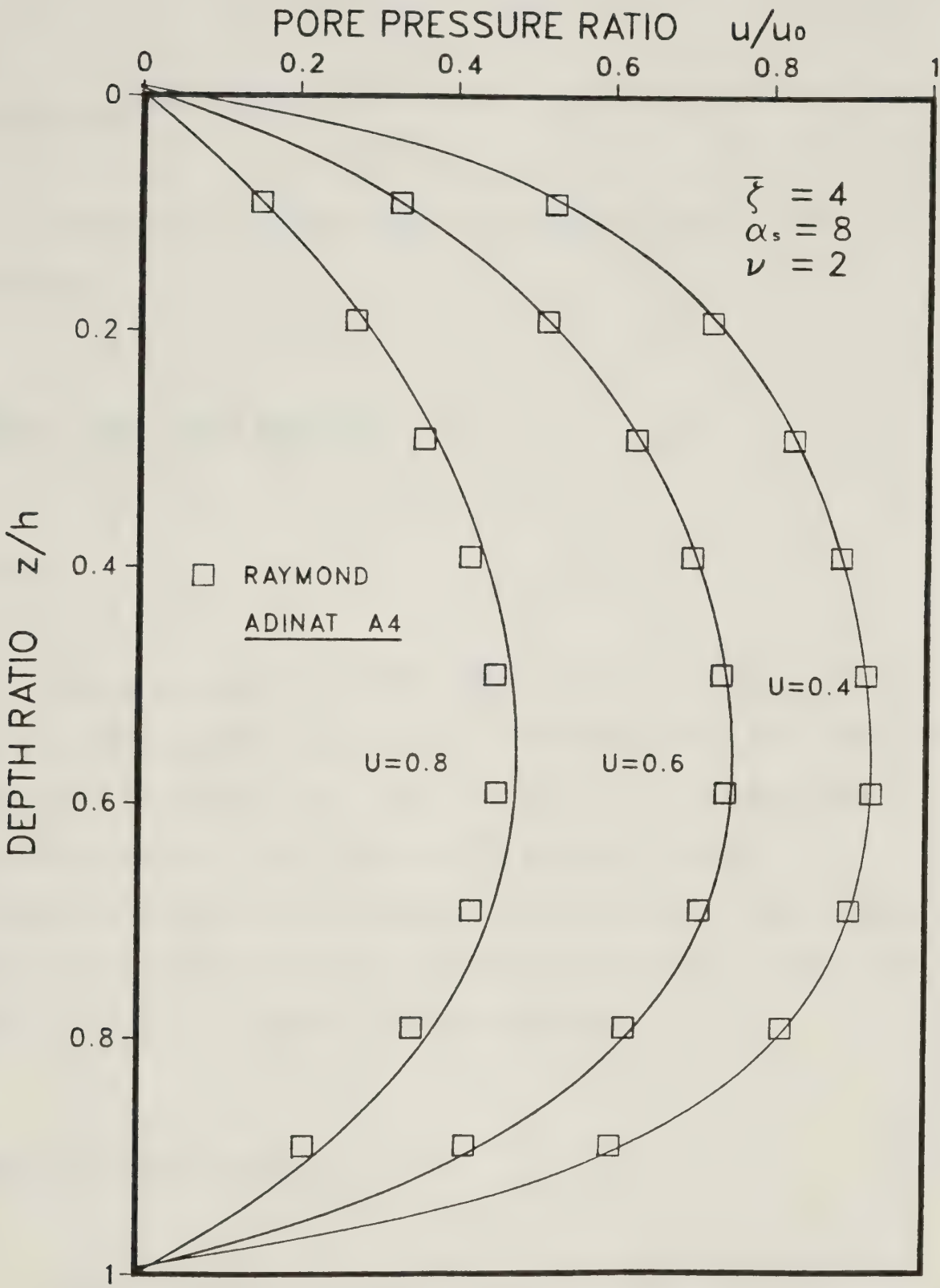


Figure 6.2 Pore pressure distribution in small strain - thick layer analysis



### 6.2.2 Davis Theory

#### Objective:

To demonstrate the use of ADINAT in Davis (1971) analysis.

#### ADINAT Nonlinear Equation: A4

#### Brief:

Davis has given a small strain theory for a thick layer. His governing equation is essentially the same as equation A4 except that the coefficient of consolidation is assumed constant and the coefficient of volume compressibility is inversely proportional to the effective stress. Closed-form solutions are provided in some limiting cases along with these simplifications.

#### Analysis and Results:

Substituting the above mentioned approximations to equation A4, which becomes

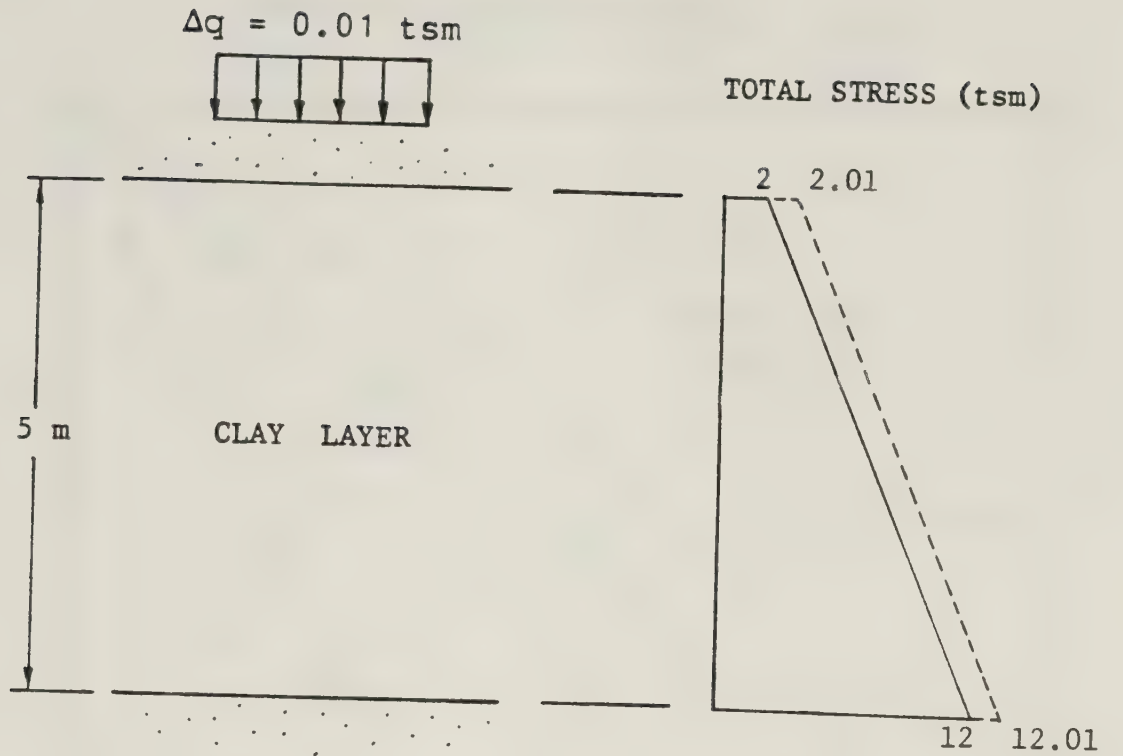


$$\frac{\partial}{\partial z} \left[ \frac{\overline{Cv}}{\sigma'} \frac{\partial u}{\partial z} \right] = \frac{1}{\sigma'} \frac{\partial u}{\partial t} \quad (6.12)$$

This equation is then used for the analysis. An example of the load ratio,  $\alpha$ , of 1.005 and the average self-weight ratio,  $\bar{\xi}$ , of 5 is given in Fig. 6.3. The input datafile is in Appendix A.18. Pore pressure results at time factors 0.2 and 0.3 given by ADINAT are plotted in Fig. 6.4. Davis' closed-form solution for  $\alpha_s = 1$  and Terzaghi's solutions are also compared in the figure.







FINITE ELEMENT MODEL  
FOR THE CLAY LAYER

MATERIAL MODEL:

$e - \log \sigma'$   
CONSTANT  $C_v$

MATERIAL PROPERTIES:

$\overline{C_v} = 1.0 \text{ m}^2/\text{day}$   
 $\overline{\gamma}_b = 2.0 \text{ t/m}^3$

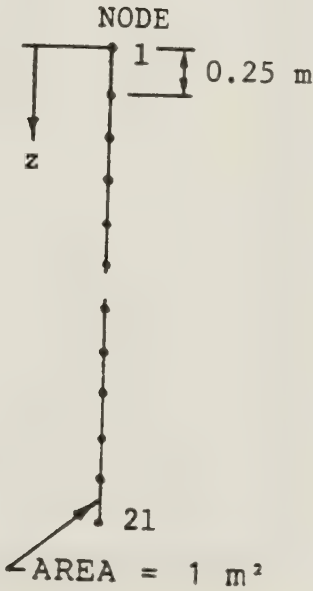


Figure 6.3 Nonlinear consolidation analysis with Davis (1971) formulation



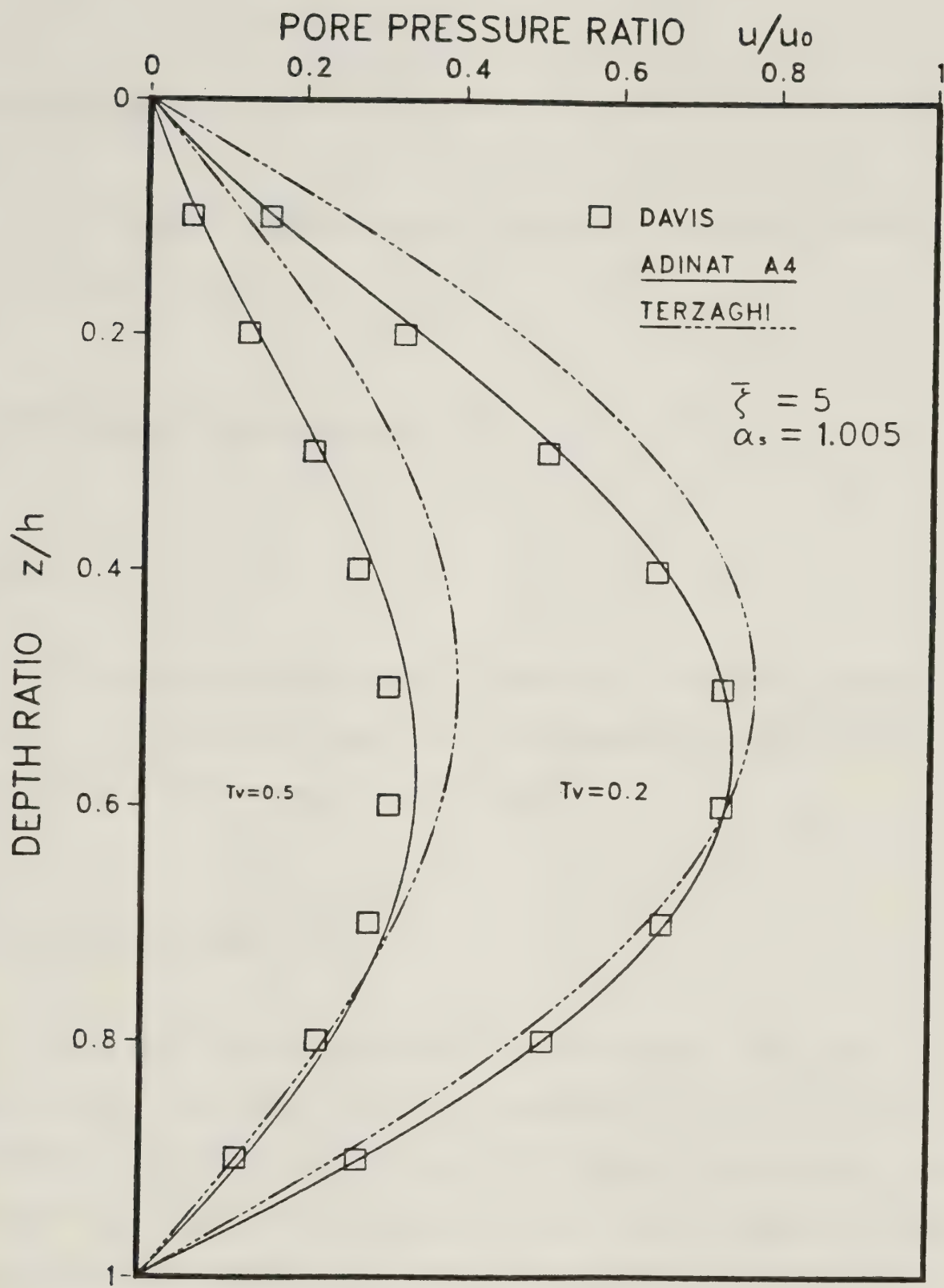


Figure 6.4 Pore pressure variation in thick layer analysis



### 6.2.3 Viggiani Numerical Solution

#### Objective:

To compare ADINAT predictions with Viggiani (1973) numerical solutions.

#### ADINAT Nonlinear Equation: A4

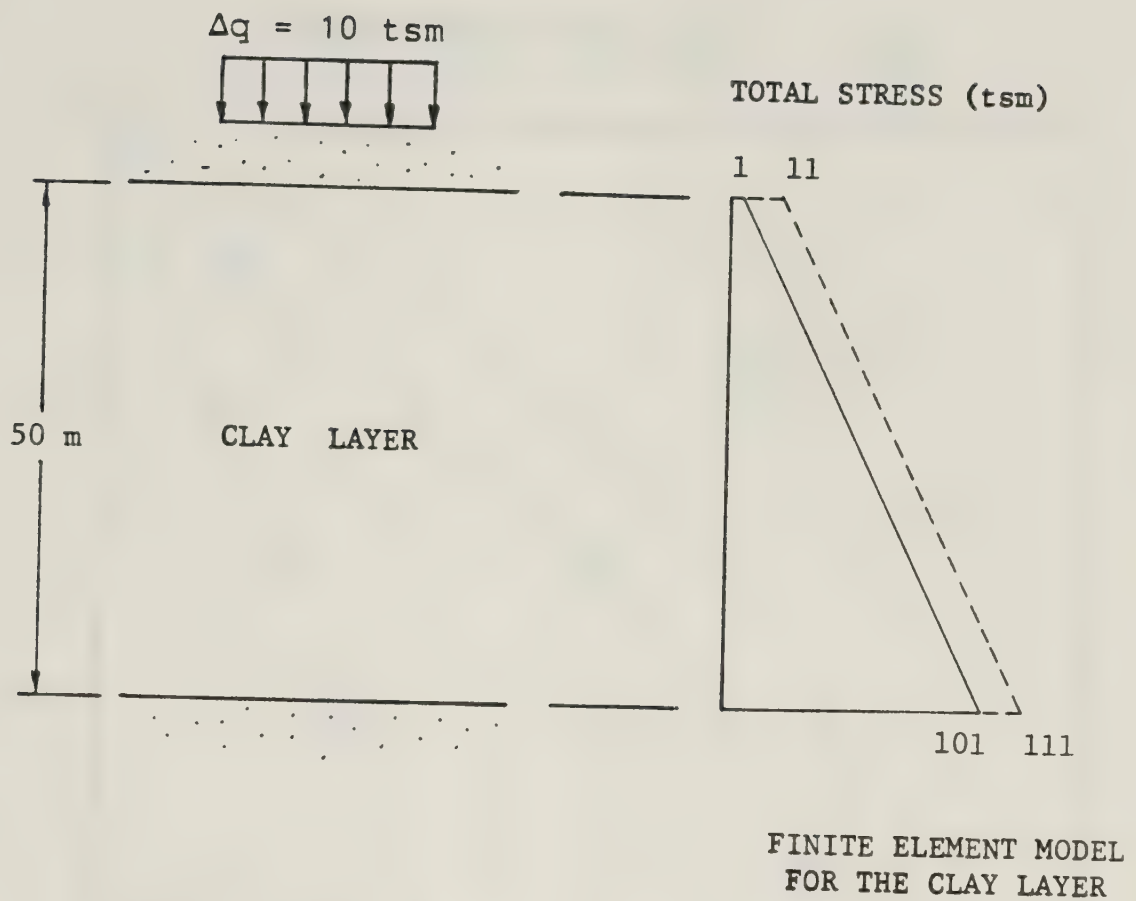
#### Brief:

This author has given a numerical treatment to Davis' (1971) formulation, and also considered the effects of variable load increment with depth.

#### Analysis and Results:

The soil layer and material parameters are shown in Fig. 6.5. The datafile is in Appendix A.19. Eq. 6.12 is used for ADINAT analysis and pore pressure results at time factor 0.2 and 0.3 are given in Fig. 6.6. Viggiani's and Terzaghi's solutions are included in the figure for comparison.





MATERIAL MODEL:

$$e - \log \sigma'$$

CONSTANT  $C_v$

MATERIAL PROPERTIES:

$$\overline{C_v} = 1.0 \text{ m}^2/\text{day}$$

$$\overline{\gamma}_b = 2.0 \text{ t/m}^3$$

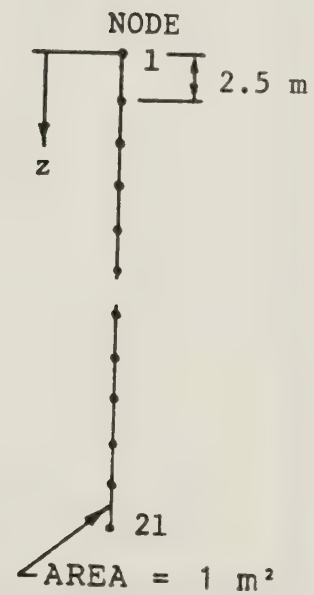


Figure 6.5 Example of nonlinear consolidation with Viggiani (1973) data





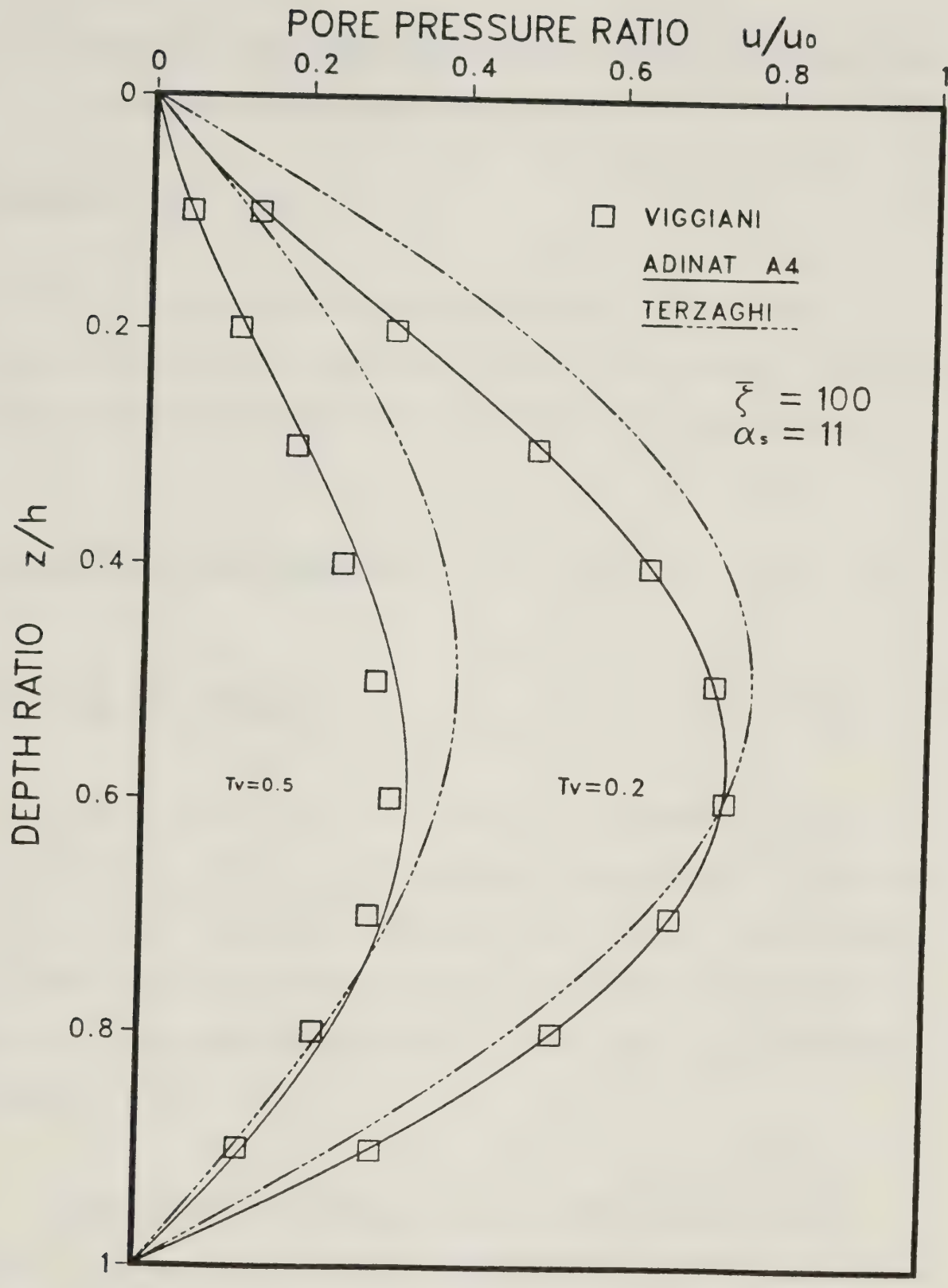


Figure 6.6 Pore pressure variation in thick layer analysis



## 6.3 General Analysis

### 6.3.1 Example with Logarithmic Material Law and Constant Coefficient of Consolidation

#### Objective:

To illustrate ADINAT application to nonlinear consolidation with logarithmic effective stress - void ratio relationship and constant coefficient of consolidation.

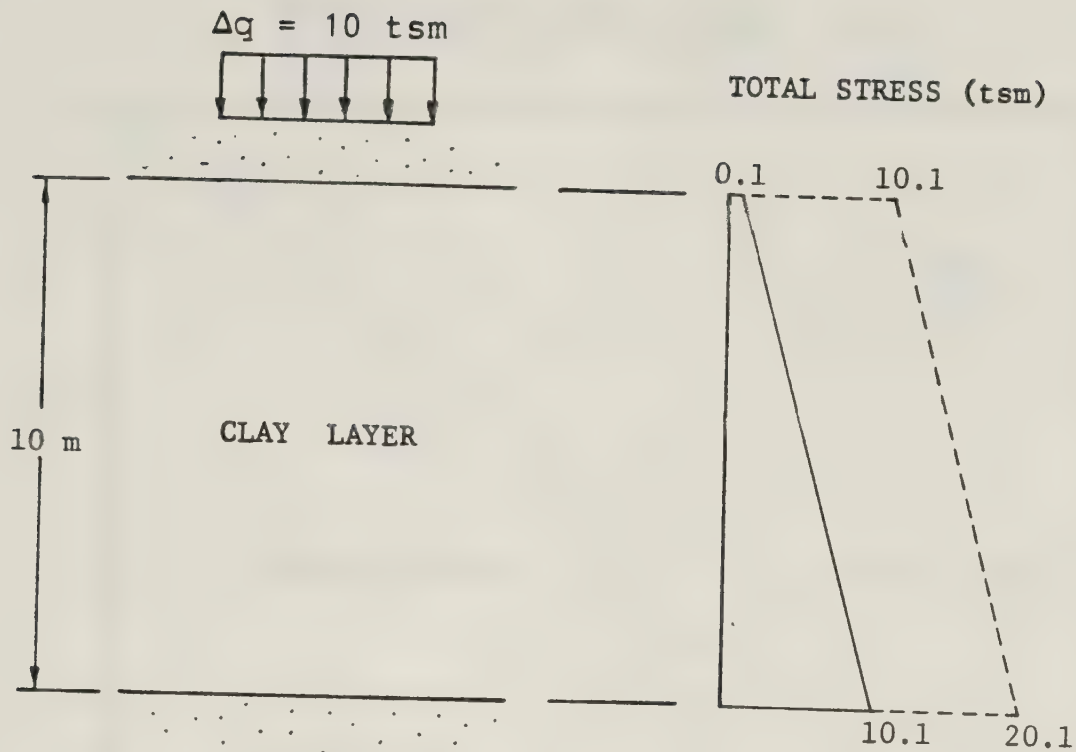
#### ADINAT Nonlinear Equation: A4, A8, A9

#### Analysis and Results:

The problem layout and material parameters are presented in Fig. 6.7. The small strain equation A4 and both of the two finite strain - thick layer formulations, A8 and A9, each of which utilizes a different coordinate system, are used for analyses. Input data can be found in Appendix A.20.

Pore pressure distributions predicted by ADINAT are shown in Fig. 6.8. Greater difference between the small strain and finite strain formulation can now be seen. Terzaghi's solution is also included in the figure.





MATERIAL MODEL:

$$e - \log \sigma'$$

CONSTANT  $C_v$

FINITE ELEMENT MODEL  
FOR THE CLAY LAYER

MATERIAL PARAMETERS:

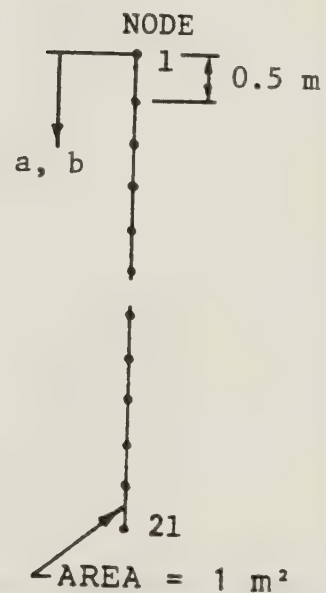
$$e_0 = 0.9; \quad C_c = 0.4$$

$$\overline{C_v} = 1.0 \text{ m}^2/\text{day}$$

$$\gamma_t = 2.0 \text{ t/m}^3$$

$$\gamma_w = 1.0 \text{ t/m}^3$$

$$\gamma_s = 2.9 \text{ t/m}^3$$



REDUCED CLAY LAYER THICKNESS  $\ell = 8.438 \text{ m}$

Figure 6.7 Consolidation analysis with logarithmic material law and constant coefficient of consolidation



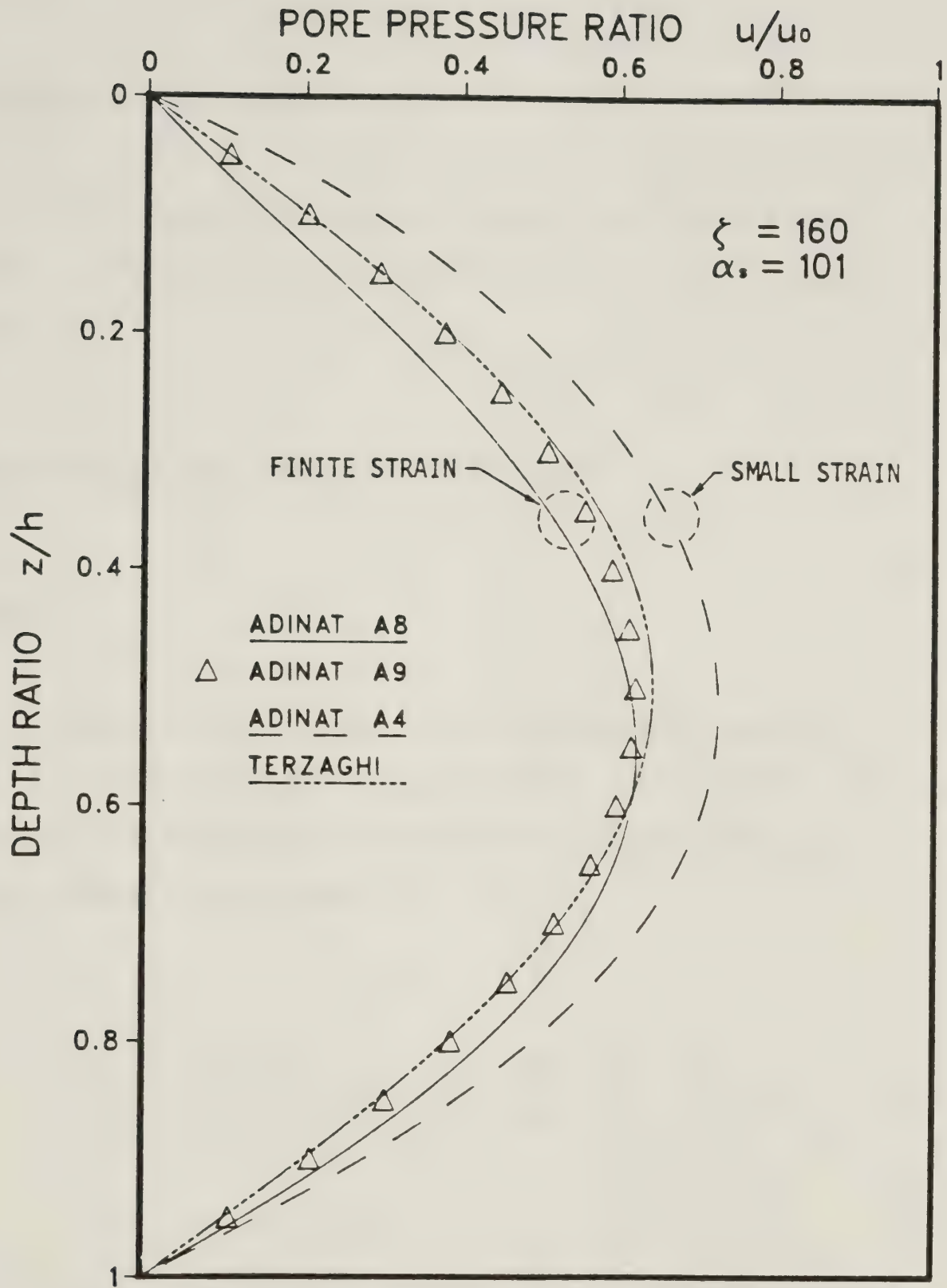


Figure 6.8 Pore pressure distribution in thick layer analysis





### 6.3.2 Gibson-Schiffman-Cargill General Analysis

#### Objective:

To illustrate the use of ADINAT in finite strain - thick layer analysis with Gibson *et al*'s (1981) general theory.

#### ADINAT Nonlinear Equation: A4, A8, A9

#### Brief:

These authors applied the same general equation, Eq. 13, from an earlier paper, Gibson *et al* (1967), but the effect of self-weight is included in their 1981 paper. Their governing equation

$$\left(1 - \frac{\gamma_s}{\gamma_w}\right) \frac{d}{de} \left[ \frac{k}{1+e} \right] \frac{\partial e}{\partial b} + \frac{\partial}{\partial b} \left[ \frac{k}{\gamma_w (1+e)} \frac{d\sigma'}{de} \right] \frac{\partial e}{\partial b} + \frac{\partial e}{\partial t} = 0 \quad (6.13)$$

is basically the same as equation A9. They made further assumptions and arrive at an approximate equation as



$$\frac{\partial^2 e}{\partial b^2} - \lambda' (\gamma_s - \gamma_w) = \frac{1}{g} \frac{\partial e}{\partial t} \quad (6.14)$$

where  $\lambda'$  and  $g$  are assumed constants, which may be computed as:

$$\lambda' = - \frac{1}{\sigma'(a)} \ln \left[ \frac{e(a) - e(\infty)}{e_0 - e(\infty)} \right] \quad (6.15)$$

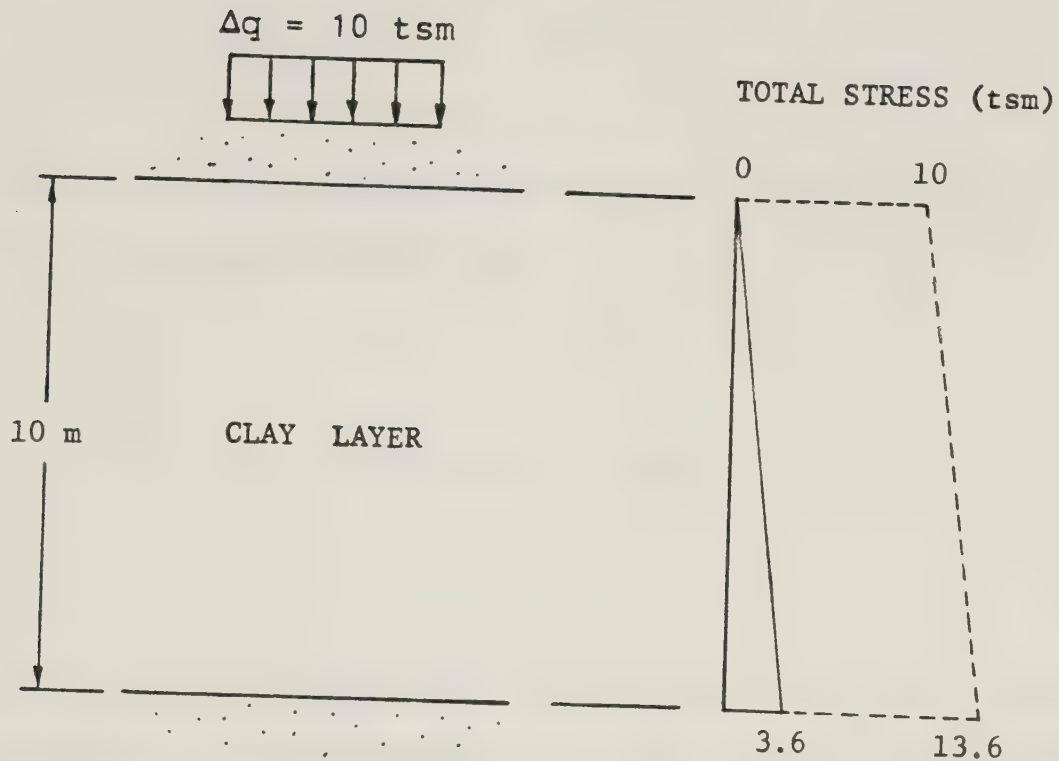
$$g = \overline{C_v} / [1 + e(a)] \quad (6.16)$$

The bracket (a) denotes the average value and ( $\infty$ ) refers to the final value at the layer bottom.

### Description and Analysis:

The soil layout and material parameters from Gibson *et al*'s (1981) example are given in Fig. 6.9. The compressibility law employed here is the same as that in the paper:





MATERIAL MODEL:

$$\sigma' - \ln e$$

CONSTANT  $C_v$

FINITE ELEMENT MODEL  
FOR THE CLAY LAYER

MATERIAL PARAMETERS:

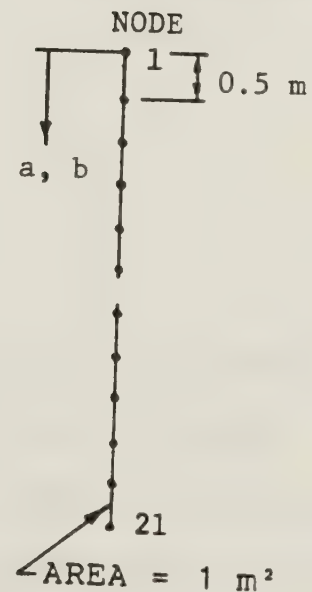
$$e_0 = 3.83; \quad e_\infty = 1.49$$

$$\lambda' = .332 \text{ m}^2/\text{t}$$

$$\bar{C}_v = .001944 \text{ m}^2/\text{day}$$

$$\gamma_t = 1.36 \text{ t/m}^3; \quad \gamma_w = 1.0 \text{ t/m}^3$$

$$\gamma_s = 2.74 \text{ t/m}^3$$



REDUCED CLAY LAYER THICKNESS  $\ell = 2.728 \text{ m}$

Figure 6.9 One dimensional nonlinear consolidation with Gibson et al (1981) analysis



$$\sigma'(b,t) = - \frac{1}{\lambda'} \ln \left[ \frac{e - e(\infty)}{e_0 - e(\infty)} \right] \quad (6.17)$$

The coefficient of volume compressibility can be determined with this relation and Eq. 3.8,

$$m_v = \lambda' [e - e(\infty)] / (1 + e) \quad (6.18)$$

Note that  $C_v$  is treated as a constant in the paper and hence the coefficient of permeability  $k$  can be calculated from Eq. 3.7.

Twenty solution time steps are used in this example. Datafiles are given in Appendix A.21.

### Results:

Pore pressures predicted by ADINAT at two different times are shown in Fig. 6.10 and Fig. 6.11 with Gibson *et al*'s (1981) solutions. Reasonable agreement can be seen. A small strain analysis of the same problem is given in the figure. Terzaghi's solutions is also included for comparisons.





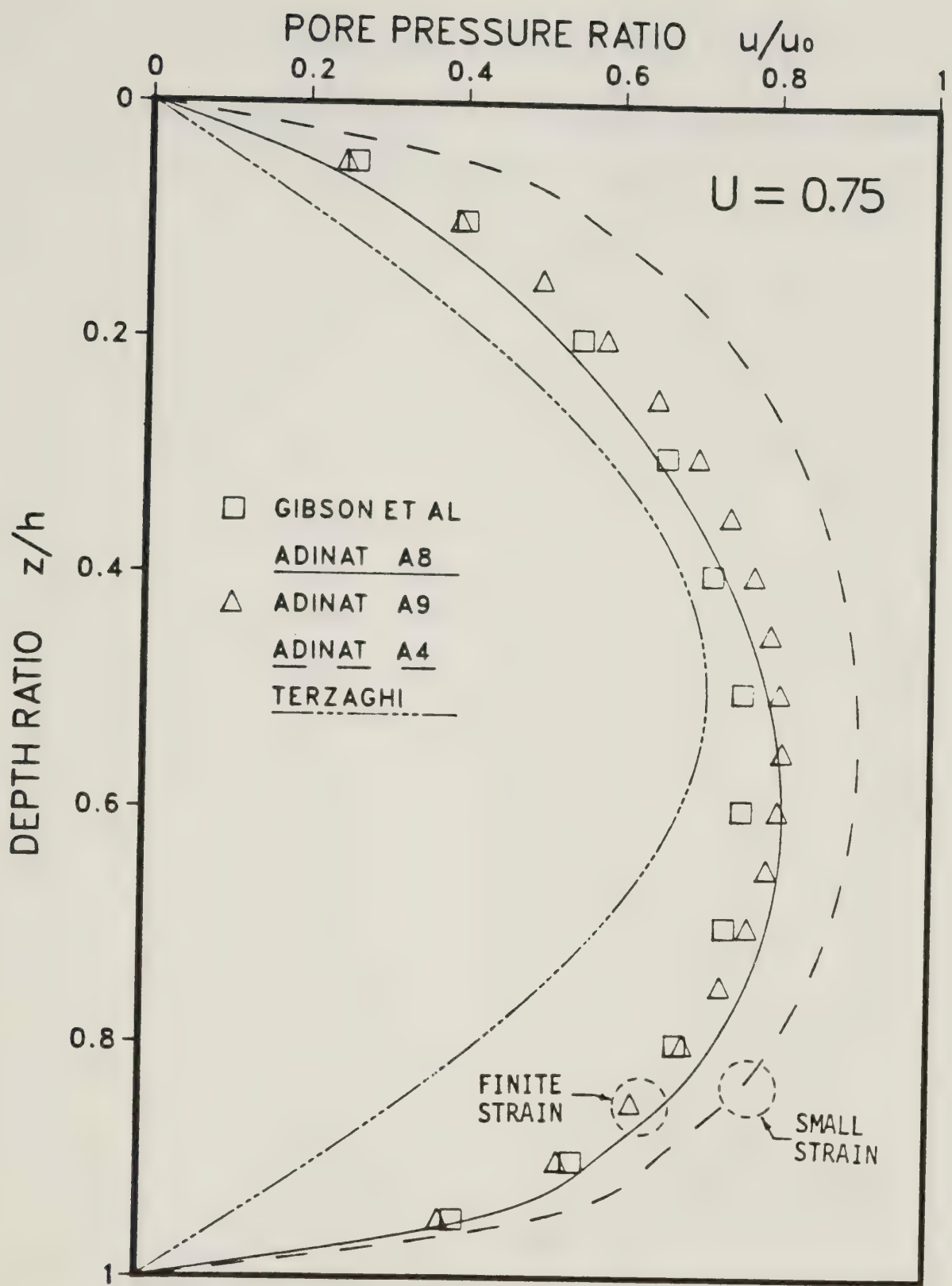


Figure 6.10 Pore pressure isochrones in thick layer analysis



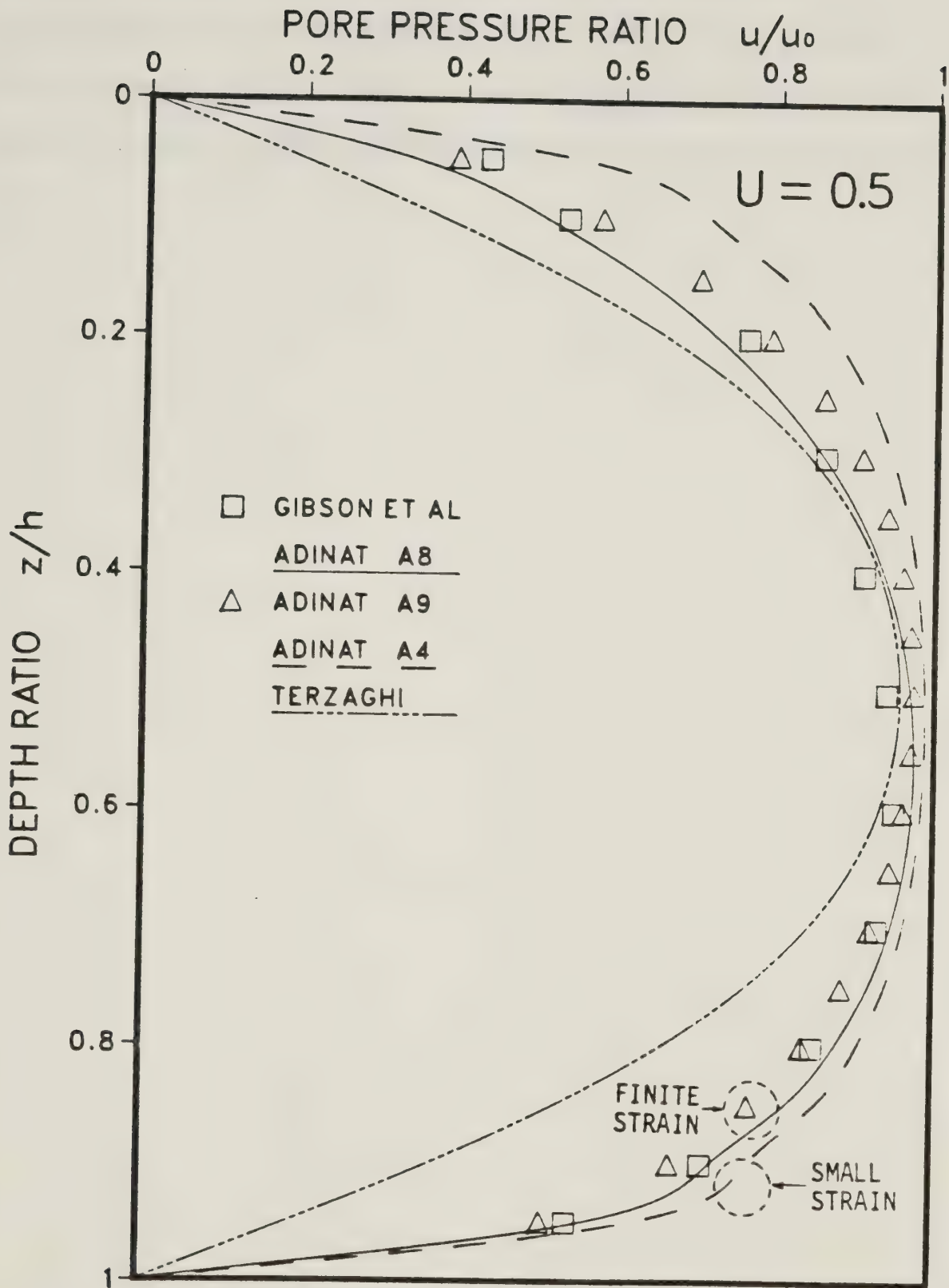


Figure 6.11 Pore pressure isochrones in thick layer analysis



It can be seen from the figures that the conventional theory is an underestimate while the small strain approximation is an overestimate of the pore pressure. Little 'bumps' on the top and bottom parts of the smooth curves are due to the limitations of the plotting machine.



## 6.4 Discussion

Since thick layer analysis allows for variation of effective stress with depth, and large changes of void ratio or strain, the difference between small strain and finite strain analysis becomes more significant as shown in the preceding examples. In addition, pore pressure isoschrones are all skewed because of the unsymmetrical stress distribution in the soil layer.

Note that the *average* submerged weight is used in the analyses by Raymond (1969), Davis (1971) and Viggiani (1973). Since this value does not consider the variation of void ratio throughout the soil depth, it is apparently inadequate to take account of the self-weight effects. The theories derived by the cited authors, therefore, are only approximate to the thick layer analyses.

In the analyses of thick deposits using ADINAT, preparation of input data may become tedious since material parameters change with depth. Simple supplemental programs can be easily developed for this purpose. Examples are given in Appendix B.





## 7. CONCLUSIONS

Application of ADINAT to various one dimensional consolidation theories are illustrated. Within the scope of the present work, finite element solutions are readily found with simple procedures for linear and nonlinear analyses. The overall simplicity in the use of the program makes it a convenient tool for general consolidation studies. Although examples are selected from a large area of ADINAT applicability, only part of the many possible applications of the program have been shown.

The high accuracy of ADINAT and its incremental solution procedure are particularly useful in soil problems with nonlinear material behaviour. The *ADINAT equations* in Fig. 4.3 and Fig. 4.4 provide a simple means for extensive research into nonlinear consolidation behaviour.

Use of the program for nonlinear analysis is straightforward once the material relationships for a given problem is determined. These relationships are then input into the program at discrete points in the range of consolidation. Linear interpolation is carried out by ADINAT to obtain values between the input points.



If the effective stress is a function of void ratio only, there should be no great difficulty in finding their functional relationship experimentally. A series of standard consolidation tests may well be performed to accomplish this. The problem of determining permeability and void ratio relationship appears to be a little more difficult. This may possibly be resolved by additional measurements attached to the consolidation equipment.

Parametric studies can be undertaken easily by the program. Effects of load ratio, self-weight ratio, soil depth and other physical conditions can be investigated. The assistance of a small supplemental program for the preparation of input may greatly simplify the use of ADINAT for this purpose.

New sets of useful consolidation parameters for nonlinear theories may also be developed by reviewing the proposed formulations. Comprehensive experimental and case studies, however, should be performed to demonstrate the usefulness of these parameters.

The operation of ADINAT can be stopped and restarted at preselected time steps to allow for changes of input options. This restart capability may be used for analysis of problems with stress history and time dependent loading.



The program also permits the elements to be active or inactive at any specific solution time. This option may be used to simulate the physical situations for a construction period or moving boundaries.

The orthotropic material model of ADINAT can be used to analyse anisotropic conditions in soils. Variable load increment can be handled easily by the program by specifying the appropriate loading at nodal points throughout the layer.

Impeded drainage boundary conditions may be treated as layered problems with the additional top and bottom layers being assigned appropriate permeability values. Radial drainage conditions can be modeled using axisymmetric elements.

Assessment of the numerical behaviour of ADINAT, such as stability in nonlinear analyses, is required in future work. A better understanding of the numerical behaviour is essential for further development of the program. This will also facilitate program use and reduce operating costs.





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## APPENDIX A

### Listings of ADINAT Input Datafiles



The ADINAT program (a Finite Element Program for Automatic Dynamic Incremental Nonlinear Analysis of Temperatures) is available on the MTS system at the University of Alberta. A copy of the user's manual can be obtained from the Department of Civil Engineering. The following commands can be used to run the program

```
$RUN 3012:ADINAT 5=input 6=output
```

Listings of datafiles corresponding to the examples discussed in the thesis can be found in this appendix. The first card in each datafile contains a brief note of the file. Computer programs used for the preparation of input data of several examples are given in Appendix B.



## LIST OF 'ADINAT' INPUT FILES

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#### A.4 Four layers

CARD NUMBER	1	2	3	4	5	6	7	8
1	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890
2	1-D CONSOLIDATION (4 LAYERS, DOUBLE DRAIN) DATA Schiffman&Stein70							
3	41	4	0	1	1	0	20	0
4	2	0	0					
5								
6								
7								
8								
9	20							
10	37							
11	1	1	0	0	0	0		
12	2	0	0	0	2	1		
13	40	0	0	0	78	0		
14	41	1	0	0	80			
15	1	1						
16	1	1	1	1	1	1	1	1
17	1	1	1	1	1	1	1	1
18	1	1	1	1	1	1	1	1
19	1	1	1	1	1	1	1	1
20	1	1	1	1	1	1	1	1
21	1	1	1	1	1	1	1	1
22	1	1	1	1	1	1	1	1
23	1	2						
24	1	2						
25			0	7500	0			
26	1	5					1	1
27	1	1						
28	00000789							
29	00018197							
30	1		2	1	1			
31	5	5	6	1				
32	1	10					1	1
33	1	1						
34	0000234							
35	0001220							
36	1	6	7	1	1			
37	10	15	16	1				
38	1	15					1	1
39	1	1						
40	00000333							
41	00006077							
42	1	16	17	1	1			
43	15	30	31	1				
44	1	10					1	1
45	1	1						
46	00000835							
47	00012172							
48	1	31	32	1	1			
49	10	40	41	1				
50	STDP							





## A.5 Linear variation of permeability

CARD NUMBER	C O L U M N   N U M B E R							
	1	2	3	4	5	6	7	8
1	1-D CONSOLIDATION (LIN K, D DR) DATA Martins65							
2	41	1	0	1	1	0	10	0
3	2	0						
4								
5								
6								
7								
8								
9	10							
10	319.47788							
11	1	1	0	0	0	0		
12	2	0	0	0	05	1		
13	40	0	0	0	1.95	0		
14	41	1	0	0	2			
15	1	1						
16		1 EO	1 EC	1 EO	1 EO	1 EO	1 EO	1 EO
17		1 EO	1 EO	1 EO	1 EO	1 EO	1 EO	1 EO
18		1 EO	1 EO	1 EO	1 EO	1 EO	1 EO	1 EO
19		1 EO	1 EO	1 EO	1 EO	1 EO	1 EO	1 EO
20		1 EO	1 EO	1 EO	1 EO	1 EO	1 EO	1 EO
21		1 EO	1 EO	1 EO	1 EO	1 EO	1 EO	1 EO
22		1 EO	1 EO	1 EO	1 EO	1 EO	1 EO	1 EO
23	1	2						
24	1	2						
25		0	3500	0				
26	1	40	0			1	40	
27	1							
28	11	875E-6						
29		1						
30	2		1					
31	15	625E-6						
32		1						
33	3		1					
34	19	375E-6						
35		1						
36	4		1					
37	23	125E-6						
38		1						
39	5		1					
40	26	875E-6						
41		1						
42	6		1					
43	30	625E-6						
44		1						
45	7		1					
46	34	375E-6						
47		1						
48	8		1					
49	38	125E-6						
50		1						
51	9		1					
52	41	875E-6						
53		1						
54	10		1					
55	45	625E-6						
56		1						
57	11		1					
58	49	375E-6						
59		1						
60	12		1					
61	53	125E-6						
62		1						
63	13		1					
64	56	875E-6						
65		1						
66	14		1					
67	60	625E-6						
68		1						
69	15		1					
70	64	375E-6						
71		1						
72	16		1					
73	68	125E-6						
74		1						
75	17		1					
76	71	875E-6						
77		1						
78	18		1					
79	75	625E-6						
80		1						
81	19		1					
82	79	375E-6						
83		1						
84	20		1					
85	83	125E-6						
86		1						
87	21		1					
88	86	875E-6						
89		1						
90	22		1					
91	90	625E-6						
92		1						
93	23		1					
94	94	375E-6						
95		1						
96	24		1					
97	98	125E-6						



98	1.							
99	25	1						
100	101.875E-6							
101	1.							
102	26	1.						
103	105.625E-6							
104	1.							
105	27	1.						
106	109.375E-6							
107	1.							
108	28	1						
109	113.125E-6							
110	1.							
111	29	1						
112	116.875E-6							
113	1							
114	30	1						
115	120.625E-6							
116	1.							
117	31	1.						
118	124.375E-6							
119	1.							
120	32	1.						
121	128.125E-6							
122	1.							
123	33	1.						
124	131.875E-6							
125	1.							
126	34	1						
127	135.625E-6							
128	1.							
129	35	1.						
130	139.375E-6							
131	1.							
132	36	1						
133	143.125E-6							
134	1.							
135	37	1.						
136	146.875E-6							
137	1.							
138	38	1.						
139	150.625E-6							
140	1.							
141	39	1.						
142	154.375E-6							
143	1.							
144	40	1.						
145	158.125E-6							
146	1.							
147	1	1	2	1				
148	2	2	3	2				
149	3	3	4	3				
150	4	4	5	4				
151	5	5	6	5				
152	6	6	7	6				
153	7	7	8	7				
154	8	8	9	8				
155	9	9	10	9				
156	10	10	11	10				
157	11	11	12	11				
158	12	12	13	12				
159	13	13	14	13				
160	14	14	15	14				
161	15	15	16	15				
162	16	16	17	16				
163	17	17	18	17				
164	18	18	19	18				
165	19	19	20	19				
166	20	20	21	20				
167	21	21	22	21				
168	22	22	23	22				
169	23	23	24	23				
170	24	24	25	24				
171	25	25	26	25				
172	26	26	27	26				
173	27	27	28	27				
174	28	28	29	28				
175	29	29	30	29				
176	30	30	31	30				
177	31	31	32	31				
178	32	32	33	32				
179	33	33	34	33				
180	34	34	35	34				
181	35	35	36	35				
182	36	36	37	36				
183	37	37	38	37				
184	38	38	39	38				
185	39	39	40	39				
186	40	40	41	40				
187	STOP							
CARD	1	2	3	4	5	6	7	8
NUMBER	123456789012345678901234567890123456789012345678901234567890							
	C O L U M N				N U M B E R			



## A.6 Polynomial variation of permeability

CARD NUMBER	1	2	3	4	5	6	7	8
1	1	2	3	4	5	6	7	8
2	1	2	3	4	5	6	7	8
3	1	2	3	4	5	6	7	8
4	1	2	3	4	5	6	7	8
5	1	2	3	4	5	6	7	8
6	1	2	3	4	5	6	7	8
7	1	2	3	4	5	6	7	8
8	1	2	3	4	5	6	7	8
9	1	2	3	4	5	6	7	8
10	1	2	3	4	5	6	7	8
11	1	2	3	4	5	6	7	8
12	1	2	3	4	5	6	7	8
13	1	2	3	4	5	6	7	8
14	1	2	3	4	5	6	7	8
15	1	2	3	4	5	6	7	8
16	1	2	3	4	5	6	7	8
17	1	2	3	4	5	6	7	8
18	1	2	3	4	5	6	7	8
19	1	2	3	4	5	6	7	8
20	1	2	3	4	5	6	7	8
21	1	2	3	4	5	6	7	8
22	1	2	3	4	5	6	7	8
23	1	2	3	4	5	6	7	8
24	1	2	3	4	5	6	7	8
25	1	2	3	4	5	6	7	8
26	1	2	3	4	5	6	7	8
27	1	2	3	4	5	6	7	8
28	1	2	3	4	5	6	7	8
29	1	2	3	4	5	6	7	8
30	1	2	3	4	5	6	7	8
31	1	2	3	4	5	6	7	8
32	1	2	3	4	5	6	7	8
33	1	2	3	4	5	6	7	8
34	1	2	3	4	5	6	7	8
35	1	2	3	4	5	6	7	8
36	1	2	3	4	5	6	7	8
37	1	2	3	4	5	6	7	8
38	1	2	3	4	5	6	7	8
39	1	2	3	4	5	6	7	8
40	1	2	3	4	5	6	7	8
41	1	2	3	4	5	6	7	8
42	1	2	3	4	5	6	7	8
43	1	2	3	4	5	6	7	8
44	1	2	3	4	5	6	7	8
45	1	2	3	4	5	6	7	8
46	1	2	3	4	5	6	7	8
47	1	2	3	4	5	6	7	8
48	1	2	3	4	5	6	7	8
49	1	2	3	4	5	6	7	8
50	1	2	3	4	5	6	7	8
51	1	2	3	4	5	6	7	8
52	1	2	3	4	5	6	7	8
53	1	2	3	4	5	6	7	8
54	1	2	3	4	5	6	7	8
55	1	2	3	4	5	6	7	8
56	1	2	3	4	5	6	7	8
57	1	2	3	4	5	6	7	8
58	1	2	3	4	5	6	7	8
59	1	2	3	4	5	6	7	8
60	1	2	3	4	5	6	7	8
61	1	2	3	4	5	6	7	8
62	1	2	3	4	5	6	7	8
63	1	2	3	4	5	6	7	8
64	1	2	3	4	5	6	7	8
65	1	2	3	4	5	6	7	8
66	1	2	3	4	5	6	7	8
67	1	2	3	4	5	6	7	8
68	1	2	3	4	5	6	7	8
69	1	2	3	4	5	6	7	8
70	1	2	3	4	5	6	7	8
71	1	2	3	4	5	6	7	8
72	1	2	3	4	5	6	7	8
73	1	2	3	4	5	6	7	8
74	1	2	3	4	5	6	7	8
75	1	2	3	4	5	6	7	8
76	1	2	3	4	5	6	7	8
77	1	2	3	4	5	6	7	8
78	1	2	3	4	5	6	7	8
79	1	2	3	4	5	6	7	8
80	1	2	3	4	5	6	7	8
81	1	2	3	4	5	6	7	8
82	1	2	3	4	5	6	7	8
83	1	2	3	4	5	6	7	8
84	1	2	3	4	5	6	7	8
85	1	2	3	4	5	6	7	8
86	1	2	3	4	5	6	7	8
87	1	2	3	4	5	6	7	8
88	1	2	3	4	5	6	7	8
89	1	2	3	4	5	6	7	8
90	1	2	3	4	5	6	7	8
91	1	2	3	4	5	6	7	8
92	1	2	3	4	5	6	7	8
93	1	2	3	4	5	6	7	8
94	1	2	3	4	5	6	7	8
95	1	2	3	4	5	6	7	8
96	1	2	3	4	5	6	7	8
97	1	2	3	4	5	6	7	8



98		1						
99	25	1.						
100	.155662633							
101	1.							
102	26	1						
103	.150318459							
104	1.							
105	27	1.						
106	.145321463							
107	1.							
108	28	1.						
109	.140639112							
110	1.							
111	29	1.						
112	.136242799							
113	1.							
114	30	1.						
115	.132107271							
116	1.							
117	31	1.						
118	.128210153							
119	1.							
120	32	1						
121	.124531550							
122	1.							
123	33	1.						
124	.121053713							
125	1.							
126	34	1.						
127	.117760758							
128	1.							
129	35	1.						
130	.114638427							
131	1.							
132	36	1.						
133	.111673888							
134	1.							
135	37	1.						
136	.108855558							
137	1.							
138	38	1.						
139	.106172958							
140	1							
141	39	1						
142	.103616582							
143	1.							
144	40	1.						
145	.101177790							
146	1.							
147	1	1	2	1				
148	2	2	3	2				
149	3	3	4	3				
150	4	4	5	4				
151	5	5	6	5				
152	6	6	7	6				
153	7	7	8	7				
154	8	8	9	8				
155	9	9	10	9				
156	10	10	11	10				
157	11	11	12	11				
158	12	12	13	12				
159	13	13	14	13				
160	14	14	15	14				
161	15	15	16	15				
162	16	16	17	16				
163	17	17	18	17				
164	18	18	19	18				
165	19	19	20	19				
166	20	20	21	20				
167	21	21	22	21				
168	22	22	23	22				
169	23	23	24	23				
170	24	24	25	24				
171	25	25	26	25				
172	26	26	27	26				
173	27	27	28	27				
174	28	28	29	28				
175	29	29	30	29				
176	30	30	31	30				
177	31	31	32	31				
178	32	32	33	32				
179	33	33	34	33				
180	34	34	35	34				
181	35	35	36	35				
182	36	36	37	36				
183	37	37	38	37				
184	38	38	39	38				
185	39	39	40	39				
186	40	40	41	40				
187	STOP							
CARD	1	2	3	4	5	6	7	8
NUMBER	1234567890123456789012345678901234567890123456789012345678901234567890							
	C O L U M N				N U M B E R			





A.7 Sinusoidal variation of permeability

		C O L U M N   N U M B E R							
CARD	1	2	3	4	5	6	7	8	
NUMBER	123456789012345678901234567890123456789012345678901234567890								
1	1-D CONSOLIDATION (SIN VAR K, D DR) DATA Schiffman & Gibson 64								
2	41	1	0	1	1	0	10	0	
3	1	0							
4									
5									
6									
7									
8									
9	10								
10	7412								
11	1	1	0	0	0	0	0		
12	2	0	0	0	0	5	1		
13	40	0	0	0	0	19	5	0	
14	41	1	0	0	0	20			
15	1	1							
16		1 EO	1 EO	1 EO	1 EO	1 EO	1 EO	1 EO	
17		1 EO	1 EO	1 EO	1 EO	1 EO	1 EO	1 EO	
18		1 EO	1 EO	1 EO	1 EO	1 EO	1 EO	1 EO	
19		1 EO	1 EO	1 EO	1 EO	1 EO	1 EO	1 EO	
20		1 EO	1 EO	1 EO	1 EO	1 EO	1 EO	1 EO	
21		1 EO	1 EO	1 EO	1 EO	1 EO	1 EO	1 EO	
22		1 EO	1 EO	1 EO	1 EO	1 EO	1 EO	1 EO	
23	1	2							
24	1	2							
25		0	0	50	0				
26	1	40					1	40	
27	1		1						
28	1	35333834							
29	1								
30	2		1						
31	2	05783658							
32	1								
33	3		1						
34	2	75581290							
35	1								
36	4		1						
37	3	44286405							
38	1								
39	11		1						
40	4	11505351							
41	1								
42	6		1						
43	4	76793764							
44	1								
45	7		1						
46	5	39759117							
47	1								
48	8		1						
49	6	00013210							
50	1								
51	11		1						
52	6	57184554							
53	1								
54	10		1						
55	7	10920671							
56	1								
57	11		1						
58	7	60890258							
59	1								
60	12		1						
61	8	06785238							
62	1								
63	13		1						
64	8	48322651							
65	1								
66	14		1						
67	8	85246406							
68	1								
69	15		1						
70	9	17328856							
71	1								
72	16		1						
73	9	44372202							
74	1								
75	17		1						
76	9	66209713							
77	1								
78	18		1						
79	9	82706752							
80	1								
81	19		1						
82	9	93761611							
83	1								
84	20		1						
85	9	99306133							
86	1								
87	21		1						
88	9	99306133							
89	1								
90	22		1						
91	9	93761611							
92	1								
93	23		1						
94	9	82706752							
95	1								
96	24		1						
97	9	66209713							







A.8 Linear variation of modulus of elasticity

CARD NUMBER	C O L U M N   N U M B E R							
	1	2	3	4	5	6	7	8
1	1-D CONSOLIDATION (VAR MV, D, DR) DATA Edelmann53							
2	41	1	0	1	1	0	10	0
3	2	0						
4								
5								
6								
7								
8								
9	10							
10	256	1171						
11	1	1	0	0	0	0		
12	2	0	0	0	.0525	1		
13	41	0	0	0	2	1		
14	1	1						
15		1 EO	1 EO	1 EO	1 EO	1 EO	1 EO	1 EO
16		1 EO	1 EO	1 EO	1 EO	1 EO	1 EO	1 EO
17		1 EO	1 EO	1 EO	1 EO	1 EO	1 EO	1 EO
18		1 EO	1 EO	1 EO	1 EO	1 EO	1 EO	1 EO
19		1 EO	1 EO	1 EO	1 EO	1 EO	1 EO	1 EO
20		1 EO	1 EO	1 EO	1 EO	1 EO	1 EO	1 EO
21		1 EO	1 EO	1 EO	1 EO	1 EO	1 EO	1 EO
22	1	2						
23	1	2						
24		0	0	3000	0			
25	1	40	0				1	40
26	1							
27	1 E-5							
28	156402737							
29	2	1						
30	1 E-5							
31	139251523							
32	3	1						
33	1 E-5							
34	125490196							
35	4	1						
36	1 E-5							
37	114204140							
38	5	1						
39	1 E-5							
40	104780616							
41	8	1						
42	1 E-5							
43	086793706							
44	7	1						
45	1 E-5							
46	089938168							
47	8	1						
48	1 E-5							
49	083889501							
50	9	1						
51	1 E-5							
52	078778927							
53	10	1						
54	1 E-5							
55	074177088							
56	11	1						
57	1 E-5							
58	070083224							
59	12	1						
60	1 E-5							
61	066417601							
62	13	1						
63	1 E-5							
64	063116371							
65	14	1						
66	1 E-5							
67	060127772							
68	15	1						
69	1 E-5							
70	057409401							
71	16	1						
72	1 E-5							
73	054926193							
74	17	1						
75	1 E-5							
76	052648898							
77	18	1						
78	1 E-5							
79	050552923							
80	19	1						
81	1 E-5							
82	048617442							
83	20	1						
84	1 E-5							
85	046824700							
86	21	1						
87	1 E-5							
88	045159469							
89	22	1						
90	1 E-5							
91	043608613							
92	23	1						
93	1 E-5							
94	042160738							
95	24	1						
96	1 E-5							
97	040805917							



98	25	1.						
99	1. E-5							
100	.039535458							
101	26	1.						
102	1. E-5							
103	.038341721							
104	27	1.						
105	1. E-5							
106	.037217958							
107	28	1.						
108	1. E-5							
109	.036158192							
110	29	1.						
111	1. E-5							
112	.035157108							
113	30	1.						
114	1. E-5							
115	.034209964							
116	31	1.						
117	1. E-5							
118	.033312513							
119	32	1.						
120	1. E-5							
121	.032460945							
122	33	1.						
123	1. E-5							
124	.031651830							
125	34	1.						
126	1. E-5							
127	.030882069							
128	35	1.						
129	1. E-5							
130	.030148860							
131	36	1.						
132	1. E-5							
133	.029449659							
134	37	1.						
135	1. E-5							
136	.028782155							
137	38	1.						
138	1. E-5							
139	.028144239							
140	39	1.						
141	1. E-5							
142	.027533987							
143	40	1.						
144	1. E-5							
145	.026949638							
146	1	1	2	1				
147	2	2	3	2				
148	3	3	4	3				
149	4	4	5	4				
150	5	5	6	5				
151	6	6	7	6				
152	7	7	8	7				
153	8	8	9	8				
154	9	9	10	9				
155	10	10	11	10				
156	11	11	12	11				
157	12	12	13	12				
158	13	13	14	13				
159	14	14	15	14				
160	15	15	16	15				
161	16	16	17	16				
162	17	17	18	17				
163	18	18	19	18				
164	19	19	20	19				
165	20	20	21	20				
166	21	21	22	21				
167	22	22	23	22				
168	23	23	24	23				
169	24	24	25	24				
170	25	25	26	25				
171	26	26	27	26				
172	27	27	28	27				
173	28	28	29	28				
174	29	29	30	29				
175	30	30	31	30				
176	31	31	32	31				
177	32	32	33	32				
178	33	33	34	33				
179	34	34	35	34				
180	35	35	36	35				
181	36	36	37	36				
182	37	37	38	37				
183	38	38	39	38				
184	39	39	40	39				
185	40	40	41	40				
186	STOP							
CARD	1	2	3	4	5	6	7	8
NUMBER	1234567890123456789012345678901234567890123456789012345678901234567890							
	C O L U M N	N U M B E R						





A.9 Sinusoidal variation of compressibility

CARD NUMBER	C O L U M N N U M B E R							
	1	2	3	4	5	6	7	8
1	1-D CONSOLIDATION (SIN VAR MV, D DR) DATA: Schiffman & Gibson 64							
2	41	1	0	1	1	0	10	0
3	1	0						
4								
5								
6								
7								
8								
9	10							
10	4316							
11	1	1	0	0	0	0	0	
12	2	0	0	0	0	5	1	
13	40	0	0	0	0	19	5	0
14	41	1	0	0	0	20		
15	1	1						
16		1.E0	1.E0	1.E0	1.E0	1.E0	1.E0	1.E0
17		1.E0	1.E0	1.E0	1.E0	1.E0	1.E0	1.E0
18		1.E0	1.E0	1.E0	1.E0	1.E0	1.E0	1.E0
19		1.E0	1.E0	1.E0	1.E0	1.E0	1.E0	1.E0
20		1.E0	1.E0	1.E0	1.E0	1.E0	1.E0	1.E0
21		1.E0	1.E0	1.E0	1.E0	1.E0	1.E0	1.E0
22		1.E0	1.E0	1.E0	1.E0			
23	1	2						
24	1	2						
25		0		50	0			
26	1	40					1	40
27	1		1					
28		1						
29	964666166							
30	2		1					
31		1						
32	894216342							
33	3		1					
34		1						
35	824418710							
36	4		1					
37		1						
38	755703595							
39	5		1					
40		1						
41	688494649							
42	E		1					
43		1						
44	623206236							
45	7		1					
46		1						
47	560240823							
48	8		1					
49		1						
50	499986790							
51	9		1					
52		1						
53	442815446							
54	10		1					
55		1						
56	389078329							
57	11		1					
58		1						
59	339109742							
60	12		1					
61		1						
62	293214762							
63	13		1					
64		1						
65	251677349							
66	14		1					
67		1						
68	214753594							
69	15		1					
70		1						
71	182671144							
72	16		1					
73		1						
74	155627798							
75	17		1					
76		1						
77	133790287							
78	18		1					
79		1						
80	117293248							
81	19		1					
82		1						
83	106238389							
84	20		1					
85		1						
86	100693867							
87	21		1					
88		1						
89	100693867							
90	22		1					
91		1						
92	106238389							
93	23		1					
94		1						
95	117293248							
96	24		1					
97		1						



98	.133790287							
99	25	1.						
100		1.						
101	.155627798							
102	26	1.						
103		1.						
104	.182671144							
105	27	1.						
106		1.						
107	.214753594							
108	28	1.						
109		1.						
110	.251677349							
111	29	1.						
112		1.						
113	.293214762							
114	30	1.						
115		1.						
116	.339109742							
117	31	1.						
118		1.						
119	.389079329							
120	32	1.						
121		1.						
122	.442815446							
123	33	1.						
124		1.						
125	.499986790							
126	34	1.						
127		1.						
128	.560240883							
129	35	1.						
130		1.						
131	.623206236							
132	36	1.						
133		1.						
134	.688494649							
135	37	1.						
136		1.						
137	.755703595							
138	38	1.						
139		1.						
140	.824418710							
141	39	1.						
142		1.						
143	.894216342							
144	40	1.						
145		1.						
146	.964666166							
147	1	1	2	1				
148	2	2	3	2				
149	3	3	4	3				
150	4	4	5	4				
151	5	5	6	5				
152	6	6	7	6				
153	7	7	8	7				
154	8	8	9	8				
155	9	9	10	9				
156	10	10	11	10				
157	11	11	12	11				
158	12	12	13	12				
159	13	13	14	13				
160	14	14	15	14				
161	15	15	16	15				
162	16	16	17	16				
163	17	17	18	17				
164	18	18	19	18				
165	19	19	20	19				
166	20	20	21	20				
167	21	21	22	21				
168	22	22	23	22				
169	23	23	24	23				
170	24	24	25	24				
171	25	25	26	25				
172	26	26	27	26				
173	27	27	28	27				
174	28	28	29	28				
175	29	29	30	29				
176	30	30	31	30				
177	31	31	32	31				
178	32	32	33	32				
179	33	33	34	33				
180	34	34	35	34				
181	35	35	36	35				
182	36	36	37	36				
183	37	37	38	37				
184	38	38	39	38				
185	39	39	40	39				
186	40	40	41	40				
187	STOP							
CARD	1	2	3	4	5	6	7	8
NUMBER	1234567890123456789012345678901234567890123456789012345678901234567890							
	C O L U M N			N U M B E R				



A.10 Constant coefficient of consolidation with variable permeability and compressibility

CARD NUMBER	C O L U M N   N U M B E R							
	1	2	3	4	5	6	7	8
1	123456789012345678901234567890123456789012345678901234567890							
2	1-D CONSOLIDATION (VAR K, MV, CONST CV, D DR) DATA: Schiffman & Gibson 64							
3	41	1	0	1	1	0.	10	0
4	1	0						
5								
6								
7								
8								
9	10							
10	.838							
11	1	1	0.		0.	0	0	
12	2	0	0.		0	.5	1	
13	41	0	0		0	19.5	0	
14	41	1	0		0	20		
15	1	1						
16		1 EO		1 EO		1 EO		1 EO
17		1 EO		1 EO		1 EO		1 EO
18		1 EO		1 EO		1 EO		1 EO
19		1 EO		1 EO		1 EO		1 EO
20		1 EO		1 EO		1 EO		1 EO
21		1 EO		1 EO		1 EO		1 EO
22		1 EO		1 EO		1 EO		1 EO
23	1	2						
24	1	2						
25		0		0	50	0		
26	1	40					1	40
27	1		1					
28	90495							
29	90495							
30	2		1					
31	75961							
32	75961							
33	3		1					
34	65383							
35	65383							
36	4		1					
37	57346							
38	57346							
39	5		1					
40	51036							
41	51036							
42	6		1					
43	45855							
44	45855							
45	7		1					
46	41776							
47	41776							
48	8		1					
49	38281							
50	38281							
51	9		1					
52	35314							
53	35314							
54	10		1					
55	32766							
56	32766							
57	11		1					
58	30555							
59	30555							
60	12		1					
61	28617							
62	28617							
63	13		1					
64	26906							
65	26906							
66	14		1					
67	25384							
68	25384							
69	15		1					
70	24021							
71	24021							
72	16		1					
73	22795							
74	22795							
75	17		1					
76	21685							
77	21685							
78	18		1					
79	20676							
80	20676							
81	19		1					
82	18755							
83	18755							
84	20		1					
85	18911							
86	18911							
87	21		1					
88	18134							
89	18134							
90	22		1					
91	17418							
92	17418							
93	23		1					
94	16755							
95	16755							
96	24		1					
97	16139							



98	.16139							
99	25	1						
100	.15566							
101	.15566							
102	26	1.						
103	.15032							
104	.15032							
105	27	1.						
106	.14532							
107	.14532							
108	28	1.						
109	.14064							
110	.14064							
111	29	1.						
112	.13624							
113	.13624							
114	30	1.						
115	.13211							
116	.13211							
117	31	1.						
118	.12821							
119	.12821							
120	32	1.						
121	.12453							
122	.12453							
123	33	1.						
124	.12105							
125	.12105							
126	34	1.						
127	.11776							
128	.11776							
129	35	1.						
130	.11464							
131	.11464							
132	36	1.						
133	.11167							
134	.11167							
135	37	1.						
136	.10886							
137	.10886							
138	38	1.						
139	.10617							
140	.10617							
141	39	1.						
142	.10362							
143	.10362							
144	40	1.						
145	.10118							
146	.10118							
147	1	1	2	1				
148	2	2	3	2				
149	3	3	4	3				
150	4	4	5	4				
151	5	5	6	5				
152	6	6	7	6				
153	7	7	8	7				
154	8	8	9	8				
155	9	9	10	9				
156	10	10	11	10				
157	11	11	12	11				
158	12	12	13	12				
159	13	13	14	13				
160	14	14	15	14				
161	15	15	16	15				
162	16	16	17	16				
163	17	17	18	17				
164	18	18	19	18				
165	19	19	20	19				
166	20	20	21	20				
167	21	21	22	21				
168	22	22	23	22				
169	23	23	24	23				
170	24	24	25	24				
171	25	25	26	25				
172	26	26	27	26				
173	27	27	28	27				
174	28	28	29	28				
175	29	29	30	29				
176	30	30	31	30				
177	31	31	32	31				
178	32	32	33	32				
179	33	33	34	33				
180	34	34	35	34				
181	35	35	36	35				
182	36	36	37	36				
183	37	37	38	37				
184	38	38	39	38				
185	39	39	40	39				
186	40	40	41	40				
187	STOP							
CARD	1	2	3	4	5	6	7	8
NUMBER	1234567890123456789012345678901234567890123456789012345678901234567890							
	C O L U M N				N U M B E R			





A.11 Thin layer: Samarasinghe et al (1982)

CARD NUMBER	C O L U M N N U M B E R															
	1	2	3	4	5	6	7	8								
1	1-D CONSOLDN (SMALL STR, THIN LVR, D.DR) DATA: Samarasinghe-etal182 EQ A1															
2	21	0	1	1	1	0	10									
3	1	0														
4																
5																
6																
7																
8																
9	10															
10	.000401															
11	1	1	0		0	0		0								
12	2	0	0		0	.00098552	1									
13	21	1	0		0	.0197104										
14	1	1														
15	47	88028	47	88028	47	88028	47	88028	47	88028	47	88028	47	88028	47	88028
16	47	88028	47	88028	47	88028	47	88028	47	88028	47	88028	47	88028	47	88028
17	47	88028	47	88028	47	88028	47	88028	47	88028	47	88028	47	88028	47	88028
18	47	88028														
19	1	2	0	0												
20	1	2														
21			0		10	0										
22	1	20	1					4	1	44	44					
23	1		1													
24		0		2	4	4	8	7	2	8	6	12	14	4	16	8
25		19	2	21	6	24	26	4		28	8	31	2	33	6	36
26		38	4	40	8	43	2	45	6	47	88028	48	2	32918E-5	2	364E-5
27	2	40014E-52	43771E-52	47678E-52	51747E-52	55988E-52	60415E-5	2	6504E-52	69881E-5						
28	2	74953E-52	80275E-5	2	8587E-52	91761E-52	97975E-53	04543E-5	3	115E-53	18887E-5					
29	3	26749E-5	3	3514E-53	43659E-53	44122E-5										
30		0	2	4	8	7	2	8	6	12	14	4	16	8		
31		19	2	21	6	24	26	4		28	8	31	2	33	6	36
32		38	4	40	8	43	2	45	6	47	88028	48	4	3872E-34	49998E-3	
33	4	61871E-34	74388E-34	87602E-35	01573E-35	16369E-35	32064E-35	48743E-35	66502E-3							
34	5	85448E-36	05706E-36	27415E-36	50739E-36	75864E-37	03007E-37	32421E-37	64404E-3							
35	7	99308E-38	37552E-3	8	7744E-38	79639E-3										
36	1		2	1	1	0										
37	20	20	21	1	0	0										
38	STOP															

CARD NUMBER	C O L U M N N U M B E R																
	1	2	3	4	5	6	7	8									
1	1234567890123456789012345678901234567890123456789012345678901234567890																
2	1-D CONSOLDN (FINITE STR, THIN LVR, D.DR) DATA: Samarasinghe-etal182 EQ A6																
3	21	0	1	1	1	0	10										
4	1	0															
5																	
6																	
7																	
8																	
9	10																
10	.000401																
11	1	1	0		0	0											
12	2	0	0		0	.00098552	1										
13	21	1	0		0	.0197104											
14	1	1															
15	47	88028	47	88028	47	88028	47	88028	47	88028							
16	47	88028	47	88028	47	88028	47	88028	47	88028							
17	47	88028	47	88028	47	88028	47	88028	47	88028							
18	47	88028															
19	1	2	0	0													
20	1	2															
21	0		0		10	0											
22	1	20	1				4	1	44	44							
23	1	0	1														
24		19	2	2	4	4	8	7	2	8	6	12	14	4	16	8	
25		38	4	40	8	43	2	45	6	47	88028	48	2	40046E-52	43362E-5		
26																	
27	2	46799E-52	50367E-52	54073E-52	57926E-52	61937E-52	66118E-52	70479E-52	75036E-5								
28	2	79803E-52	84797E-52	90038E-52	95546E-53	01345E-53	07463E-5	3	1393E-53	20782E-5							
29	3	28059E-53	35809E-53	43659E-53	44085E-5												
30		0	2	4	8	7	2	8	6	12	14	4	16	8			
31		19	2	21	6	24	26	4	28	8	31	2	33	6	36		
32		38	4	40	8	43	2	45	6	47	88028	48	4	3872E-34	49998E-3		
33	4	61871E-34	74388E-34	87602E-35	01573E-35	16369E-35	32064E-35	48743E-35	66502E-3								
34	5	85448E-36	05706E-36	27415E-36	50739E-36	75864E-37	03007E-37	32421E-37	64404E-3								
35	7	99308E-38	37552E-3	8	7744E-38	79639E-3											
36	1	2	0	1	1	0											
37	20	20	21	1	0	0											
38	STOP																

CARD NUMBER	C O L U M N N U M B E R															
	1	2	3	4	5	6	7	8	9	10	11	12	13	14	15	16
1	1-D CONSOLDN (FINITE STR, THIN LVR, D.DR) DATA: Samarasinghe-etal182 EQ A6															
2	21	0	1	1	1	0	10									
3	1	0														
4																
5																
6																
7																
8																
9	10															
10	000401															
11	1	1	0		0	0										
12	2	0	0		0	00098552	1									
13	21	1	0		0	.0197104										
14	1	1														
15	47	88028	47	88028	47	88028	47	88028	47	88028	47	88028	47	88028	47	88028
16	47	88028	47	88028	47	88028	47	88028	47	88028	47	88028	47	88028	47	88028
17	47	88028	47	88028	47	88028	47	88028	47	88028	47	88028	47	88028	47	88028
18	47	88028														
19	1	2	0	0												
20	1	2														
21		0		0		10	0									
22	1	20	1										4	1	44	44
23	1		1													
24		0		2	4	4	8	7	2	8	6	12	14	4	16	8
25		18	2	21	6	24	26	4	28	8	31	2	33	6	36	
26		38	4	40	8	43	2	45	6	47	88028	48	2	40046E-52	43362E-5	
27	2	46799E-52	50367E-52	54073E-52	57926E-52	61937E-52	66118E-52	70479E-52	75036E-5							
28	2	79803E-52	84797E-52	90038E-52	95546E-53	01345E-53	07463E-5	3	1393E-53	20782E-5						
29	3	28059E-53	35809E-53	43659E-53	44085E-5											
30		0	2	4	8	7	2	8	6	12	14	4	16	8		
31		18	2	21	6	24	26	4	28	8	31	2	33	6	36	
32		38	4	40	8	43	2	45	6	47	88028	48	4	3872E-34	49998E-3	
33	4	61871E-34	74388E-34	87602E-35	01573E-35	16369E-35	32064E-35	48743E-35	66502E-3							
34	5	85448E-36	05706E-36	27415E-36	50739E-36	75864E-37	03007E-37	32421E-37	64404E-3							
35	7	99308E-38	37552E-3	8	7744E-38	79639E-3										
36	1	2	0	1	1	0										
37	20	20	21	1	0	0										
38	STOP															

CARD NUMBER	C O L U M N N U M B E R								
	1	2	3	4	5	6	7	8	
1	1234567890123456789012345678901234567890123456789012345678901234567890								
2	1-D CONSOLDN (FINITE STR, THIN LVR, D.DR) DATA: Samarasinghe-etal182 EQ A6								
3	21	0	1	1	1	0	10		
4	1	0							
5									
6									
7									
8									
9	10								
10	.000401								
11	1	1	0		0	0			
12	2	0	0		0	.00098552	1		
13	21	1	0		0	.0197104			
14	1	1							
15	47	88028	47	88028	47	88028	47	88028	47
16	47	88028	47	88028	47	88028	47	88028	47
17	47	88028	47	88028	47	88028	47	88028	47
18	47	88028							
19	1	2	0	0					



CARD NUMBER	1	2	3	4	5	6	7	8
1	1-D CONSOLDN (FINITE STR, THIN LVR, D.DR) DATA:Samarasinghe-etal82 EQ:B4							
2	21	0	1	1	1	0	10	
3	1	0						
4								
5								
6								
7								
8								
9	10							
10	.000401							
11	1	1	0.	0.	0.	0		
12	2	0	0.	0.	.00098552	1		
13	21	1	0.	0.	.0197104			
14	1	1						
15	47.88028	47.88028	47.88028	47.88028	47.88028	47.88028	47.88028	
16	47.88028	47.88028	47.88028	47.88028	47.88028	47.88028	47.88028	
17	47.88028	47.88028	47.88028	47.88028	47.88028	47.88028	47.88028	
18	47.88028							
19	1	2	0	0				
20	1	2						
21	1	0.	0.	10.	0.			
22	1	20	1				4	1 44 44
23	1	1						
24	0.	2.4	4.8	7.2	9.6	12.	14.4	16.8
25	19.2	21.6	24.	26.4	28.8	31.2	33.6	36.
26	38.4	40.8	43.2	45.6	47.88028	48.5	71374E-55	79266E-5
27	5.87449E-5	5.9594E-56	0.4761E-56	1.3934E-56	2.3481E-56	3.3431E-56	4.3813E-5	6.5466E-5
28	6.66006E-56	7.77894E-56	9.0367E-57	0.3478E-57	1.7282E-57	3.1844E-57	4.7238E-57	6.3547E-5
29	7.80869E-57	9.9315E-58	1.8001E-58	1.9016E-5				
30	0.	2.4	4.8	7.2	9.6	12.	14.4	16.8
31	19.2	21.6	24.	26.4	28.8	31.2	33.6	36.
32	38.4	40.8	43.2	45.6	47.88028	48.	010442713	010711161
33	.010993776	.011291708	.011606238	.011938793	.012290968	.012664551	.013061555	.013484256
34	.013935231	.014417415	.014934164	.015489333	.016087371	.016733444	.017433582	.018194866
35	.019025672	.019935982	.020885425	.020937778				
36	1	1	2	1	1	0.		
37	20	20	21	1	0	0.		
38	STOP							

CARD NUMBER	1	2	3	4	5	6	7	8
1	C O L U M N N U M B E R							

CARD NUMBER	1	2	3	4	5	6	7	8
1	1-D CONSOLDN (FINITE STR, THIN LVR, D.DR) DATA:Samarasinghe-etal82 EQ:B5							
2	21	0	1	1	1	0	10	
3	1	0						
4								
5								
6								
7								
8								
9	10							
10	.000401							
11	1	0	0.	0.	0.	1		
12	21	0	0.	0.	.0197104			
13	1	1						
14	.622	.622	.622	.622	.622	.622	.622	
15	.622	.622	.622	.622	.622	.622	.622	
16	.622	.622	.622	.622	.622	.622	.622	
17	.622	.622	.622	.622	.622	.622	.622	
18	1	2	0	2				
19	1	2						
20	0.	1.	10.	1.				
21	1	1.573835201	0.					
22	21	1.573835201	0.					
23	1	20	1				2	1 44 0
24	1	1.						
25	.573	.573835201	.57545	.5779	.58035	.5828	.58525	.5877
26	.59015	.5926	.59505	.5975	.59995	.6024	.60485	.6073
27	.60975	.6122	.61465	.6171	.61955	.622	.005501723	.005471507
28	.005413397	.005326013	.005239574	.005154086	.005069551	.004985972	.004903351	.00482169
29	.00474089	.004661251	.004582475	.00450466	.004427807	.004351914	.004276979	.004203002
30	.004129979	.004057908	.003986787	.003916611				
31	1.							
32	1	1	2	1	1	0.		
33	20	20	21	1	0	0.		
34	STOP							

CARD NUMBER	1	2	3	4	5	6	7	8
1	C O L U M N N U M B E R							



CARD NUMBER	1		2		3		4		5		6		7		8	
	1	2	3	4	5	6	7	8	9	10	11	12	13	14	15	16
1	I-D CONSOLDN (FINITE STR, THIN LTR, D.DR) DATA: Samarasinghe-etal82 EO:B6															
2	21	0	1	1	1	0	10									
3	1	0														
4																
5																
6																
7																
8																
9	10															
10	.000401															
11	1	0		0		0		0		1						
12	21	0		0		0		.0197104								
13	0	1														
14	1	2		0		2										
15	1	2														
16		0			1		10		1							
17	1	1		1		0										
18	21	1		1		0										
19	1	20	1													
20	1		1											2	1	42 0
21		0		.05		.1		.15		.2		.25		.3		.35
22		.4		.45		.5		.55		.6		.65		.7		.75
23		.8		.85		.9		.95		1		.003916611		.003985582		.004055468
24	.004126271	.004197993	.004270637	.004344206	.004418702	.004494124	.004570475	.004647756								
25	.004725966	.004805105	.004885173	.004966169	.005048091	.005130937	.005214705	.005299392								
26	.005384994	.005471507														
27	1															
28	1		1	2		1	1	0								
29	20		20	21		1	0									
30	STOP															



## A.12 Thin layer: Barden-Berry (1965)

CARD NUMBER	1	2	3	4	5	6	7	8
1	1234567890123456789012345678901234567890123456789012345678901234567890							
2	1-D CONSOLDN (SMALL STR, THIN Lyr, D DR) DATA Barden&Berry65 (n=1) EO A1							
3	21	0	1	1	1	0	10	
4	1	0						
5								
6								
7								
8								
9	10							
10	.0055437							
11	1	1	0	0	0	0		
12	2	0	0	0	.003175	1		
13	21	1	0	0	.0635			
14	1	1						
15	107	2517799	107	2517799	107	2517799	107	2517799
16	107	2517799	107	2517799	107	2517799	107	2517799
17	107	2517799	107	2517799	107	2517799	107	2517799
18	107	2517799						
19	1	2	0	0				
20	1	2						
21		0		0	10	0		
22	1	20	1				4	1 4 46
23	1		1					
24		0		110 2	08767E-54	22884E-5		
25		0		5	10	15	20	25 30 35
26		40		45	50	55	60	65 70 75
27		80		85	90	95	100	105 110 4 05411E-3
28	4	15087E-34	25236E-34	35893E-34	47098E-34	58895E-34	71331E-34	8446E-34 98341E-3
29	5	13041E-35	28634E-35	45205E-35	62849E-35	81673E-35	1018E-36	23369E-36 46542E-3
30	6	71504E-36	98472E-37	27696E-37	59472E-37	9415E-38	32146E-3	
31	1	1	2	1	1	0		
32	20	20	21	1	0	0		
33	STDP							

CARD NUMBER	1	2	3	4	5	6	7	8
1	1234567890123456789012345678901234567890123456789012345678901234567890							
2	1-D CONSOLDN (SMALL STR, THIN Lyr, D DR) DATA Barden&Berry65 (n=1/2) EO A1							
3	21	0	1	1	1	0	10	
4	1	0						
5								
6								
7								
8								
9	10							
10	.0055437							
11	1	1	0	0	0	0		
12	2	0	0	0	.003175	1		
13	21	1	0	0	.0635			
14	1	1						
15	107	2517799	107	2517799	107	2517799	107	2517799
16	107	2517799	107	2517799	107	2517799	107	2517799
17	107	2517799	107	2517799	107	2517799	107	2517799
18	107	2517799						
19	1	2	0	0				
20	1	2						
21		0		0	10	0		
22	1	20	1				4	1 46 46
23	1		1					
24		0		5	10	15	20	25 30 35
25		40		45	50	55	60	65 70 75
26		80		85	90	95	100	105 110 2 08767E-5
27	2	53843E-52	72514E-52	86841E-52	98919E-53	10856E-53	1818E-53	28027E-53 36261E-5
28	3	43985E-53	5131E-53	582E7E-53	64815E-53	71291E-53	77426E-53	83346E-53 89071E-5
29	3	94620E-54	00008E-54	05249E-54	10353E-54	15331E-54	20192E-5	
30		0		5	10	15	20	25 30 35
31		40		45	50	55	60	65 70 75
32		80		85	90	95	100	105 110 4 05411E-3
33	4	15087E-34	25236E-34	35893E-34	47098E-34	58895E-34	71331E-34	8446E-34 98341E-3
34	5	13041E-35	28634E-35	45205E-35	62849E-35	81673E-35	1018E-36	23369E-36 46542E-3
35	6	71504E-36	98472E-37	27696E-37	59472E-37	9415E-38	32146E-3	
36	1	1	2	1	1	0		
37	20	20	21	1	0	0		
38	STDP							













CARD NUMBER	C O L U M N   N U M B E R							
	1	2	3	4	5	6	7	8
1	1	2	3	4	5	6	7	8
2	1	2	3	4	5	6	7	8
3	1	2	3	4	5	6	7	8
4	1	2	3	4	5	6	7	8
5	1	2	3	4	5	6	7	8
6	1	2	3	4	5	6	7	8
7	1	2	3	4	5	6	7	8
8	1	2	3	4	5	6	7	8
9	1	2	3	4	5	6	7	8
10	1	2	3	4	5	6	7	8
11	1	2	3	4	5	6	7	8
12	1	2	3	4	5	6	7	8
13	1	2	3	4	5	6	7	8
14	1	2	3	4	5	6	7	8
15	1	2	3	4	5	6	7	8
16	1	2	3	4	5	6	7	8
17	1	2	3	4	5	6	7	8
18	1	2	3	4	5	6	7	8
19	1	2	3	4	5	6	7	8
20	1	2	3	4	5	6	7	8
21	1	2	3	4	5	6	7	8
22	1	2	3	4	5	6	7	8
23	1	2	3	4	5	6	7	8
24	1	2	3	4	5	6	7	8
25	1	2	3	4	5	6	7	8
26	1	2	3	4	5	6	7	8
27	1	2	3	4	5	6	7	8
28	1	2	3	4	5	6	7	8
29	1	2	3	4	5	6	7	8
30	1	2	3	4	5	6	7	8

1-D CONSOLDN (FINITE STR, THIN LVR, D.DR) DATA: Barden&Berry65 EQ:B6

10

0055437

1 0 0. 0. 0. 1

21 0 0. 0. .0635

0 1

1 2 0 2

1 2

0. 1. 10. 1

1 1 1. 0. 0

21 1 1. 0.

1 20 1 2 1 42 0

1 1.

0. .05 .1 .15 .2 .25 .3 .35

.4 .45 .5 .55 .6 .65 .7 .75

.8 .85 .9 .95 1. .002574756 .002673772 .002776622

.002883456 .002994429 .003109703 .003229447 .003353834 .003483047 .003617273 .00375671

.003901561 .004052038 .004208362 .004370761 .004539474 .004714748 .004896841 .00508602

.005282564 .005486762

1.

1 1 2 1 1 0.

20 20 21 1 0 0.

STOP

C O L U M N   N U M B E R



## A.13 Thin layer: Davis-Raymond (1965)

CARD NUMBER	1	2	3	4	5	6	7	8
1	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890
2	1-D CONSOLDN (SMALL STR, THIN LVR, D.DR) DATA: Davis&Raymond65 EQ: A1							
3	21	0	1	1	1	0	10	
4	1	0						
5								
6								
7								
8								
9	10							
10	.005995618							
11	1	1	0	0	0	0	1	
12	2	0	0	0	.001155			
13	21	1	0	0	.0231			
14	1	1						
15		562.5	562.5	562.5	562.5	562.5	562.5	562.5
16		562.5	562.5	562.5	562.5	562.5	562.5	562.5
17		562.5	562.5	562.5	562.5	562.5	562.5	562.5
18		562.5						
19	1	2	0	0				
20	1	2						
21	1	0	0	10	0			
22	1	20	1			4	1	42 42
23	1	0	1					
24		0	28.5	57	85.5	114	142.5	171
25		228	255.5	285	313.5	342	370.5	399
26		456	484.5	513	541.5	570	598.5	627
27	1	80135E-51	88207E-51	97038E-52	06735E-52	17438E-5	2.2931E-52	42552E-52
28	2	74225E-52	83381E-53	15414E-53	41024E-53	71152E-54	07141E-54	50846E-55
29	5	74099E-55	64998E-5					
30		0	28.5	57	85.5	114	142.5	171
31		228	255.5	285	313.5	342	370.5	399
32		456	484.5	513	541.5	570	598.5	627
33	2	02399E-32	11468E-32	21389E-32	32287E-32	44312E-32	57651E-32	72531E-32
34	3	08119E-33	28542E-33	54398E-33	83174E-34	17035E-34	57462E-35	08568E-35
35	5	45055E-37	47188E-3					
36	1	1	2	1	1	0		
37	20	20	21	1	0			
38	STDP							

CARD NUMBER	1	2	3	4	5	6	7	8
1	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890
2	C D L U M N N U M B E R							

CARD NUMBER	1	2	3	4	5	6	7	8
1	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890
2	1-D CONSOLDN (SMALL STR, THIN LVR, D.DR) DATA: Davis&Raymond65 EQ: A2							
3	21	1	0	1	1	0	10	
4	1	0						
5								
6								
7								
8								
9	10							
10	.005995618							
11	1	0	0	0	0	0	1	
12	21	0	0	0	.0231			
13	1	1						
14		9	9	9	9	9	9	9
15		9	9	9	9	9	9	9
16		9	9	9	9	9	9	9
17		9						
18	1	2	0	2				
19	1	2						
20	1	0	1	10	1			
21	1	1	538764005	0				
22	21	1	538764005	0				
23	1	20	0			1	1	0 0
24	1	20	0					
25	.00089							
26	1							
27	1	1	2	1	1	0		
28	20	20	21	1	0			
29	STDP							

CARD NUMBER	1	2	3	4	5	6	7	8
1	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890
2	C D L U M N N U M B E R							





CARD NUMBER	C O L U M N   N U M B E R							
	1	2	3	4	5	6	7	8
1	1-D CONSOLDN (FINITE STR, THIN LVR, D.DR) DATA: Davis&Raymond65 EQ:B2							
2	21	0	1	1	1	0.	10	
3	1	0						
4								
5								
6								
7								
8								
9	10							
10	.005995618							
11	1	0	0.	0.	0.	1		
12	21	0	0.	0.	.0231			
13	1	1						
14		.9	.9	.9	.9	.9	.9	.9
15		.9	.9	.9	.9	.9	.9	.9
16		.9	.9	.9	.9	.9	.9	.9
17		9	.9	.9	.9			
18	1	2	0	2				
19	1	2						
20	1	0.	1.	10.	1.			
21	1	1.538764005	0.					
22	21	1.538764005	0.					
23	1	20	1			2	1	44
24	1	1.						
25	.537.538764005		.54	.56	.58	.6	.62	.64
26	.66	.68	.7	.72	.74	.76	.78	.8
27	.82	.84	.86	.88	.9	.9011.36003E-31	.35692E-3	
28	1.35474E-31.32023E-31.28701E-31.25504E-31.22424E-31.19456E-31.16595E-31.13836E-3							
29	1.11173E-31.08603E-31.0612E-31.03722E-31.01404E-39.91636E-49.69961E-4							
30	9.28691E-49.09037E-4		8.9E-48.89064E-4					
31	1.							
32	1	1	2	1	1	0.		
33	20	20	21	1	0	0.		
34	STOP							
CARD NUMBER	C O L U M N   N U M B E R							
	1	2	3	4	5	6	7	8

CARD NUMBER	C O L U M N   N U M B E R							
	1	2	3	4	5	6	7	8
1	1-D CONSOLDN (FINITE STR, THIN LVR, D.DR) DATA: Davis&Raymond65 EQ:B3							
2	21	0	1	1	1	0.	10	
3	1	0						
4								
5								
6								
7								
8								
9	10							
10	.005995618							
11	1	0	0.	0.	0.	1		
12	21	0	0.	0.	.0231			
13	0	1						
14	1	2	0	2				
15	1	2						
16	0.		1.	10.	1.			
17	1	1	1.	0.				
18	21	1	1.	0.				
19	1	20	1			2	1	42
20	1	1.						
21	0.	.05	.1	.15	.2	.25	.3	.35
22	.4	.45	.5	.55	.6	.65	.7	.75
23	.8	.85	.9	.95	1.	8.9E-49.07165E-49.24832E-4		
24	9.4302E-4	9.6175E-49.81044E-41.00092E-31.02141E-31.04254E-31.06433E-31.08681E-3						
25	1.11001E-31.13395E-31.15869E-31.18424E-31.21064E-31.23794E-31.26617E-31.29538E-3							
26	1.32561E-31.35692E-3							
27	1.							
28	1	1	2	1	1	0.		
29	20	20	21	1	0	0.		
30	STOP							
CARD NUMBER	C O L U M N   N U M B E R							
	1	2	3	4	5	6	7	8





















A.15 Finite strain-thin layer: Poskitt (1969)

CARD NUMBER	1	2	3	4	5	6	7	8
1	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890
2	1-D CONSOLDN (FINITE STR, THIN LVR, D.DR) DATA POSKIT 69 EQ B10							
3	21	0	1	1	1	0	10	
4	1	0						
5								
6								
7								
8								
9	10							
10		.05						
11	1	0	0		0		1	
12	21	0	0		0	3.14159265		
13	0	1						
14	1	2	0	2				
15	1	2						
16		0		1	50		1	
17	1	1	1		0	0		
18	21	1	1		0			
19	1	20	1				2	1 42 0
20	1		1					
21		0 0	0 05	0 10	0 15	0 20	0 25	0 30 0 35
22		0 40	0 45	0 50	0 55	0 60	0 65	0 70 0 75
23		0 80	0 85	0 90	0 95	1 002	467401102	436756182 40651093
24	2	376654142	347182672	318091442	289375442	261029742	233049452	205429772 17816596
25	2	151253322	124687262	058463192	072576642	047023172	021798391	.996897981 .97231770
26	1	948053331	92410073					
27		1						
28	1		1	2	1	1	0	
29	20		20	21	1	0	0	
30	STOP							
CARD NUMBER	1	2	3	4	5	6	7	8
1	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890
C O L U M N N U M B E R								







A.17 Small strain-thick layer: Raymond (1969)

CARD NUMBER	1	2	3	4	5	6	7	8
1	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890
2	1-D CONSOLID (SMALL STR, THICK LVR, APPR: e-logp & e-logk, D.DR) DATA: Raymond69 EO: A4							
3	21	0	1	1	1	0	10	
4	1	0						
5								
6								
7								
8								
9								
10	10							
11	.038578							
12	1	1	0	0	0	0		
13	2	0	0	0	.05	1		
14	21	1	0	0	1			
15	1	1						
16		7	7	7	7	7	7	7
17		7	7	7	7	7	7	7
18		7	7	7	7	7	7	7
19	1	2	0	0				
20	1	2						
21	1	0	0	10	0			
22	1	20	1			4	20	42 42
23	1							
24	0.0	0.350	0.700	1.050	1.400	1.750	2.100	2.450
25	2.800	3.150	3.500	3.850	4.200	4.550	4.900	5.250
26	5.600	5.950	6.300	6.650	7.000	7.350	7.700	8.050
27	.020119712	.022276676	.024800060	.027777778	.031325867	.035598858	.040812162	.047258979
28	.055363322	.067746220	.079348335	.087656250	.123114805	.180000000	.216333153	.308641975
29	.475624257	.826446281						
30	0.0	0.350	0.700	1.050	1.400	1.750	2.100	2.450
31	2.800	3.150	3.500	3.850	4.200	4.550	4.900	5.250
32	5.600	5.950	6.300	6.650	7.000	7.350	7.700	8.050
33	.141843972	.149253731	.157480315	.166666667	.176981150	.188679245	.202020202	.217391304
34	.235294118	.256410256	.281690141	.312500000	.350877193	.400000000	.466116279	.556555556
35	.689655172	.809090909						
36	2							
37	0.0	0.350	0.700	1.050	1.400	1.750	2.100	2.450
38	2.800	3.150	3.500	3.850	4.200	4.550	4.900	5.250
39	5.600	5.950	6.300	6.650	7.000	7.350	7.700	8.050
40	.018024970	.021003991	.023308565	.026014668	.029220542	.033057851	.037703836	.043402778
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42	.357309458	.591715976						
43	0.0	0.350	0.700	1.050	1.400	1.750	2.100	2.450
44	2.800	3.150	3.500	3.850	4.200	4.550	4.900	5.250
45	5.600	5.950	6.300	6.650	7.000	7.350	7.700	8.050
46	.137931034	.144927636	.152671765	.161290323	.170940171	.181818182	.194174767	.208333333
47	.224719101	.243802439	.265688667	.294117647	.327868852	.370370370	.425631915	.500000000
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51	2.800	3.150	3.500	3.850	4.200	4.550	4.900	5.250
52	5.600	5.950	6.300	6.650	7.000	7.350	7.700	8.050
53	.018017206	.019837334	.021847874	.024414062	.027320538	.030778701	.034937549	.040000000
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55	.282184076	.444444444						
56	0.0	0.350	0.700	1.050	1.400	1.750	2.100	2.450
57	2.800	3.150	3.500	3.850	4.200	4.550	4.900	5.250
58	5.600	5.950	6.300	6.650	7.000	7.350	7.700	8.050
59	.134228188	.140845070	.148148148	.156250000	.165289256	.175438596	.188915888	.200000000
60	.215053763	.232558140	.253164567	.277777778	.307692308	.344827586	.392156863	.464546455
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63	0.0	0.350	0.700	1.050	1.400	1.750	2.100	2.450
64	2.800	3.150	3.500	3.850	4.200	4.550	4.900	5.250
65	5.600	5.950	6.300	6.650	7.000	7.350	7.700	8.050
66	.017087445	.018765247	.020702862	.022955841	.025600000	.028727377	.032464897	.036982248
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69	0.0	0.350	0.700	1.050	1.400	1.750	2.100	2.450
70	2.800	3.150	3.500	3.850	4.200	4.550	4.900	5.250
71	5.600	5.950	6.300	6.650	7.000	7.350	7.700	8.050
72	.130718954	.136988301	.143884892	.151515152	.160000000	.169491525	.180180180	.192307692
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75	5							
76	0.0	0.350	0.700	1.050	1.400	1.750	2.100	2.450
77	2.800	3.150	3.500	3.850	4.200	4.550	4.900	5.250
78	5.600	5.950	6.300	6.650	7.000	7.350	7.700	8.050
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83	2.800	3.150	3.500	3.850	4.200	4.550	4.900	5.250
84	5.600	5.950	6.300	6.650	7.000	7.350	7.700	8.050
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90	2.800	3.150	3.500	3.850	4.200	4.550	4.900	5.250
91	5.600	5.950	6.300	6.650	7.000	7.350	7.700	8.050
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239	2.800	3.150	3.500	3.850	4.200	4.550	4.900	5.250	
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244	18	1.							
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246	2.800	3.150	3.500	3.850	4.200	4.550	4.900	5.250	
247	5.600	5.950	6.300	6.650	7.000	.007561437	.008043596	.008573388	
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251	0.0	0.350	0.700	1.050	1.400	1.750	2.100	2.450	
252	2.800	3.150	3.500	3.850	4.200	4.550	4.900	5.250	
253	5.600	5.950	6.300	6.650	7.000	.086956522	.089686099	.092592593	
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256	.206185567	.222222222							
257	19	1.							
258	0.0	0.350	0.700	1.050	1.400	1.750	2.100	2.450	
259	2.800	3.150	3.500	3.850	4.200	4.550	4.900	5.250	
260	5.600	5.950	6.300	6.650	7.000	.007305136	.007762619	.008264463	
261	.008816593	.009425959	.010100755	.010850694	.011687363	.012624669	.013679423	.014872100	
262	.016227839	.017777778	.019560859	.021626298	.024037017	.026874496	.030245747	.034293553	
263	.039211842	.045269353							
264	0.0	0.350	0.700	1.050	1.400	1.750	2.100	2.450	
265	2.800	3.150	3.500	3.850	4.200	4.550	4.900	5.250	
266	5.600	5.950	6.300	6.650	7.000	.085470085	.088105727	.090909091	
267	.093896714	.097087379	.100502513	.104166667	.108108108	.112359551	.116959064	.121951220	
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269	.198019802	.212765957							
270	20	1.							
271	0.0	0.350	0.700	1.050	1.400	1.750	2.100	2.450	
272	2.800	3.150	3.500	3.850	4.200	4.550	4.900	5.250	
273	5.600	5.950	6.300	6.650	7.000	.007061648	.007496111	.007971939	
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277	0.0	0.350	0.700	1.050	1.400	1.750	2.100	2.450	
278	2.800	3.150	3.500	3.850	4.200	4.550	4.900	5.250	
279	5.600	5.950	6.300	6.650	7.000	.084033613	.086580087	.089285714	
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288	6	6	7	6	0				
289	7	7	8	7	0				
290	8	8	9	8	0				
291	9	9	10	9	0				
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297	15	15	16	15	0				
298	16	16	17	16	0				
299	17	17	18	17	0				
300	18	18	19	18	0				
301	19	19	20	19	0				
302	20	20	21	20	0				
303	STOP								
CARD	1	2	3	4	5	6	7	8	
NUMBER	123456789012345678901234567890123456789012345678901234567890								
	C O L U M N	N U M B E R							





A.18 Small strain-thick layer: Davis (1971)

CARD NUMBER	C O L U M N   N U M B E R							
	1	2	3	4	5	6	7	8
1	12345678901	2345678901	2345678901	2345678901	2345678901	2345678901	2345678901	2345678901
2	1-D CONSDLDN (SMALL STR, THICK LVR, APPR:CONST Cv & e-logp, D.DR) DATA: Davis73(Tve.2) EO:AA							
3	21	0	1	1	1	0.	10	
4	1	0						
5								
6								
7								
8								
9	10							
10		.125						
11	1	1	0.	0.	0.	0		
12	2	0	0.	0.	.25	1		
13	21	1	0.	0.	5.			
14	1	1						
15		.01	.01	.01	.01	.01	.01	
16		.01	.01	.01	.01	.01	.01	
17		.01	.01	.01	.01	.01	.01	
18		.01						
19	1	2	0	0				
20	1	2						
21	0.							
22	1	20	1		10.	0.		
23	1						4	20
24	0	0	0.00050	0.00100	0.00150	0.00200	0.00250	0.00300
25	0	0.00400	0.00450	0.00500	0.00550	0.00600	0.00650	0.00700
26	0	0.00800	0.00850	0.00900	0.00950	0.01000	442477876	442575781
27	.442771751	.442869796	.442967885	.443066017	.443164192	.443262411	.443360674	.443458980
28	.443557330	.443655723	.443754160	.443852641	.443951165	.444049734	.444148346	.444247001
29	.444345701	.444444444						
30	0	0	0.00050	0.00100	0.00150	0.00200	0.00250	0.00300
31	0	0.00400	0.00450	0.00500	0.00550	0.00600	0.00650	0.00700
32	0	0.00800	0.00850	0.00900	0.00950	0.01000	442477876	442575781
33	.442771751	.442869796	.442967885	.443066017	.443164192	.443262411	.443360674	.443458980
34	.443557330	.443655723	.443754160	.443852641	.443951165	.444049734	.444148346	.444247001
35	.444345701	.444444444						
36	2	1						
37	0	0	0.00050	0.00100	0.00150	0.00200	0.00250	0.00300
38	0	0.00400	0.00450	0.00500	0.00550	0.00600	0.00650	0.00700
39	0	0.00800	0.00850	0.00900	0.00950	0.01000	362318841	362384490
40	.362515860	.362581581	.362647325	.362713094	.362778886	.362844702	.362910543	.362976407
41	.363042294	.363108206	.363174142	.363240102	.363306085	.363372093	.363438125	.363504180
42	.363570260	.363636364						
43	0	0	0.00050	0.00100	0.00150	0.00200	0.00250	0.00300
44	0	0.00400	0.00450	0.00500	0.00550	0.00600	0.00650	0.00700
45	0	0.00800	0.00850	0.00900	0.00950	0.01000	362318841	362384490
46	.362515860	.362581581	.362647325	.362713094	.362778886	.362844702	.362910543	.362976407
47	.363042294	.363108206	.363174142	.363240102	.363306085	.363372093	.363438125	.363504180
48	.363570260	.363636364						
49	3	1						
50	0	0	0.00050	0.00100	0.00150	0.00200	0.00250	0.00300
51	0	0.00400	0.00450	0.00500	0.00550	0.00600	0.00650	0.00700
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55	.307644878	.307692308						
56	0	0	0.00050	0.00100	0.00150	0.00200	0.00250	0.00300
57	0	0.00400	0.00450	0.00500	0.00550	0.00600	0.00650	0.00700
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63	0	0	0.00050	0.00100	0.00150	0.00200	0.00250	0.00300
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299 17 17 18 17 0 0.
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301 19 19 20 19 0 0.
302 20 20 21 20 0 0.
303 STOP

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CARD 1 2 3 4 5 6 7 8
NUMBER 1234567890123456789012345678901234567890123456789012345678901234567890
C O L U M N N U M B E R

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CARD NUMBER	1								2								3								4								5								6								7								8							
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1	1-D CONSOLDN (Terzaghi's Analysis, D.DR) DATA: Davis71(Tv=.2)																																																															
2	21	1	0	1	1	0	10																																																									
3	1	0																																																														
4																																																																
5																																																																
6																																																																
7																																																																
8																																																																
9	10																																																															
10																																																																
11	1	1	0	0	0	0	0																																																									
12	2	0	0	0	.25	1																																																										
13	21	1	0	0	5																																																											
14	1	1																																																														
15		.01		.01	.01	.01	.01	.01	.01	.01																																																						
16		.01		.01	.01	.01	.01	.01	.01	.01																																																						
17		.01		.01	.01	.01	.01	.01	.01	.01																																																						
18		.01																																																														
19	1	2	0	0																																																												
20	1	2																																																														
21		0		0	10	0																																																										
22	1	20	0	0																																																												
23	1		1																																																													
24		1																																																														
25		1																																																														
26	1		1	2	1	1	0	0																																																								
27	20		20	21	1	0	0																																																									
28	STOP																																																															



## A.19 Small strain-thick layer: Viggiani (1973)

CARD NUMBER	1	2	3	4	5	6	7	8
1	12345678901234567890123456789012345678901234567890123456789012345678901234567890							
2	1-D CONSOLDN (SMALL STR, THICK LVR, APPR:CONST CV & e-logp, D.DR) DATA:Viggiani73(Tvs.2) EQ:A4							
3	21 0 1 1 1 0 10							
4	1 0							
5								
6								
7								
8								
9	10							
10	12.5							
11	1 1 0 0 0 0							
12	2 0 0 0 2.5 1							
13	21 1 0 0 50							
14	1 1							
15	10 10 10 10 10 10							
16	10 10 10 10 10 10							
17	10 10 10 10 10 10							
18	10 10 10 10 10 10							
19	1 2 0 0							
20	1 2							
21	1 0 0 500 0							
22	1 20 1 0 4 20 42 42							
23	1 1							
24	0 0 0 500 1 000 1 500 2 000 2 500 3 000 3 500							
25	4 000 4 500 5 000 5 500 6 000 6 500 7 000 7 500							
26	8 000 8 500 9 000 9 500 10 000 074074074 075923077 080000000							
27	083333333 088955522 090909091 095238085 100000000 105263158 111111111 117647059							
28	125000000 133333333 142857143 153846154 165836567 181818182 200000000 222222222							
29	250000000 285714285							
30	0 0 0 500 1 000 1 500 2 000 2 500 3 000 3 500							
31	4 000 4 500 5 000 5 500 6 000 6 500 7 000 7 500							
32	8 000 8 500 9 000 9 500 10 000 074074074 075923077 080000000							
33	083333333 088955522 090909091 095238085 100000000 105263158 111111111 117647059							
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35	250000000 285714285							
36	2 1							
37	0 0 0 500 1 000 1 500 2 000 2 500 3 000 3 500							
38	4 000 4 500 5 000 5 500 6 000 6 500 7 000 7 500							
39	8 000 8 500 9 000 9 500 10 000 054054054 055555555 057142857							
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62	4 1							
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97	8 000 8 500 9 000 9 500 10 000 025974026 026315789 026666667							





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104	8.000	8.500	9.000	9.500	10.000	.022988506	.023255814	.023529412
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108	0.0	0.500	1.000	1.500	2.000	2.500	3.000	3.500
109	4.000	4.500	5.000	5.500	6.000	6.500	7.000	7.500
110	8.000	8.500	9.000	9.500	10.000	.022988506	.023255814	.023529412
111	.023809524	.024096386	.024390244	.024691358	.025000000	.025316456	.025641026	.025974026
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116	4.000	4.500	5.000	5.500	6.000	6.500	7.000	7.500
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119	.023255814	.023529412	.023809524	.024096386	.024390244	.024691358	.025000000	.025316456
120	.025641026	.025974026						
121	0.0	0.500	1.000	1.500	2.000	2.500	3.000	3.500
122	4.000	4.500	5.000	5.500	6.000	6.500	7.000	7.500
123	8.000	8.500	9.000	9.500	10.000	.020618557	.020833333	.021052632
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126	.025641026	.025974026						
127	9	1						
128	0.0	0.500	1.000	1.500	2.000	2.500	3.000	3.500
129	4.000	4.500	5.000	5.500	6.000	6.500	7.000	7.500
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133	.022727273	.022988506						
134	0.0	0.500	1.000	1.500	2.000	2.500	3.000	3.500
135	4.000	4.500	5.000	5.500	6.000	6.500	7.000	7.500
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139	.022727273	.022988506						
140	10	1						
141	0.0	0.500	1.000	1.500	2.000	2.500	3.000	3.500
142	4.000	4.500	5.000	5.500	6.000	6.500	7.000	7.500
143	8.000	8.500	9.000	9.500	10.000	.017094017	.017241379	.017391304
144	.017543860	.017699115	.017857143	.018018018	.018181818	.018348624	.018518519	.018691589
145	.018867925	.019047619	.019230769	.019417476	.019607843	.019801980	.020000000	.020202020
146	.020408163	.020618557						
147	0.0	0.500	1.000	1.500	2.000	2.500	3.000	3.500
148	4.000	4.500	5.000	5.500	6.000	6.500	7.000	7.500
149	8.000	8.500	9.000	9.500	10.000	.017094017	.017241379	.017391304
150	.017543860	.017699115	.017857143	.018018018	.018181818	.018348624	.018518519	.018691589
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152	.020408163	.020618557						
153	11	1						
154	0.0	0.500	1.000	1.500	2.000	2.500	3.000	3.500
155	4.000	4.500	5.000	5.500	6.000	6.500	7.000	7.500
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159	.018518519	.018691589						
160	0.0	0.500	1.000	1.500	2.000	2.500	3.000	3.500
161	4.000	4.500	5.000	5.500	6.000	6.500	7.000	7.500
162	8.000	8.500	9.000	9.500	10.000	.015748031	.015873016	.016000000
163	.016129032	.016260163	.016393443	.016528926	.016666667	.016806723	.016949153	.017094017
164	.017241379	.017391304	.017543860	.017699115	.017857143	.018018018	.018181818	.018348624
165	.018518519	.018691589						
166	12	1						
167	0.0	0.500	1.000	1.500	2.000	2.500	3.000	3.500
168	4.000	4.500	5.000	5.500	6.000	6.500	7.000	7.500
169	8.000	8.500	9.000	9.500	10.000	.014598540	.014705882	.014814815
170	.014925373	.015037594	.015151515	.015267176	.015384615	.015503876	.015625000	.015748031
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172	.016949153	.017094017						
173	0.0	0.500	1.000	1.500	2.000	2.500	3.000	3.500
174	4.000	4.500	5.000	5.500	6.000	6.500	7.000	7.500
175	8.000	8.500	9.000	9.500	10.000	.014598540	.014705882	.014814815
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178	.016949153	.017094017						
179	13	1						
180	0.0	0.500	1.000	1.500	2.000	2.500	3.000	3.500
181	4.000	4.500	5.000	5.500	6.000	6.500	7.000	7.500
182	8.000	8.500	9.000	9.500	10.000	.013605442	.013698630	.013793103
183	.013888889	.013986014	.014084507	.014184397	.014285714	.014388489	.014492754	.014598540
184	.014705882	.014814815	.014925373	.015037594	.015151515	.015267176	.015384615	.015503876
185	.015625000	.015748031						
186	0.0	0.500	1.000	1.500	2.000	2.500	3.000	3.500
187	4.000	4.500	5.000	5.500	6.000	6.500	7.000	7.500
188	8.000	8.500	9.000	9.500	10.000	.013605442	.013698630	.013793103
189	.013888889	.013986014	.014084507	.014184397	.014285714	.014388489	.014492754	.014598540
190	.014705882	.014814815	.014925373	.015037594	.015151515	.015267176	.015384615	.015503876
191	.015625000	.015748031						
192	14	1						
193	0.0	0.500	1.000	1.500	2.000	2.500	3.000	3.500
194	4.000	4.500	5.000	5.500	6.000	6.500	7.000	7.500
195	8.000	8.500	9.000	9.500	10.000	.012738854	.012820513	.012903226
196	.012987013	.013071895	.013157895	.013245033	.013333333	.013422819	.013513514	.013605442
197	.013698630	.013793103	.013888889	.013986014	.014084507	.014184397	.014285714	.014388489
198	.014492754	.014598540						
199	0.0	0.500	1.000	1.500	2.000	2.500	3.000	3.500
200	4.000	4.500	5.000	5.500	6.000	6.500	7.000	7.500
201	8.000	8.500	9.000	9.500	10.000	.012738854	.012820513	.012903226
202	.012987013	.013071895	.013157895	.013245033	.013333333	.013422819	.013513514	.013605442
203	.013698630	.013793103	.013888889	.013986014	.014084507	.014184397	.014285714	.014388489





204	.014492754	.014598540							
205	15	1.							
206	0.0	0.500	1.000	1.500	2.000	2.500	3.000	3.500	
207	4.000	4.500	5.000	5.500	6.000	6.500	7.000	7.500	
208	8.000	8.500	9.000	9.500	10.000	.011976048	.012048193	.012121212	
209	.012195122	.012269939	.012345679	.012422360	.012500000	.012578616	.012658228	.012738854	
210	.012820513	.012903226	.012987013	.013071895	.013157895	.013245033	.013333333	.013422819	
211	.013513514	.013605442							
212	0.0	0.500	1.000	1.500	2.000	2.500	3.000	3.500	
213	4.000	4.500	5.000	5.500	6.000	6.500	7.000	7.500	
214	8.000	8.500	9.000	9.500	10.000	.011976048	.012048193	.012121212	
215	.012195122	.012269939	.012345679	.012422360	.012500000	.012578616	.012658228	.012738854	
216	.012820513	.012903226	.012987013	.013071895	.013157895	.013245033	.013333333	.013422819	
217	.013513514	.013605442							
218	16	1.							
219	0.0	0.500	1.000	1.500	2.000	2.500	3.000	3.500	
220	4.000	4.500	5.000	5.500	6.000	6.500	7.000	7.500	
221	8.000	8.500	9.000	9.500	10.000	.011299435	.011363636	.011428571	
222	.011494253	.011560694	.011627907	.011695906	.011764706	.011834320	.011904762	.011976048	
223	.012048193	.012121212	.012195122	.012269939	.012345679	.012422360	.012500000	.012578616	
224	.012658228	.012738854							
225	0.0	0.500	1.000	1.500	2.000	2.500	3.000	3.500	
226	4.000	4.500	5.000	5.500	6.000	6.500	7.000	7.500	
227	8.000	8.500	9.000	9.500	10.000	.011299435	.011363636	.011428571	
228	.011494253	.011560694	.011627907	.011695906	.011764706	.011834320	.011904762	.011976048	
229	.012048193	.012121212	.012195122	.012269939	.012345679	.012422360	.012500000	.012578616	
230	.012658228	.012738854							
231	17	1.							
232	0.0	0.500	1.000	1.500	2.000	2.500	3.000	3.500	
233	4.000	4.500	5.000	5.500	6.000	6.500	7.000	7.500	
234	8.000	8.500	9.000	9.500	10.000	.010695187	.010752688	.010810811	
235	.010869565	.010928962	.010989011	.011049724	.011111111	.011173184	.011235955	.011299435	
236	.011363636	.011428571	.011494253	.011560694	.011627907	.011695906	.011764706	.011834320	
237	.011904762	.011976048							
238	0.0	0.500	1.000	1.500	2.000	2.500	3.000	3.500	
239	4.000	4.500	5.000	5.500	6.000	6.500	7.000	7.500	
240	8.000	8.500	9.000	9.500	10.000	.010695187	.010752688	.010810811	
241	.010869565	.010928962	.010989011	.011049724	.011111111	.011173184	.011235955	.011299435	
242	.011363636	.011428571	.011494253	.011560694	.011627907	.011695906	.011764706	.011834320	
243	.011904762	.011976048							
244	18	1							
245	0.0	0.500	1.000	1.500	2.000	2.500	3.000	3.500	
246	4.000	4.500	5.000	5.500	6.000	6.500	7.000	7.500	
247	8.000	8.500	9.000	9.500	10.000	.010152284	.010204082	.010256410	
248	.010309278	.010362694	.010416667	.010471204	.010526316	.010582011	.010638298	.010695187	
249	.010752688	.010810811	.010869565	.010928962	.010989011	.011049724	.011111111	.011173184	
250	.011235955	.011299435							
251	0.0	0.500	1.000	1.500	2.000	2.500	3.000	3.500	
252	4.000	4.500	5.000	5.500	6.000	6.500	7.000	7.500	
253	8.000	8.500	9.000	9.500	10.000	.010152284	.010204082	.010256410	
254	.010309278	.010362694	.010416667	.010471204	.010526316	.010582011	.010638298	.010695187	
255	.010752688	.010810811	.010869565	.010928962	.010989011	.011049724	.011111111	.011173184	
256	.011235955	.011299435							
257	19	1.							
258	0.0	0.500	1.000	1.500	2.000	2.500	3.000	3.500	
259	4.000	4.500	5.000	5.500	6.000	6.500	7.000	7.500	
260	8.000	8.500	9.000	9.500	10.000	.009661836	.009708738	.009756098	
261	.009803922	.009852217	.009900990	.009950249	.010000000	.010050251	.010101010	.010152284	
262	.010204082	.010256410	.010309278	.010362694	.010416667	.010471204	.010526316	.010582011	
263	.010638298	.010695187							
264	0.0	0.500	1.000	1.500	2.000	2.500	3.000	3.500	
265	4.000	4.500	5.000	5.500	6.000	6.500	7.000	7.500	
266	8.000	8.500	9.000	9.500	10.000	.009661836	.009708738	.009756098	
267	.009803922	.009852217	.009900990	.009950249	.010000000	.010050251	.010101010	.010152284	
268	.010204082	.010256410	.010309278	.010362694	.010416667	.010471204	.010526316	.010582011	
269	.010638298	.010695187							
270	20	1.							
271	0.0	0.500	1.000	1.500	2.000	2.500	3.000	3.500	
272	4.000	4.500	5.000	5.500	6.000	6.500	7.000	7.500	
273	8.000	8.500	9.000	9.500	10.000	.009216590	.009259259	.009302326	
274	.009345794	.009389671	.009433962	.009478673	.009523810	.009569378	.009615385	.009661836	
275	.009708738	.009756098	.009803922	.009852217	.009900990	.009950249	.010000000	.010050251	
276	.010101010	.010152284							
277	0.0	0.500	1.000	1.500	2.000	2.500	3.000	3.500	
278	4.000	4.500	5.000	5.500	6.000	6.500	7.000	7.500	
279	8.000	8.500	9.000	9.500	10.000	.009216590	.009259259	.009302326	
280	.009345794	.009389671	.009433962	.009478673	.009523810	.009569378	.009615385	.009661836	
281	.009708738	.009756098	.009803922	.009852217	.009900990	.009950249	.010000000	.010050251	
282	.010101010	.010152284							
283	1	2	1	0	0.				
284	2	3	2	0	0.				
285	3	4	3	0	0.				
286	4	5	4	0	0.				
287	5	6	5	0	0.				
288	6	7	6	0	0.				
289	7	8	7	0	0.				
290	8	9	8	0	0.				
291	9	10	9	0	0.				
292	10	11	10	0	0.				
293	11	12	11	0	0.				
294	12	13	12	0	0.				
295	13	14	13	0	0.				
296	14	15	14	0	0.				
297	15	16	15	0	0.				
298	16	17	16	0	0.				
299	17	18	17	0	0.				
300	18	19	18	0	0.				
301	19	20	19	0	0.				
302	20	20	20	0	0.				
303	STOP								
CARD	1	2	3	4	5	6	7	8	
NUMBER	1234567890123456789012345678901234567890123456789012345678901234567890								
	C O L U M N			N U M B E R					



CARD NUMBER	C O L U M N										N U M B E R									
	1		2		3		4		5		6		7		8					
	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890				
1	I-D CONSOLDN (Terzaghi's Analysis, D.DR) DATA:Viggiani73(Tv=.2)																			
2	21	1	0	1	1	0.	10													
3	1	0																		
4																				
5																				
6																				
7																				
8																				
9	10																			
10	31.25																			
11	1	1	0.		0.		0.	0												
12	2	0	0.		0.		2.5	1												
13	21	1	0.		0.		50.													
14	1	1																		
15		10.		10.		10.		10.		10.		10.		10.		10.				
16		10.		10.		10.		10.		10.		10.		10.		10.				
17		10.		10.		10.		10.		10.		10.		10.		10.				
18		10.		10.		10.		10.		10.		10.		10.		10.				
19	1	2	0	0																
20	1	2																		
21		0.		0.		500.		0.												
22	1	20	0									1	1	0	0					
23	1		1.																	
24		1.																		
25		1																		
26	1		1	2	1	1		0.												
27	20		20	21	1	0		0.												
28	STOP																			



# A.20 Thick layer: logarithmic material law and constant coefficient of consolidation

CARD NUMBER	1	2	3	4	5	6	7	8
1	1234567890123456789012345678901234567890123456789012345678901234567890							
2	1-D CONSOLDN (SMALL STR, THICK LVR, APPR CONST CV & E-100P, D DR, EQ A4							
3	1	0	1	1	0	10		
4								
5								
6								
7								
8								
9								
10	10							
11	711974814							
12	1	1	0	0	0	0		
13	21	1	0	0	5	1		
14	1	1			10			
15		10		10	10		10	
16		10		10	10		10	
17		10		10	10		10	
18		10		10	10		10	
19	1	2	0	0				
20	1	2						
21		0		100	0			
22	1	20	1				4	20
23	1		1				42	42
24	0	0	0	5	1	0	1	5
25	4	0	4	5	5	0	2	5
26	8	0	8	5	9	0	6	0
27	011707634	012407048	013195336	014090590	015116164	016302748	010521404	011082868
28	021324531	023764626	026835301	030817261	036186851	043822464	055542145	075818773
29	119412322	280950595						
30	0	0	0	5	1	0	1	5
31	4	0	4	5	5	0	2	5
32	8	0	8	5	9	0	6	0
33	011707634	012407048	013195336	014090590	015116164	016302748	010521404	011082868
34	021324531	023764626	026835301	030817261	036186851	043822464	055542145	075818773
35	119412322	280950595						
36	2	1						
37	0	0	0	5	1	0	1	5
38	4	0	4	5	5	0	2	5
39	8	0	8	5	9	0	6	0
40	012186754	012865873	013625147	014479658	015448524	016556345	017835325	019328451
41	021094419	023215535	025810911	029059626	033243902	038835849	046689472	058524692
42	078397523	118705562						
43	0	0	0	5	1	0	1	5
44	4	0	4	5	5	0	2	5
45	8	0	8	5	9	0	6	0
46	012186754	012865873	013625147	014479658	015448524	016556345	017835325	019328451
47	021094419	023215535	025810911	029059626	033243902	038835849	046689472	058524692
48	078397523	118705562						
49	3	1						
50	0	0	0	5	1	0	1	5
51	4	0	4	5	5	0	2	5
52	8	0	8	5	9	0	6	0
53	012119738	012749352	013447965	014227580	015103150	016093553	017222965	018522861
54	020034992	021815960	023944450	026533179	029749517	033853182	039270125	046750861
55	057752357	075525048						
56	0	0	0	5	1	0	1	5
57	4	0	4	5	5	0	2	5
58	8	0	8	5	9	0	6	0
59	012119738	012749352	013447965	014227580	015103150	016093553	017222965	018522861
60	020034992	021815960	023944450	026533179	029749517	033853182	039270125	046750861
61	057752357	075525048						
62	4	1						
63	0	0	0	5	1	0	1	5
64	4	0	4	5	5	0	2	5
65	8	0	8	5	9	0	6	0
66	011863045	012438273	013072129	013774056	014555643	015431265	016418980	017541763
67	018829423	020321076	022069396	024146867	026656102	029747310	033649517	038730070
68	045617620	055484729						
69	0	0	0	5	1	0	1	5
70	4	0	4	5	5	0	2	5
71	8	0	8	5	9	0	6	0
72	011863045	012438273	013072129	013774056	014555643	015431265	016418980	017541763
73	018829423	020321076	022069396	024146867	026656102	029747310	033649517	038730070
74	045617620	055484729						
75	5	1						
76	0	0	0	5	1	0	1	5
77	4	0	4	5	5	0	2	5
78	8	0	8	5	9	0	6	0
79	011525954	012048635	012620973	013250397	013945897	014718452	015581622	016552343
80	017652048	018908277	020357007	022046156	024040993	026432746	029352865	032998567
81	037678139	043904264						
82	0	0	0	5	1	0	1	5
83	4	0	4	5	5	0	2	5
84	8	0	8	5	9	0	6	0
85	011525954	012048635	012620973	013250397	013945897	014718452	015581622	016552343
86	017652048	018908277	020357007	022046156	024040993	026432746	029352865	032998567
87	037678139	043904264						
88	6	1						
89	0	0	0	5	1	0	1	5
90	4	0	4	5	5	0	2	5
91	8	0	8	5	9	0	6	0
92	011156020	011630204	012146487	012710737	013329964	014012615	014768852	015611605
93	016556234	017622536	018835643	020228114	021842903	023737873	025992872	028721273
94	032089633	036353022						
95	0	0	0	5	1	0	1	5
96	4	0	4	5	5	0	2	5
97	8	0	8	5	9	0	6	0





98	.011156020	.011630204	.012146487	.012710737	.013329964	.014012615	.014768958	.015611609
99	.016556234	.017622536	.018835643	.020228114	.021842903	.023737873	.025992872	.028721273
100	.032089633	.036353028						
101	7	1						
102	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
103	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
104	8.0	8.5	9.0	9.5	10.0	.009663960	.010008577	.010378681
105	.010777208	.011207563	.011673717	.012180332	.012732913	.013338015	.014003498	.014738876
106	.015555769	.016468527	.017495077	.018658112	.019986791	.021519215	.023306140	.025416706
107	.027947595	.031038254						
108	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
109	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
110	8.0	8.5	9.0	9.5	10.0	.009663960	.010008577	.010378681
111	.010777208	.011207563	.011673717	.012180332	.012732913	.013338015	.014003498	.014738876
112	.015555769	.016468527	.017495077	.018658112	.019986791	.021519215	.023306140	.025416706
113	.027947595	.031038254						
114	8	1						
115	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
116	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
117	8.0	8.5	9.0	9.5	10.0	.009382159	.009699174	.010038362
118	.010402133	.010793260	.011214950	.011670929	.012165559	.012703971	.013292246	.013937649
119	.014648925	.015436701	.016314022	.017297074	.018406196	.019667302	.021113929	.022790266
120	.024755744	.027092231						
121	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
122	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
123	8.0	8.5	9.0	9.5	10.0	.009382159	.009699174	.010038362
124	.010402133	.010793260	.011214950	.011670929	.012165559	.012703971	.013292246	.013937649
125	.014648925	.015436701	.016314022	.017297074	.018406196	.019667302	.021113929	.022790266
126	.024755744	.027092231						
127	9	1						
128	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
129	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
130	8.0	8.5	9.0	9.5	10.0	.009101981	.009393885	.009705133
131	.010037712	.010393894	.010776284	.011187885	.011632177	.012113215	.012635755	.013205411
132	.013828854	.014514082	.015270757	.016110668	.017048349	.018101925	.019294301	.020654836
133	.022221805	.024046045						
134	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
135	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
136	8.0	8.5	9.0	9.5	10.0	.009101981	.009393885	.009705133
137	.010037712	.010393894	.010776284	.011187885	.011632177	.012113215	.012635755	.013205411
138	.013828854	.014514082	.015270757	.016110668	.017048349	.018101925	.019294301	.020654836
139	.022221805	.024046045						
140	10	1						
141	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
142	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
143	8.0	8.5	9.0	9.5	10.0	.008827683	.009096834	.009382913
144	.009687570	.010012675	.010360358	.010733055	.011133568	.011565130	.012031498	.012537059
145	.013086971	.013687338	.014345437	.015070016	.015871685	.016763438	.017761362	.018885620
146	.020161822	.021623004						
147	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
148	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
149	8.0	8.5	9.0	9.5	10.0	.008827683	.009096834	.009382913
150	.009687570	.010012675	.010360358	.010733055	.011133568	.011565130	.012031498	.012537059
151	.013086971	.013687338	.014345437	.015070016	.015871685	.016763438	.017761362	.018885620
152	.020161822	.021623004						
153	11	1						
154	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
155	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
156	8.0	8.5	9.0	9.5	10.0	.008561706	.008810274	.009073707
157	.009353379	.009650840	.009967842	.010306377	.010668715	.011057459	.011475604	.011926617
158	.012414531	.012944069	.013520794	.014151309	.014843504	.015606899	.016453073	.017396262
159	.018454167	.019649069						
160	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
161	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
162	8.0	8.5	9.0	9.5	10.0	.008561706	.008810274	.009073707
163	.009353379	.009650840	.009967842	.010306377	.010668715	.011057459	.011475604	.011926617
164	.012414531	.012944069	.013520794	.014151309	.014843504	.015606899	.016453073	.017396262
165	.018454167	.019649069						
166	12	1						
167	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
168	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
169	8.0	8.5	9.0	9.5	10.0	.008305447	.008535407	.008778464
170	.009035770	.009308615	.009598451	.009906916	.010235866	.010587410	.010963961	.011368284
171	.011803569	.012273516	.012782435	.013335384	.013938335	.014598393	.015324073	.016125674
172	.017015767	.018009862						
173	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
174	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
175	8.0	8.5	9.0	9.5	10.0	.008305447	.008535407	.008778464
176	.009035770	.009308615	.009598451	.009906916	.010235866	.010587410	.010963961	.011368284
177	.011803569	.012273516	.012782435	.013335384	.013938335	.014598393	.015324073	.016125674
178	.017015767	.018009862						
179	13	1						
180	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
181	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
182	8.0	8.5	9.0	9.5	10.0	.008059583	.008272714	.008497424
183	.008734683	.008985571	.009251298	.009533220	.009832865	.010151958	.010492455	.010856586
184	.011246899	.011666324	.012118244	.012606586	.013135939	.013711697	.014340239	.015029175
185	.015787648	.016626744						
186	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
187	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
188	8.0	8.5	9.0	9.5	10.0	.008059583	.008272714	.008497424
189	.008734683	.008985571	.009251298	.009533220	.009832865	.010151958	.010492455	.010856586
190	.011246899	.011666324	.012118244	.012606586	.013135939	.013711697	.014340239	.015029175
191	.015787648	.016626744						
192	14	1						
193	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
194	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
195	8.0	8.5	9.0	9.5	10.0	.007824270	.008022163	.008230326
196	.008449579	.008680833	.008925103	.009183517	.009457341	.009747997	.010057084	.010386415
197	.010738044	.011114316	.011517915	.011951932	.012419938	.012926089	.013475248	.014073139
198	.014726549	.015443589						
199	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
200	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
201	8.0	8.5	9.0	9.5	10.0	.007824270	.008022163	.008230326
202	.008449579	.008680833	.008925103	.009183517	.009457341	.009747997	.010057084	.010386415
203	.010738044	.011114316	.011517915	.011951932	.012419938	.012926089	.013475248	.014073139





204	.014726549	.015443589							
205	15	1							
206	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5	
207	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5	
208	8.0	8.5	9.0	9.5	10.0	.007599559	.007783641	.007976861	
209	.008179919	.008393585	.008618713	.008856249	.009107250	.009372894	.009654500	.009953552	
210	.010271723	.010610906	.010973255	.011361225	.011777636	.012225733	.012709275	.013232641	
211	.013800963	.014420293							
212	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5	
213	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5	
214	8.0	8.5	9.0	9.5	10.0	.007599559	.007783641	.007976861	
215	.008179919	.008393585	.008618713	.008856249	.009107250	.009372894	.009654500	.009953552	
216	.010271723	.010610906	.010973255	.011361225	.011777636	.012225733	.012709275	.013232641	
217	.013800963	.014420293							
218	16	1							
219	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5	
220	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5	
221	8.0	8.5	9.0	9.5	10.0	.007385087	.007556625	.007736321	
222	.007924772	.008122633	.008330627	.008549553	.008780296	.009023840	.009281280	.009553840	
223	.009842893	.010149982	.010476851	.010825472	.011198093	.011597281	.012025981	.012487591	
224	.012986054	.013525964							
225	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5	
226	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5	
227	8.0	8.5	9.0	9.5	10.0	.007385087	.007556625	.007736321	
228	.007924772	.008122633	.008330627	.008549553	.008780296	.009023840	.009281280	.009553840	
229	.009842893	.010149982	.010476851	.010825472	.011198093	.011597281	.012025981	.012487591	
230	.012986054	.013525964							
231	17	1							
232	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5	
233	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5	
234	8.0	8.5	9.0	9.5	10.0	.007180628	.007340762	.007508201	
235	.007683455	.007867087	.008059711	.008262005	.008474715	.008698667	.008934777	.009184061	
236	.009447656	.009726828	.010023002	.010337779	.010672968	.011030621	.011413076	.011823004	
237	.012263477	.012738039							
238	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5	
239	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5	
240	8.0	8.5	9.0	9.5	10.0	.007180628	.007340762	.007508201	
241	.007683455	.007867087	.008059711	.008262005	.008474715	.008698667	.008934777	.009184061	
242	.009447656	.009726828	.010023002	.010337779	.010672968	.011030621	.011413076	.011823004	
243	.012263477	.012738039							
244	18	1							
245	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5	
246	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5	
247	8.0	8.5	9.0	9.5	10.0	.006985365	.007135089	.007291371	
248	.007454654	.007625416	.007804186	.007991539	.008188109	.008394592	.008611759	.008840461	
249	.009081641	.009336350	.009605758	.009891177	.010194076	.010516113	.010859160	.011225343	
250	.011617084	.012037156							
251	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5	
252	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5	
253	8.0	8.5	9.0	9.5	10.0	.006985365	.007135089	.007291371	
254	.007454654	.007625416	.007804186	.007991539	.008188109	.008394592	.008611759	.008840461	
255	.009081641	.009336350	.009605758	.009891177	.010194076	.010516113	.010859160	.011225343	
256	.011617084	.012037156							
257	19	1							
258	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5	
259	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5	
260	8.0	8.5	9.0	9.5	10.0	.006799811	.006940064	.007086224	
261	.007238674	.007397827	.007564136	.007733895	.007920243	.008111173	.008311536	.008522048	
262	.008743501	.008976771	.009222828	.009482755	.009757757	.010049187	.010358560	.010687587	
263	.011028201	.011412601							
264	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5	
265	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5	
266	8.0	8.5	9.0	9.5	10.0	.006799811	.006940064	.007086224	
267	.007238674	.007397827	.007564136	.007733895	.007920243	.008111173	.008311536	.008522048	
268	.008743501	.008976771	.009222828	.009482755	.009757757	.010049187	.010358560	.010687587	
269	.011028201	.011412601							
270	20	1							
271	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5	
272	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5	
273	8.0	8.5	9.0	9.5	10.0	.006622605	.006754183	.006891097	
274	.007033675	.007182279	.007337297	.007499154	.007668313	.007845280	.008030608	.008224904	
275	.008428835	.008643135	.008868617	.009106178	.009356817	.009621643	.009901897	.010198967	
276	.010514413	.010849995							
277	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5	
278	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5	
279	8.0	8.5	9.0	9.5	10.0	.006622605	.006754183	.006891097	
280	.007033675	.007182279	.007337297	.007499154	.007668313	.007845280	.008030608	.008224904	
281	.008428835	.008643135	.008868617	.009106178	.009356817	.009621643	.009901897	.010198967	
282	.010514413	.010849995							
283	1	1	2	1	0				
284	2	2	3	2	0				
285	3	3	4	3	0				
286	4	4	5	4	0				
287	5	5	6	5	0				
288	6	6	7	6	0				
289	7	7	8	7	0				
290	8	8	9	8	0				
291	9	9	10	9	0				
292	10	10	11	10	0				
293	11	11	12	11	0				
294	12	12	13	12	0				
295	13	13	14	13	0				
296	14	14	15	14	0				
297	15	15	16	15	0				
298	16	16	17	16	0				
299	17	17	18	17	0				
300	18	18	19	18	0				
301	19	19	20	19	0				
302	20	20	21	20	0				
303	STOP								
CARD	1	2	3	4	5	6	7	8	
NUMBER	1234567890123456789012345678901234567890123456789012345678901234567890								
	C O L U M N	N U M B E R							



CARD NUMBER	C O L U M N   N U M B E R							
	1	2	3	4	5	6	7	8
1	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890	1234567890
2	1-D CONSOLDN (FINITE STR, THICK LVR, APPR:CONST CV & e-logp, D DR) EQ:A8							
3	21	0	1	1	1	0	10	
4	1	0						
5								
6								
7								
8								
9	10							
10	.711974814							
11	1	1	0	0	0	0		
12	2	0	0	0	.5	1		
13	21	1	0	0	10			
14	1	1						
15		10		10	10	10	10	
16		10		10	10	10	10	
17		10		10	10	10	10	
18		10		10	10	10	10	
19	1	2	0	0				
20	1	2						
21		0		50	0			
22	1	20	1					
23	1		1			4	20	42 42
24	0	0	0.5	1.0	1.5	2.0	2.5	3.0 3.5
25	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
26	8.0	8.5	9.0	9.5	10.0	0.23429339	0.24234334	0.25114347
27	.026081020	.027148607	.028334762	.029661636	.031157432	.032858664	.034813518	.037087013
28	.039769207	.042988782	.046936706	.051909988	.058399038	.067280215	.080301304	.101569944
29	.143835636	.280950595						
30	0	0	0.5	1.0	1.5	2.0	2.5	3.0 3.5
31	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
32	8.0	8.5	9.0	9.5	10.0	0.10014084	0.10521404	0.11082868
33	.011707634	.012407048	.013195336	.014090590	.015116164	.016302748	.017691490	.019338882
34	.021324531	.023764626	.026835301	.030817261	.036186851	.043822464	.055542145	.075818773
35	.119412322	.280950595						
36	2	1						
37	0	0	0.5	1.0	1.5	2.0	2.5	3.0 3.5
38	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
39	8.0	8.5	9.0	9.5	10.0	0.20280151	0.20934269	0.21645637
40	.022422574	.023275138	.024215618	.025259194	.026424861	.027736730	.029225900	.030933241
41	.032913637	.035242667	.038027586	.041426225	.045681482	.051188992	.058642944	.069392308
42	.085482511	.118705562						
43	0	0	0.5	1.0	1.5	2.0	2.5	3.0 3.5
44	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
45	8.0	8.5	9.0	9.5	10.0	0.10520756	0.11023061	0.11575735
46	.012186754	.012865872	.013625147	.014479652	.015443524	.016556345	.017835325	.019328451
47	.021094419	.023215535	.025810911	.029059626	.033243902	.038835649	.046689472	.058524692
48	.078397523	.118705562						
49	3	1						
50	0	0	0.5	1.0	1.5	2.0	2.5	3.0 3.5
51	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
52	8.0	8.5	9.0	9.5	10.0	0.18346067	0.18900821	0.19501130
53	.020153184	.020864383	.021643649	.022501845	.023452345	.024511816	.025701317	.027047872
54	.028586784	.030365135	.032447235	.034923464	.037925256	.041651926	.046422071	.052780900
55	.061751505	.075525048						
56	0	0	0.5	1.0	1.5	2.0	2.5	3.0 3.5
57	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
58	8.0	8.5	9.0	9.5	10.0	0.10555868	0.11030299	0.11549384
59	.012119738	.012749352	.013447965	.014227580	.015103150	.016093553	.017222965	.018522861
60	.020034993	.021815960	.023944450	.026533179	.029749517	.033853182	.039270125	.046750861
61	.057752357	.075525048						
62	4	1						
63	0	0	0.5	1.0	1.5	2.0	2.5	3.0 3.5
64	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
65	8.0	8.5	9.0	9.5	10.0	0.16885246	0.17364118	0.17879910
66	.018437319	.019041901	.019700270	.020420378	.021211868	.022086565	.023059131	.024147898
67	.025376671	.026775650	.028385279	.030260149	.032476125	.035142050	.038420241	.042564715
68	.047998278	.055484729						
69	0	0	0.5	1.0	1.5	2.0	2.5	3.0 3.5
70	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
71	8.0	8.5	9.0	9.5	10.0	0.10417696	0.10858690	0.11338670
72	.011863045	.012438273	.013072129	.013774056	.014555643	.015431265	.016418980	.017541783
73	.018829423	.020321076	.022069396	.024146867	.026656102	.029747310	.033649517	.038730070
74	.045617620	.055484729						
75	5	1						
76	0	0	0.5	1.0	1.5	2.0	2.5	3.0 3.5
77	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
78	8.0	8.5	9.0	9.5	10.0	0.15696034	0.16114017	0.16562298
79	.017044494	.017564837	.018128312	.018740841	.019409509	.020142880	.020951398	.021847956
80	.022848659	.023973915	.025250004	.026711388	.028404228	.030391914	.032764088	.035652105
81	.039257053	.043904264						
82	0	0	0.5	1.0	1.5	2.0	2.5	3.0 3.5
83	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
84	8.0	8.5	9.0	9.5	10.0	0.10198672	0.10605778	0.11046737
85	.011525954	.012048635	.012620973	.013250397	.013945897	.014718453	.015581623	.016552343
86	.017652048	.018908277	.020357007	.022046158	.024040993	.026432746	.029352969	.032998567
87	.037678139	.043904264						
88	6	1						
89	0	0	0.5	1.0	1.5	2.0	2.5	3.0 3.5
90	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
91	8.0	8.5	9.0	9.5	10.0	0.14689882	0.15057744	0.15450719
92	.015871625	.016323727	.016810835	.017337423	.017908786	.018531234	.019212357	.019961372
93	.020789590	.021711056	.022743445	.023909340	.025238086	.026768584	.028533599	.030666664
94	.033213632	.036353028						
95	0	0	0.5	1.0	1.5	2.0	2.5	3.0 3.5
96	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
97	8.0	8.5	9.0	9.5	10.0	0.09940184	0.10314906	0.10718988



98	.011156020.	.011630204.	.012146487.	.012710737.	.013329964.	.014012615.	.014768958.	.015611609
99	.016556234.	.017622536.	.018835643.	.020228114.	.021842903.	.023737873.	.025992872.	.028721273
100	.032089633.	.036353028						
101	7	1.						
102	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
103	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
104	8.0	8.5	9.0	9.5	10.0	.013818975.	.014144928.	.014491879
105	.014862042.	.015257969.	.015682607.	.016139383.	.016632309.	.017166113.	.017746409.	.018379915
106	.019074752.	.019840827.	.020690366.	.021638647.	.022705034.	.023914461.	.025299636.	.026904377
107	.028788871.	.031038254						
108	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
109	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
110	8.0	8.5	9.0	9.5	10.0	.009663960.	.010008577.	.010378681
111	.010777208.	.011207563.	.011673717.	.012180332.	.012732913.	.013338015.	.014003498.	.014738876
112	.015555769.	.016468527.	.017495077.	.018658112.	.019986791.	.021519215.	.023306140.	.025416706
113	.027947595.	.031038254						
114	8	1.						
115	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
116	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
117	8.0	8.5	9.0	9.5	10.0	.013053142.	.013343640.	.013651828
118	.013979467.	.014328569.	.014701440.	.015100738.	.015529543.	.015991449.	.016490674.	.017032205
119	.017621983.	.018267150.	.018976374.	.019760282.	.020632060.	.021608284.	.022710112.	.023965003
120	.025409304.	.027092231						
121	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
122	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
123	8.0	8.5	9.0	9.5	10.0	.009382159.	.009699174.	.010038362
124	.010402133.	.010793260.	.011214950.	.011670929.	.012165559.	.012703971.	.013292246.	.013937649
125	.014648925.	.015436701.	.016314022.	.017297074.	.018406196.	.019667302.	.021113929.	.022790266
126	.024755744.	.027092231						
127	9	1.						
128	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
129	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
130	8.0	8.5	9.0	9.5	10.0	.012372143.	.012632376.	.012907616
131	.013199277.	.013508961.	.013838491.	.014189952.	.014565737.	.014968616.	.015401802.	.015869055
132	.016374802.	.016924291.	.017523801.	.018180903.	.018904827.	.019706943.	.020601433.	.021606234
133	.022744390.	.024046045						
134	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
135	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
136	8.0	8.5	9.0	9.5	10.0	.009101981.	.009393885.	.009705133
137	.010037712.	.010393894	.010776284.	.011187885.	.011632177.	.012113215.	.012635755.	.013205411
138	.013828854.	.014514082.	.015270757.	.016110668.	.017048349.	.018101925.	.019294301.	.020654836
139	.022221805.	.024046045						
140	10	1.						
141	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
142	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
143	8.0	8.5	9.0	9.5	10.0	.011761303.	.011995505.	.012242521
144	.012503495.	.012779713.	.013072632.	.013383903.	.013715409.	.014069306.	.014448079.	.014854603
145	.015292226.	.015764874.	.016277182.	.016834663.	.017443932.	.018112999.	.018851664.	.019672059
146	.020589390.	.021623004						
147	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
148	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
149	8.0	8.5	9.0	9.5	10.0	.008827683.	.009096834.	.009388293
150	.009687570.	.010012675.	.010360358.	.010733055.	.011133568.	.011565130.	.012031498.	.012537059
151	.013086971.	.013687338.	.014345437.	.015070016.	.015871685.	.016763438.	.017761362.	.018885620
152	.020161822.	.021623004						
153	11	1.						
154	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
155	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
156	8.0	8.5	9.0	9.5	10.0	.011209459.	.011421120.	.011643788
157	.011878395.	.012125984.	.012387728.	.012664947.	.012959136.	.013271997.	.013605471.	.013961785
158	.014343510.	.014753627.	.015195618.	.015673571.	.016192329.	.016757667.	.017376538.	.018057389
159	.018810599.	.019649069						
160	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
161	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
162	8.0	8.5	9.0	9.5	10.0	.008561706.	.008810274.	.009073707
163	.009353379.	.009650840.	.009967842.	.010306377.	.010668715.	.011057459.	.011475604.	.011926617
164	.012414531.	.012944069.	.013520794.	.014151309.	.014843504.	.015606899.	.016453073.	.017396262
165	.018454167.	.019649069						
166	12	1.						
167	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
168	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
169	8.0	8.5	9.0	9.5	10.0	.010708026.	.010900057.	.011101594
170	.011313403.	.011536337.	.011771347.	.012019499.	.012281992.	.012560181.	.012855600.	.013169996
171	.013505367.	.013864011.	.014248580.	.014662162.	.015108365.	.015591440.	.016116431.	.016689375
172	.017317557.	.018009862						
173	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
174	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
175	8.0	8.5	9.0	9.5	10.0	.008305447.	.008535407.	.008778464
176	.009035770.	.009330815.	.009598451.	.009906916.	.010235866.	.010587410.	.010963961.	.011368284
177	.011803569.	.012273516.	.012782435.	.013335384.	.013938335.	.014598393.	.015324073.	.016125674
178	.017015767.	.018009862						
179	13	1.						
180	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
181	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
182	8.0	8.5	9.0	9.5	10.0	.010250120.	.010424963.	.010608059
183	.010800042.	.011001612.	.011213549.	.011436720.	.011672095.	.011920762.	.012183948.	.012463036
184	.012759599.	.013075429.	.013412578.	.013773409.	.014180657.	.014577508.	.015027696.	.015515628
185	.016046548.	.016626744						
186	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
187	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
188	8.0	8.5	9.0	9.5	10.0	.008059583.	.008272714.	.008497424
189	.008734683.	.008985571.	.009251298.	.009533220.	.009832865.	.010151958.	.010492455.	.010856586
190	.011246899.	.011666324.	.012118244.	.012606586.	.013135939.	.013711697.	.014340239.	.015029175
191	.015787648.	.016626744						
192	14	1.						
193	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
194	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
195	8.0	8.5	9.0	9.5	10.0	.009830009.	.009989726.	.010156642
196	.010331286.	.010514236.	.010706135.	.010907696.	.011119708.	.011343055.	.011578724.	.011827824
197	.012091604.	.012371479.	.012669056.	.012986169.	.013324921.	.013687738.	.014077433.	.014497287
198	.014951153.	.015443589						
199	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
200	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
201	8.0	8.5	9.0	9.5	10.0	.007824270.	.008022163.	.008230326
202	.008449579.	.008680833.	.008925103.	.009183517.	.009457341.	.009747997.	.010057084.	.010386415
203	.010738044.	.011114316.	.011517915.	.011951932.	.012419938.	.012926089.	.013475248.	.014073139





CARD NUMBER	1	2	3	4	5	6	7	8
204	014726549	015443589						
205	15	1						
206	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
207	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
208	8.0	8.5	9.0	9.5	10.0	009443260	009589611	009742272
209	009901682	010066328	010242740	010425507	010617278	010818777	011030805	011254258
210	011490140	011739579	012003847	012284385	012582834	012901067	013241235	013605825
211	013997720	014420293						
212	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
213	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
214	8.0	8.5	9.0	9.5	10.0	007599559	007783641	007976861
215	008179919	008393585	008618713	008856249	009107250	009372894	009654500	009953552
216	010271723	010610906	010973255	011361225	011777636	012225733	012709275	013232641
217	013800963	014420293						
218	16	1						
219	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
220	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
221	8.0	8.5	9.0	9.5	10.0	009085792	009220277	009360313
222	009506276	009658566	009817631	009983958	010158087	010340610	010532187	010733547
223	010945503	011168963	011404946	011654598	011919210	012200249	012499385	012818528
224	013159875	013525964						
225	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
226	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
227	8.0	8.5	9.0	9.5	10.0	007385087	007556625	007736321
228	007924772	008122633	008330627	008549553	008780296	009023840	009281280	009553840
229	009842893	010148982	010476851	010825472	011198093	011597281	012025981	012487591
230	012986054	013525964						
231	17	1						
232	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
233	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
234	8.0	8.5	9.0	9.5	10.0	008754547	008878463	009007286
235	0095141328	009280932	009426472	009578357	009737034	009902998	010076790	010259011
236	010450325	010651470	010863268	011086639	011322613	011572353	011837168	012118549
237	012418191	012738039						
238	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
239	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
240	8.0	8.5	9.0	9.5	10.0	007180629	007340762	007508201
241	007683455	007867087	008059711	008262005	008474715	008698667	008934777	009184061
242	009447656	009726828	010023002	010337779	010672968	011030621	011413076	011823004
243	012263477	012738039						
244	18	1						
245	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5
246	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5
247	8.0	8.5	9.0	9.5	10.0	008446092	008560542	008679342
248	008802759	008931082	009064628	009203740	009348795	009500203	009658416	0





CARD NUMBER	C O L U M N   N U M B E R									
	1	2	3	4	5	6	7	8		
1	1234567890123456789012345678901234567890123456789012345678901234567890									
2	1-D CONSOLDN (FINITE STR, THICK LVR, APPR:CONST Cv & e-logp, D.DR) EQ:A9									
3	1	0	1	1	1	0	10			
4										
5										
6										
7										
8										
9	10									
10	.711974814									
11	1	1	0.	0.	0.	0				
12	2	0	0.	0.	.421893	1				
13	21	1	0.	0.	8.43786					
14	1	1								
15		10.		10.	10.	10.	10.	10.		
16		10.		10.	10.	10.	10.	10.		
17		10.		10.	10.	10.	10.	10.		
18		10.								
19	1	2	0	0						
20	1	2								
21		0.	0.	100.	0.					
22	1	20	1				4	20	42	42
23	1		1.							
24	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5		
25	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5		
26	8.0	8.5	9.0	9.5	10.0	0.13885633	0.14356058	0.14869713		
27	.015433210	.016054619	.016743904	.017513527	.018379285	.019361539	.020487019	.021791567		
28	.023324469	.025155535	.027387266	.030176958	.033779899	.038641859	.045623208	.056646258		
29	.077160580	.132153627								
30	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5		
31	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5		
32	8.0	8.5	9.0	9.5	10.0	0.16543294	0.17370393	0.18284547		
33	.019300265	.020435468	.021712557	.023159907	.024813998	.026722532	.028949114	.031580469		
34	.034738012	.038597106	.043420782	.049622336	.057890525	.069464934	.086824238	.115750256		
35	.173579216	.346881719								
36	2	1.								
37	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5		
38	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5		
39	8.0	8.5	9.0	9.5	10.0	0.13208298	0.13620444	0.14067500		
40	.014554367	.015086935	.015672351	.016319378	.017038887	.017844543	.018753785	.019789240		
41	.020980856	.022369172	.024010558	.025985937	.028416085	.031490200	.035523605	.041086938		
42	.049339102	.063082846								
43	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5		
44	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5		
45	8.0	8.5	9.0	9.5	10.0	0.15369998	0.16081415	0.16861884		
46	.017721974	.018674523	.019735287	.020923817	.022264675	.023789153	.025537739	.027563774		
47	.029938981	.032762137	.036173158	.040376997	.045686414	.052603577	.061989045	.075450906		
48	.096381641	.133383333								
49	3	1.								
50	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5		
51	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5		
52	8.0	8.5	9.0	9.5	10.0	0.12610329	0.12974820	0.13367941		
53	.013793407	.014255624	.014759857	.015312459	.015921170	.016595519	.017347392	.018191814		
54	.019148079	.020241414	.021505493	.022986361	.024748813	.026887225	.029544992	.032951897		
55	.037502450	.043940322								
56	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5		
57	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5		
58	8.0	8.5	9.0	9.5	10.0	0.14352108	0.14970520	0.15644626		
59	.016382302	.017192986	.018088082	.019081497	.020190372	.021436077	.022845605	.024453547		
60	.026304968	.028459703	.030998938	.034035673	.037731995	.042328982	.048201496	.055965933		
61	.066712113	.082565815								
62	#	1.								
63	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5		
64	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5		
65	8.0	8.5	9.0	9.5	10.0	0.12077988	0.12402970	0.12751748		
66	.013127193	.013532669	.013972147	.014450351	.014972948	.015546799	.016180298	.016883833		
67	.017870423	.018556630	.019563864	.020720340	.022064064	.023647572	.025545746	.027869384		
68	.030790160	.034590168								
69	0.0	0.5	1.0	1.5	2.0	2.5	3.0	3.5		
70	4.0	4.5	5.0	5.5	6.0	6.5	7.0	7.5		
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## A.21 Thick layer: Gibson et al (1981)

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93	.120006E-3	.135313E-3	.151512E-3	.168332E-3	.185405E-3	.202268E-3	.218380E-3	.233139E-3
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CARD 1 2 3 4 5 6 7 8
NUMBER 123456789012345678901234567890123456789012345678901234567890
C O L U M N N U M B E R

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7	21	1	0.	0.	2.728			
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13	10.	10.	10.	10.	10.	10.	10.	10.
14	10.	10.	10.	10.	10.	10.	10.	10.
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16	10.	10.	10.	10.	10.	10.	10.	10.
17	10.	10.	10.	10.	10.	10.	10.	10.
18	10.	10.	10.	10.	10.	10.	10.	10.
19	1	2	0	0				
20	1	2						
21	0.	0.	5000.	0.				
22	1	20	1					
23	1	1				4	20	42
24	0.0	0.50	1.00	1.50	2.00	2.50	3.00	3.50
25	4.00	4.50	5.00	5.50	6.00	6.50	7.00	7.50
26	8.00	8.50	9.00	9.50	10.00	7.93902E-5	9.26644E-5	1.07945E-4
27	.125456E-4	.145419E-4	.168038E-4	.193481E-4	.221861E-4	.253205E-4	.287425E-4	.324279E-4
28	.363344E-4	.403980E-4	.445319E-4	.486253E-4	.525511E-4	.561620E-4	.593090E-4	.618481E-4
29	.636542E-4	.646333E-4						
30	0.0	0.50	1.00	1.50	2.00	2.50	3.00	3.50
31	4.00	4.50	5.00	5.50	6.00	6.50	7.00	7.50
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36	2	1.						
37	0.0	0.50	1.00	1.50	2.00	2.50	3.00	3.50
38	4.00	4.50	5.00	5.50	6.00	6.50	7.00	7.50
39	8.00	8.50	9.00	9.50	10.00	.737277E-5	.861270E-5	1.00427E-4
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41	.344560E-4	.384545E-4	.425674E-4	.466956E-4	.507182E-4	.544967E-4	.578825E-4	.607272E-4
42	.628951E-4	.642761E-4						
43	0.0	0.50	1.00	1.50	2.00	2.50	3.00	3.50
44	4.00	4.50	5.00	5.50	6.00	6.50	7.00	7.50
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50	0.0	0.50	1.00	1.50	2.00	2.50	3.00	3.50
51	4.00	4.50	5.00	5.50	6.00	6.50	7.00	7.50
52	8.00	8.50	9.00	9.50	10.00	.684448E-5	.800178E-5	.933883E-5
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56	0.0	0.50	1.00	1.50	2.00	2.50	3.00	3.50
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67	.308353E-4	.346546E-4	.386608E-4	.427771E-4	.469030E-4	.509167E-4	.546789E-4	.580408E-4
68	.608543E-4	.629848E-4						
69	0.0	0.50	1.00	1.50	2.00	2.50	3.00	3.50
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71	8.00	8.50	9.00	9.50	10.00	.021316927	.025166191	.029710528
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78	8.00	8.50	9.00	9.50	10.00	.589310E-5	.689912E-5	.806501E-5
79	.941174E-5	.109614E-4	.127364E-4	.147587E-4	.170485E-4	.196222E-4	.224902E-4	.256544E-4
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81	.595962E-4	.620663E-4						
82	0.0	0.50	1.00	1.50	2.00	2.50	3.00	3.50
83	4.00	4.50	5.00	5.50	6.00	6.50	7.00	7.50
84	8.00	8.50	9.00	9.50	10.00	.019701718	.023259320	.027459328
85	.032417744	.038271516	.045182322	.053341035	.062972991	.074344219	.087768785	.103617467
86	.122327994	.144417140	.170494990	.201281800	.237627878	.280537082	.331194533	.390999357
87	.461603323	.544956467						
88	6	1.						
89	0.0	0.50	1.00	1.50	2.00	2.50	3.00	3.50
90	4.00	4.50	5.00	5.50	6.00	6.50	7.00	7.50
91	8.00	8.50	9.00	9.50	10.00	.546584E-5	.640289E-5	.749036E-5
92	.874856E-5	.101991E-4	.118643E-4	.137666E-4	.159271E-4	.183643E-4	.210918E-4	.241159E-4
93	.274324E-4	.310233E-4	.348536E-4	.388674E-4	.429868E-4	.471101E-4	.511145E-4	.548601E-4
94	.581977E-4	.609797E-4						
95	0.0	0.50	1.00	1.50	2.00	2.50	3.00	3.50
96	4.00	4.50	5.00	5.50	6.00	6.50	7.00	7.50
97	8.00	8.50	9.00	9.50	10.00	.018208898	.021496935	.025378703





```

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100 .426627114.503664495
101 7 1.
102 0.0 0.50 1.00 1.50 2.00 2.50 3.00 3.50
103 4.00 4.50 5.00 5.50 6.00 6.50 7.00 7.50
104 8.00 8.50 9.00 9.50 10.00.506823E-5.594053E-5.695416E-5
105 812870E-5.948517E-5.110457E-4.128327E-4.148681E-4.171719E-4.197603E-4.226434E-4
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107 .566738E-4.597374E-4
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109 4.00 4.50 5.00 5.50 6.00 6.50 7.00 7.50
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135 4.00 4.50 5.00 5.50 6.00 6.50 7.00 7.50
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161 4.00 4.50 5.00 5.50 6.00 6.50 7.00 7.50
162 8.00 8.50 9.00 9.50 10.00.012279517.014496868.017114612
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165 .287704123.339655749
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168 4.00 4.50 5.00 5.50 6.00 6.50 7.00 7.50
169 8.00 8.50 9.00 9.50 10.00.346270E-5.406805E-5.477509E-5
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181 4.00 4.50 5.00 5.50 6.00 6.50 7.00 7.50
182 8.00 8.50 9.00 9.50 10.00.320691E-5.376891E-5.442586E-5
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185 .457862E-4.498436E-4
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187 4.00 4.50 5.00 5.50 6.00 6.50 7.00 7.50
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192 14 1.
193 0.0 0.50 1.00 1.50 2.00 2.50 3.00 3.50
194 4.00 4.50 5.00 5.50 6.00 6.50 7.00 7.50
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200 4.00 4.50 5.00 5.50 6.00 6.50 7.00 7.50
201 8.00 8.50 9.00 9.50 10.00.009694376.011444920.013511564
202 .015951390.018831781.022232295.026246849.030986324.036581621.043187277.050985738
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```





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206	0.0	0.50	1.00	1.50	2.00	2.50	3.00	3.50
207	4.00	4.50	5.00	5.50	6.00	6.50	7.00	7.50
208	8.00	8.50	9.00	9.50	10.00	274937E-5	3.23330E-5	3.79979E-5
209	.446192E-5.523445E-5.613389E-5.717849E-5.838815E-5.978416E-5.113887E-4.132243E-4							
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213	4.00	4.50	5.00	5.50	6.00	6.50	7.00	7.50
214	8.00	8.50	9.00	9.50	10.00	008959822	010577725	012487778
215	.014742735.017404876.020547728.024258095.028638455.033809790.039914928.047122491							
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218	16 1.							
219	0.0	0.50	1.00	1.50	2.00	2.50	3.00	3.50
220	4.00	4.50	5.00	5.50	6.00	6.50	7.00	7.50
221	8.00	8.50	9.00	9.50	10.00	254518E-5	2.99404E-5	3.51981E-5
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224	.399044E-4.440353E-4							
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230	.194018757.229053326							
231	17 1.							
232	0.0	0.50	1.00	1.50	2.00	2.50	3.00	3.50
233	4.00	4.50	5.00	5.50	6.00	6.50	7.00	7.50
234	8.00	8.50	9.00	9.50	10.00	235587E-5	2.77209E-5	3.25991E-5
235	.383092E-5.449827E-5.527681E-5.618315E-5.723563E-5.845420E-5.986023E-5.114760E-4							
236	.133239E-4.154255E-4.178001E-4.204624E-4.234208E-4.266736E-4.302062E-4.339875E-4							
237	.379666E-4.420705E-4							
238	0.0	0.50	1.00	1.50	2.00	2.50	3.00	3.50
239	4.00	4.50	5.00	5.50	6.00	6.50	7.00	7.50
240	8.00	8.50	9.00	9.50	10.00	007653471	009035482	010667047
241	.012593229.014867227.017551849.020721241.024462939.028880288.034095291.040251984							
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243	.179317735.211697695							
244	18 1.							
245	0.0	0.50	1.00	1.50	2.00	2.50	3.00	3.50
246	4.00	4.50	5.00	5.50	6.00	6.50	7.00	7.50
247	8.00	8.50	9.00	9.50	10.00	218039E-5	2.56625E-5	3.01873E-5
248	.354871E-5.416858E-5.489239E-5.573588E-5.671660E-5.785374E-5.916806E-5.106815E-4							
249	.124164E-4.143950E-4.166378E-4.191621E-4.219795E-4.250935E-4.284960E-4.321643E-4							
250	.360572E-4.401125E-4							
251	0.0	0.50	1.00	1.50	2.00	2.50	3.00	3.50
252	4.00	4.50	5.00	5.50	6.00	6.50	7.00	7.50
253	8.00	8.50	9.00	9.50	10.00	007073559	008350853	009858793
254	.011639026.013740721.016221925.019151169.022609355.026691996.031511852.037202045							
255	.043919734.05185045							



CARD NUMBER	1								2								3								4								5								6								7								8							
	1	2	3	4	5	6	7	8	1	2	3	4	5	6	7	8	1	2	3	4	5	6	7	8	1	2	3	4	5	6	7	8	1	2	3	4	5	6	7	8	1	2	3	4	5	6	7	8	1	2	3	4	5	6	7	8								
1	1-D CONSOLDN (Terzaghi's Analysis, D.DR) DATA:Gibson-eta181																																																															
2	21	1							0	1	1						0.								0.								0.								0																							
3	1	0																																																														
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13									2	0							0.				0.				0.				.5				1																															
14									21	1							0.				0.				0.				10.																																			
15									1	1																																																						
16																																																																



## APPENDIX B

### Computer Programs



## LIST OF PROGRAMS

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Note: Most of these computer programs are written for the illustrated examples in this text. Modifications of these programs are necessary for their general use.





## B.1 Terzaghi's classical analysis

```

C *****
C *
C *      *** TERZAGHI25 ***
C *      This program is to compute the excess pore pressure
C *      at various depth and time according to Terzaghi's
C *      classical theory of consolidation
C *
C *****
C
C DOUBLE PRECISION ZH(25),TV(15),UE(15,25),AM,AMT,AMZ
C DOUBLE PRECISION DSIN,DEXP,DLOAT
C TV(1)=0.05D0
C DO 8 IJ=2,4
C   8 TV(IJ)=TV(IJ-1)+0.05D0
C DO 10 II=5,11
C   10 TV(II)=TV(II-1)+0.1D0
C   ZH(1)=0.05D0
C   DO 15 JJ=2,20
C     15 ZH(JJ)=ZH(JJ-1)+0.05D0
C   DO 20 J=1,11
C     DO 20 J=1,20
C       UE(I,J)=0.0D0
C     DO 20 K=1,50
C       KM=K-1
C       AM=(2.0D0+DLOAT(KM)+1.0D0)*3.141592854D0/2.0D0
C       AMZ=AM*ZH(J)
C       AMT=-1.0D0*TV(I)+AM**2
C       UE(I,J)=UE(I,J)+DSIN(AMZ)*DEXP(AMT)*2.0D0/AM
C   20 CONTINUE
C   WRITE(6,50)
C   WRITE(6,51)(TV(I),I=1,11)
C   DO 25 L=1,20
C     25 WRITE(6,52)ZH(L),(UE(I,L),I=1,11)
C 60 FORMAT('0'/' Depth',44X,'Time Factor Tv')
C 51 FORMAT(' Ratio',11F9.2//)
C 52 FORMAT('F5.2,2X,11F9.6)
C STOP
C END

```

## B.2 Gray's two-layered analysis

```

C *****
C *
C *      *** Gray45 ***
C *      This program is to find the analytical solution of
C *      the excess pore water pressure for two-layered soils
C *      according to Gray H. (1945) for certain values of sigma,
C *      mu, lambda as given below
C *
C *****
C
C DOUBLE PRECISION AN(10),TV(5),ZH(25),UE(5,25),CN,PARC,PARMA,TNPP
C DOUBLE PRECISION PARAZ,AMU,SIGMA,ALAMDA
C AN(1)=0.0714326347
C AN(2)=0.215
C AN(3)=0.3604
C AN(4)=0.5079
C AN(5)=0.6571
C AN(6)=0.8076
C AN(7)=0.9593
C AN(8)=1.1117
C SIGMA=2.0
C ALAMDA=4.0
C AMU=5.0
C TV(1)=20.0
C TV(2)=100.0
C TV(3)=200.0
C ZH(1)=5.0
C DO 10 II=2,21
C   10 ZH(II)=ZH(II-1)-0.25
C   DO 20 J=1,3
C     DO 20 J=1,21
C       UE(I,J)=0.0
C     DO 20 K=1,8
C       PARC=AMU+ALAMDA*AN(K)
C       CN=2.0*DCOS(AN(K))/(SIGMA*AN(K)+DSIN(PARC)**2+PARC*DCOS(AN(K))**2)
C       TNPP=-1.0*TV(I)+AN(K)*AN(K)
C       IF(J GE 18) GO TO 16
C       PARMA=AMU+AN(K)*(1.0+ALAMDA-ZH(J))
C       UE(I,J)=UE(I,J)+CN*DEXP(TNPP)*DCOS(AN(K))*DSIN(PARMA)
C     GO TO 20
C   16 PARAZ=AN(K)*ZH(J)
C     UE(I,J)=UE(I,J)+CN*DEXP(TNPP)*DCOS(PARAZ)*DSIN(PARC)
C 20 CONTINUE
C   WRITE(6,53)
C   WRITE(6,50)
C   WRITE(6,51)(TV(I),I=1,3)
C   DO 25 L=1,21
C     25 WRITE(6,52)ZH(L),(UE(I,L),I=1,3)
C STOP
C 60 FORMAT(' Depth',8X,'Time Factor Tv')
C 51 FORMAT(' Ratio',3F9.0//)
C 52 FORMAT('F5.2,2X,3F9.6)
C 53 FORMAT('0'/'4X,'Analytical Solution by Gray'//)
C END

```



## B.3 Martins' closed-form solution

```

C *****
C *
C *          *** MARTINS65 ***
C *
C * # One dimensional consolidation with linear spatial
C *   variation of permeability #
C *
C * This program is to calculate the excess pore pressure
C *   at various depth at time t according to Martins'
C *   analytical solution in 1965. The soil layer is assumed
C *   to drain freely at both top and bottom.
C *
C * Input variables:
C *   T = Time required, sec
C *   K0 = The coefficient of permeability, cm/sec
C *   MV = The coefficient of volume compressibility, cm2/g
C *   BETA = A dimensionless constant relates to the variation
C *     of permeability by the eqn:  $K = K_0(1+BETA*Z)$ 
C *   GAMMA = Unit weight of soil, g/cm3
C *   GAMMAW = Unit weight of water, g/cm3
C *   H2 = Total thickness of soil, cm
C *   UO = Initial uniform pore pressure, g/cm2
C *   NOEL = Number of equal segments of the soil profile
C *     specified for the outputs
C *   AN(n) = Roots of the eqn.  $J_0(A) \cdot Y_0(KA) - J_0(KA) \cdot Y_0(A) = 0$ 
C *     where k is the variable LAMBDA in the program
C *   NORT = Total number of the above roots used
C *
C * Use of the program
C *   1 First to compile the program by:
C *   RUN *FORTGTEST SCARDS=MARTINS65 SPRINT=*1 PAR=ID,SOURCE
C *   2 Repeat the following for subsequent executions
C *   DEBUG -LOAD*+*IMSLDPLIB 5=datafile 6=output
C *
C * Remarks:
C *   The routines MMBSJO and MMBSYN in the IMSL Library will
C *   be used. Their descriptions are shown at the back of this
C *   program
C *****
C
C IMPLICIT REAL*8(A-H,D-Z)
C REAL*8 KO,MV,LAMBDA,DSORT,DEXP,DFLOAT,MMBSJO
C DIMENSION AN(10),ALA(10),BJOLA(10),UT(10),Z(100),ZH(100),UE(100)
C
C *****
C *
C * Read all input data
C *
C *****
C
C READ(5,50)T,KO,MV,BETA,H2
C READ(5,51)NOEL,NORT,UO,GAMMA,GAMMAW
C READ(5,52)(AN(I),I=1,NORT)
C
C *****
C *
C * Calculation of the average consolidation at the given
C *   time T
C *
C *****
C
C LAMBDA=DSORT(1,DO+BETA)
C CH=BETA*BETA*KO/(4 DO+H2*H2*GAMMAW*MV)
C ACT=0 DO
C
C DO 15 I=1,NORT
C   ALA(I)=LAMBDA*AN(I)
C   BJOA=MMBSJO(ALA(I),1EJ1)
C   BJOLA(I)=MMBSJO(ALA(I),1EJ2)
C   EU=DEXP(-(1 DO+CH)*AN(I)*AN(I))*T/(1BJOA+BJOLA(I))
C   UT(I)=BJOA*EU
C   ACT=ACT+(BJOA-BJOLA(I))*EU/(AN(I)*AN(I))
C 15 CONTINUE
C
C AVECN=4 DO*ACT/BETA
C
C *****
C *
C * Compute the excess pore pressure at various depth
C *
C *****
C
C P1=3 14159265358979323846264338327950288419716939937511D0
C NONOD=NOEL+1
C DIV=H2/DFLOAT(NOEL)
C Z(1)=0 DO
C
C DO 20 J=1,NONOD
C   Z(J+1)=Z(J)+DIV
C   ZH(J)=Z(J)/H2
C   R=DSORT(1 DO+BETA*ZH(J))
C   UET=0 DO
C
C DO 18 L=1,NORT
C   RA=R*AN(L)
C   BJORA=MMBSJO(RA,1EJ3)
C   CALL MMBSYN(RA,0 DO,1,YORA,1EY1)
C   CALL MMBSYN(ALA(L),0 DO,1,YOLA,1EY2)
C   UET=UET+UT(L)*(1BJORA*YOLA-BJOLA(L)*YORA)
C 18 CONTINUE

```



```

C      UE(J)=UO*PI*UET
C      20 CONTINUE
C
C      *****
C      *
C      *   Output quantities
C      *
C      *****
C
C      WRITE(6,60)
C      WRITE(6,61)AVCDN
C      WRITE(6,62)KO,BETA,MV,H2,GAMMA,GAMMAW,UO,LAMBDA,CHI,NORT,NDEL
C      WRITE(6,63)T,(Z(I),ZH(I),I,UE(I),I=1,NONOD)
C      WRITE(6,64)IEJ1,IEJ2,IEJ3,IEY1,IEY2
C
C      STOP
C
C      *****
C      *
C      *   Format statments
C      *
C      *****
C
C      50 FORMAT(5G20.0)
C      51 FORMAT(2I5,3G20.0)
C      52 FORMAT(3F20.0)
C
C      60 FORMAT('1'////6X'The One Dimensional Consolidation with ',
C      1'Linear Spatial'/18X'Variation of Permeability'
C      1//13X'Analytical Solution by Martins (1965)'//)
C      61 FORMAT(4X'Average Consolidation =',F12.8//)
C      62 FORMAT(4X'KO      =',E15.6/4X'Beta    =',F15.5/
C      14X'Mv      =',E15.6/4X'2H      =',F15.5/
C      14X'GAMMA   =',F15.5/4X'GAMMAW =',F15.5/
C      14X'UO      =',F15.5/4X'Lambda =',F15.5/
C      14X'Chi     =',E15.7/4X'No of roots used =',I5/
C      14X'No of elements =',I5//)
C      63 FORMAT(13X'Excess pore pressure vs depth: '//
C      113X'Depth',6X'Depth',4X'NDE',9X'Time t'/24X'Ratio',10XG20.9
C      1//9XF11.5,F10.4,3X13,4XE15.8))
C      64 FORMAT(///5X'Error message from the *IMSLDPLIB for the final'
C      1,' iteration: '//5X'IEJ1 IEJ2 IEJ3 IEY1 IEY2'/3X5I6)
C
C      *****
C      *
C      *   # IMSL Library #
C      *
C      *   The following is a brief description of the Fortran
C      *   routines used in this program. This information is
C      *   available in the IMSL Library reference manual, edition
C      *   8, 1980. The routines with single precision are stored
C      *   in the public file *IMSLLIB, those with double precision
C      *   are stored in *IMSLDPLIB.
C      *
C      *
C      *   Routine name:  MMBSJO
C      *
C      *   Usage:  FUNCTION MMBSJO(ARG,IER)
C      *
C      *   Description of arguments:
C      *   ARG - Input argument
C      *   IER - Error parameter, terminal error
C      *
C      *
C      *   Routine name:  MMBSYN
C      *
C      *   Usage:  CALL MMBSYN(ARG,ORDER,N,YN,IER)
C      *
C      *   Description of arguments:
C      *   ARG - Input argument
C      *   ORDER - Desired order, which must be greater than or
C      *   equal to zero and less than one
C      *   N - Number of results required, with which order
C      *   start from 'ORDER' to 'ORDER+1' to 'ORDER+2', etc
C      *   YN - Output vector of above results
C      *   IER - Error parameter where
C      *   '129': ORDER is out of range or ARG is less than
C      *   the minimum allowable value
C      *   '130': N is less than or equal to zero
C      *   '131': YN overflow
C      *   '132': ARG is greater than the maximum allowable
C      *   value
C      *
C      *****
C
C      END

```



## B.4 Edelman's closed-form solution

```

C *****
C
C      *** EDELMANN53 ***
C
C      # One dimensional consolidation with the modulus of
C      elasticity proportional to the depth #
C
C      This program is to compute the excess pore pressure
C      at various depth at time t according to Edelman's
C      analytical solution in 1953. The soil layer is assumed
C      to drain at the top only.
C
C      Input variables:
C      T = Time required, sec
C      SIGMA0 = A small initial effective stress, g/cm**2
C      C = A dimensionless constant, relates to compressibility
C      by the eqn.  $MV=1/(C \cdot \text{GAMMA}+2 \cdot \text{SIGMA0})$ 
C      where Z is the distance in cm from the surface
C      GAMMA = Unit weight of soil, g/cm**3
C      GAMMAW = Unit weight of water, g/cm**3
C      H = Soil thickness, cm
C      K = The coefficient of permeability, cm/sec
C      UO = Initial uniform pore pressure, g/cm**2
C      NOEL = Number of equal segments of the soil profile
C      specified for the outputs
C      VDN(n) = Roots of eqn:  $J_0(V) \cdot Y_1(kV) - Y_0(V) \cdot J_1(kV) = 0$ 
C      where k is the variable LAMBDA in the program
C      NORT = Total number of the above roots used
C
C      Use of the program:
C      1 First to compile the program by:
C      RUN *FORTGTEST SCARDS=EDELMANN53 SPRINT=-1 PAR=ID,SOURCE
C      2 Repeat the following for subsequent executions:
C      DEBUG -LOAD***IMSLDPLIB 5=datafile 6=output
C
C      Remarks:
C      The % consolidation computed is in terms of settlement
C      The routines MMBSJO, MMBSJ1 and MMBSYN in the IMSL
C      Library will be used Their descriptions are shown at the
C      back of this program
C *****
C
C      IMPLICIT REAL*8(A-H,D-Z)
C      REAL*8 LAMBDA,K,LIM,DSORT,DFOAT,DEXP,DLOG,MMBSJO,MMBSJ1
C      INTEGER DIGIT(10)/'1','2','3','4','5','6','7','8','9','10'/
C      INTEGER VF1(7)/'(5X,'X-7M','n' // ' = ', // '(6X,'12)')'/
C      INTEGER VF2(8)/'(5X,'7MTH','eta' // ' = 3X' // 'F8.2',')'/
C      DIMENSION VN(10),Z(100),BJOVN(10),YOVN(10),AN(10),EXPT(10)
C      DIMENSION ALVN(10),AE(10),THETA(10),ZH(100),UE(100),YVN(2),VV(2)
C
C *****
C
C      * Read all input data
C
C *****
C
C      READ(5,50)T,SIGMA0,C,H,K
C      READ(5,51)NOEL,NORT,UO,GAMMA,GAMMAW
C      READ(5,52)(VN(I),I=1,NORT)
C
C *****
C
C      * Variable format input
C
C *****
C
C      VF1(5)=DIGIT(NORT)
C      VF2(6)=DIGIT(NORT)
C
C *****
C
C      * Compute the average consolidation at the given time T
C      * and the % error at time 0 The theta terms are also
C      * calculated
C
C *****
C
C      ALPHA=SIGMA0/(C*GAMMA)
C      HA1=H/ALPHA+1 DO
C      LAMBDA=1 DO/DSORT(HA1)
C      TT=C*K=GAMMA*T/GAMMAW
C
C      DHT=0. DO
C      CONV=0 DO
C
C      DO 20 I=1,NORT
C      ALVN(I)=LAMBDA*VN(I)
C      BJOVN(I)=MMBSJO(VN(I),IEJ1)
C      BJOLVN=MMBSJO(ALVN(I),IEJ2)
C      BJ1VN=MMBSJ1(VN(I),IEJ3)
C      CALL MMBSYN(VN(I),0. DO,2,YVN,IEY1)
C      CALL MMBSYN(ALVN(I),0. DO,1,YOLVN,IEY2)
C      YOYN(I)=YVN(1)
C      Y1VN=YVN(2)
C      ZOLVN=BJOLVN-BJOVN(I)*YOLVN/YOYN(I)
C      Z1VN=BJ1VN-BJOVN(I)*Y1VN/YOYN(I)
C      AN(I)=2 DO*UO*ZOLVN/(YVN(I)*Z1VN**2*(ALVN(I)*ZOLVN**2)
C      THETA(I)=0.25DO*ALVN(I)*ALVN(I)*TT/ALPHA
C      EXPT(I)=DEXP(-THETA(I))
C      AE(I)=AN(I)*EXPT(I)

```





```

      DHT=DHT+AE(I)*ZOLVN
      CONV=CONV+AN(I)*ZOLVN/(2.DO*UO)
20  CONTINUE
C
      LIM=DLOG(1.DO/LAMBDA)/2.DO
      ERR=(CONV-LIM)*100/LIM
C
      TEM=2.DO/(GAMMA*C)
      DELHF=TEM*UO*LIM*2.DO
      DELHT=DELHF-TEM*DHT
C
      AVCON=DELHT/DELHF
C
      *****
      *
      *   Calculation of the excess pore pressure at various depth
      *
      *****
C
      NONOD=NOEL+1
      DIV=H/DFLOAT(NOEL)
      Z(1)=0.DO
C
      DO 25 J=1,NONOD
      Z(J+1)=Z(J)+DIV
      ZH(J)=Z(J)/H
      ZA1=Z(J)/ALPHA+1.DO
      UE(J)=0.DO
C
      DO 25 L=1,NORT
      V=ALVN(L)*DSORT(ZA1)
      BJ1V=MMBSJ1(V,IEJ4)
      CALL MMBSYN(V,0.DO,2,YV,IEY3)
      Y1V=YV(2)
      Z1V=BJ1V-BJOVN(L)*Y1V/YOVN(L)
      UE(J)=UE(J)+AE(L)*V+Z1V
25  CONTINUE
C
      *****
      *
      *   Output quantities
      *
      *****
C
      WRITE(6,60)
      WRITE(6,61) AVCON,ERR
      WRITE(6,65) K,C,SIGMAO,H,GAMMA,GAMMAW,UO,ALPHA,LAMBDA,NORT,NOEL
      WRITE(6,62) T,(Z(I),ZH(I),I,UE(I),I=1,NONOD)
      WRITE(6,63) IEJ1,IEJ2,IEJ3,IEJ4,IEY1,IEY2,IEY3
      WRITE(6,64)
      WRITE(6,VF1)(N,N=1,NORT)
      WRITE(6,VF2)(THETA(J),J=1,NORT)
C
      STOP
C
      *****
      *
      *   Format statments
      *
      *****
C
      50 FORMAT(5G20.0)
      51 FORMAT(2I5,3G20.0)
      52 FORMAT(3F20.0)
C
      60 FORMAT('1'////4X'The One Dimensional Consolidation with ',
     1'the Spatial Variation'/23X'of Compressibility'
     1//13X'Analytical Solution by Edelman (1953)'////)
      62 FORMAT(13X'Excess pore pressure vs depth:'//
     113X'Depth',6X'Depth',4X'NODE',9X'Time t'/24X'Ratio',10XG20.9
     1//((9XF11.5,F10.4,3X13,4XE15.8))
      61 FORMAT(4X'Average Consolidation =',F12.8
     1//4X'(Error of ave consol at time 0 =',F8.1,' %)'////)
      63 FORMAT(////5X'Error message from the *IMSLDPLIB for the final'
     1,' iteration:'//5X'IEJ1 IEJ2 IEJ3 IEJ4 IEY1 IEY2 IEY3'/
     12X7(2XI4))
      64 FORMAT(////5X'The theta terms:')
      65 FORMAT(4X'K          =',E15.6/4X'C          =',F15.5/
     14X'SIGMAO =',F15.5/4X'H          =',F15.5/
     14X'GAMMA  =',F15.5/4X'GAMMAW =',F15.5/
     14X'UO      =',F15.5//4X'Alpha =',F15.5/
     14X'Lambda =',F15.5/4X'No of roots used =',I5/
     14X'No of elements  =',I5//)
C
      *****
      *
      *   # IMSL Library #
      *
      *   The following is a brief description of the Fortran
      *   routines used in this program. This information is
      *   available in the IMSL Library reference manual, edition
      *   8, 1980. The routines with single precision are stored
      *   in the public file *IMSLLIB, those with double precision
      *   are stored in *IMSLDPLIB.
      *
      *
      *   Routine name:  MMBSJO
      *
      *   Usage:  FUNCTION MMBSJO(ARG,IER)
      *
      *   Description of arguments:
      *   ARG - Input argument
      *   IER - Error parameter, terminal error
      *
      *
      *   Routine name:  MMBSJ1

```



```
C      *
C      * Usage: FUNCTION MMBSJ1(ARG,IER)
C      *
C      * Description of arguments:
C      * ARG - Input argument
C      * IER - Error parameter, terminal error
C      *
C      *
C      * Routine name: MMBSYN
C      *
C      * Usage: CALL MMBSYN(ARG,ORDER,N,YN,IER)
C      *
C      * Description of arguments:
C      * ARG - Input argument
C      * ORDER - Desired order, which must be greater than or
C      *          equal to zero and less than one
C      * N      - Number of results required, with which order
C      *          start from 'ORDER' to 'ORDER+1' to 'ORDER+2', etc
C      * YN     - Output vector of above results
C      * IER    - Error parameter where
C      *          '129': ORDER is out of range or ARG is less than
C      *                the minimum allowable value
C      *          '130': N is less than or equal to zero
C      *          '131': YN overflow
C      *          '132': ARG is greater than the maximum allowable
C      *                value
C      *
C      *****
C
END
```



## B.5 Datafile A.20: Eq. A4

```

C *****
C *
C *      #This program is to prepare part of the data for A.20(Eq.A4)
C *      according to ADINAT users manual input section X.3 & X.4
C * Variables: CC      = Cc/ln10
C *              DIVU   = load increment / number of material input points
C *              E       = void ratio
C *              EO      = initial void ratio
C *              K       = coefficient of permeability, k;
C *                      or variable of "gamma*mv", where mv is the
C *                      coefficient of volume compressibility
C *              P       = effective stress
C *              PO      = initial effective stress
C *              U       = excess pore pressure
C *              XL      = reduced depth, computed from Eq. 6.3 in text
C * To use:  RUN *FORTGTEST SCARDS=file SPRINT=-1 PAR=id,source,map
C *          DEBUG -LOAD# 5=data 6=output
C *****
C IMPLICIT REAL*8(A-H,O-Z)
C REAL*8 DLOG,DLOG10,K(500)
C DIMENSION U(500),XL(500)
C A=1.DO
C XL(1)=.1419DO
C XL(2)=.4592DO
C XL(3)=.8025DO
C XL(4)=1.1640DO
C XL(5)=1.5399DO
C XL(6)=1.9281DO
C XL(7)=2.3270DO
C XL(8)=2.7356DO
C XL(9)=3.1530DO
C XL(10)=3.5785DO
C XL(11)=4.0116DO
C XL(12)=4.4518DO
C XL(13)=4.8987DO
C XL(14)=5.3521DO
C XL(15)=5.8115DO
C XL(16)=6.2769DO
C XL(17)=6.7478DO
C XL(18)=7.2250DO
C XL(19)=7.7059DO
C XL(20)=8.1926DO
C CC=.4DO/DLOG(10.DO)
C DIVU=10.DO/20.DO
C U(1)=0.DO
C DO10II=1,20
10 U(II+1)=U(II)+DIVU
C DO20I=1,20
C PO=.1DO+1.9DO*XL(I)
C EO=.9DO-.4DO*DLOG10(10.DO*PO)
C EO1=1.DO+EO
C DO30J=1,21
C P=PO+10.DO-U(J)
C K(J)=CC/(P*EO1)
30 CONTINUE
C WRITE(6,60)I,A
C WRITE(6,61)(U(LL),LL=1,21),(K(JL),JL=1,21)
C WRITE(6,61)(U(LL),LL=1,21),(K(JL),JL=1,21)
20 CONTINUE
C STOP
60 FORMAT(3X,I2,F10.0)
61 FORMAT(2(8F10.1/),5F10.1,3F10.9/3(8F10.9/))
C END

```



## B.6 Datafile A.20: Eq. A8

```

C *****
C *
C *      #This program is to prepare part of the data for A.20(Eq.A8)
C *      according to ADINAT users manual input section X.3 & X.4
C * Variables: CC      = Cc/ln10
C *              DIVU   = load increment / number of material input points
C *              E      = void ratio
C *              EO     = initial void ratio
C *              K      = variable of "k*(1+eo) / 1+e"
C *              MV     = variable of "gamma*mv*(1+e) / 1+eo"
C *              P      = effective stress
C *              PO     = initial effective stress
C *              U      = excess pore pressure
C *              XL     = reduced depth, computed from Eq. 6.3 in text
C * To use:  RUN *FORTGTEST SCARDS=file SPRINT=-1 PAR=id,source,map
C *          DEBUG -LOAD# 5=data 6=output
C *
C *****
IMPLICIT REAL*8(A-H,O-Z)
REAL*8 DLOG,DLOG10,K(500),MV(500)
DIMENSION U(500),XL(500)
A=1.DO
XL(1)=.1419DO
XL(2)=.4592DO
XL(3)=.8025DO
XL(4)=1.1640DO
XL(5)=1.5399DO
XL(6)=1.9281DO
XL(7)=2.3270DO
XL(8)=2.7356DO
XL(9)=3.1530DO
XL(10)=3.5785DO
XL(11)=4.0116DO
XL(12)=4.4518DO
XL(13)=4.8987DO
XL(14)=5.3521DO
XL(15)=5.8115DO
XL(16)=6.2769DO
XL(17)=6.7478DO
XL(18)=7.2250DO
XL(19)=7.7059DO
XL(20)=8.1926DO
CC=.4DO/DLOG(10.DO)
DIVU=10.DO/20.DO
U(1)=0.DO
D01OII=1,20
10 U(II+1)=U(II)+DIVU
D02OI=1,20
PO=.1DO+1.9DO*XL(I)
EO=.9DO-.4DO*DLOG10(10*PO)
EO1=1.DO+EO
D03OJ=1,21
P=PO+10.DO-U(J)
E=.9DO-.4DO*DLOG10(10*P)
E1=1.DO+E
MV(J)=CC/(P*EO1)
K(J)=CC*EO1/(P*E1*E1)
30 CONTINUE
WRITE(6,60)I,A
WRITE(6,61)(U(LL),LL=1,21),(K(JL),JL=1,21)
WRITE(6,62)(U(LL),LL=1,21),(MV(JL),JL=1,21)
20 CONTINUE
STOP
60 FORMAT(3X,I2,F10.0)
61 FORMAT(2(8F10.1/),5F10.1,3F10.9/3(8F10.9/))
62 FORMAT(2(8F10.1/),5F10.1,3F10.9/3(8F10.9/))
END

```





## B.7 Datafile A.20: Eq. A9

```

C *****
C *
C *      #This program is to prepare part of the data for A.20(Eq.A9)
C *      according to ADINAT users manual input section X.3 & X.4
C * Variables: CC   = Cc/ln10
C *              DIVU = load increment / number of interpolation points
C *              DIVZ = reduced layer depth / number of elements
C *              E    = void ratio, e
C *              K    = variable of "k / 1+e"
C *              MV   = variable of "gamma*mv*(1+e)"
C *              P    = effective stress
C *              U    = excess pore pressure
C *              Z    = reduced coordinate, 'b'
C * To use:  RUN *FORTGTEST SCARDS=file SPRINT=-1 PAR=id,source,map
C *          DEBUG -LOAD# 5=data 6=output
C *****
C IMPLICIT REAL*8(A-H,O-Z)
C REAL*8 DLOG,DLOG10,K(500),MV(500)
C DIMENSION U(500)
C A=1.DO
C DIVU=10.DO/20.DO
C U(1)=0.DO
C DO10II=1,20
10  U(II+1)=U(II)+DIVU
C CC=.4DO/DLOG(10.DO)
C DIVZ=8.43786DO/20.DO
C Z=DIVZ/2.DO
C DO20I=1,20
C DO30J=1,21
C P=10.1DO+1.9DO*Z-U(J)
C E=.9DO-.4DO*DLOG10(10.DO*P)
C E1=1.DO+E
C MV(J)=CC/P
C K(J)=MV(J)/(E1*E1)
30  CONTINUE
C WRITE(6,60)I,A
C WRITE(6,62)(U(LL),LL=1,21),(K(JJ),JJ=1,21)
C WRITE(6,61)(U(LL),LL=1,21),(MV(JJ),JJ=1,21)
C Z=Z+DIVZ
20  CONTINUE
C STOP
60  FORMAT(3X,I2,F10.0)
61  FORMAT(2(8F10.1/),5F10.1,3F10.9/3(8F10.9/))
62  FORMAT(2(8F10.1/),5F10.1,3F10.9/3(8F10.9/))
C END

```



## B.8 Datafile A.21: Eq. A4

```

C *****
C *
C *      #This program is to prepare part of the data for A.21(Eq.A4)
C *      according to ADINAT users manual input section X.3 & X.4
C * Variables: DIVU = load increment / number of material input points
C *               E   = void ratio
C *               EO  = initial void ratio
C *               K   = k ( = gamma*mv*Cv )
C *               MV  = variable of "gamma*mv"
C *               P   = effective stress
C *               POO = variable of "-lambda * initial effective stress"
C *               U   = excess pore pressure
C *               XL  = reduced depth: Eq.31 from Gibson etal (1981)
C * To use:  RUN *FORTGTEST SCARDS=file SPRINT=-1 PAR=id,source,map
C *          DEBUG -LOAD# 5=data 6=output
C *****
IMPLICIT REAL*8(A-H,O-Z)
REAL*8 DEXP,K(500),MV(500),XL(500)
DIMENSION U(500)
A=1.DO
XL(1)=.0521DO
XL(2)=.1587DO
XL(3)=.2684DO
XL(4)=.3813DO
XL(5)=.4974DO
XL(6)=.6168DO
XL(7)=.7396DO
XL(8)=.8658DO
XL(9)=.995DO
XL(10)=1.128DO
XL(11)=1.265DO
XL(12)=1.405DO
XL(13)=1.548DO
XL(14)=1.695DO
XL(15)=1.845DO
XL(16)=1.998DO
XL(17)=2.155DO
XL(18)=2.315DO
XL(19)=2.477DO
XL(20)=2.642DO
DIVU=10.DO/20.DO
U(1)=0.DO
DO10II=1,20
10 U(II+1)=U(II)+DIVU
DO20I=1,20
POO=-.332DO*1.74DO*XL(I)
EO=1.49DO+2.34DO*DEXP(POO)
EO1=1.DO+EO
DO30J=1,21
P=-.332DO*(1.74DO*XL(I)+10.DO-U(J))
E=1.49DO+2.34DO*DEXP(P)
AM=.332DO*(E-1.49DO)
MV(J)=AM/EO1
K(J)=MV(J)*.001944DO
30 CONTINUE
WRITE(6,60)I,A
WRITE(6,61)(U(LL),LL=1,21),(K(JL),JL=1,21)
WRITE(6,62)(U(LL),LL=1,21),(MV(JL),JL=1,21)
20 CONTINUE
STOP
60 FORMAT(3X,I2,F10.0)
61 FORMAT(2(8F10.1/),5F10.1,3E11.6/3(8E11.6/))
62 FORMAT(2(8F10.1/),5F10.1,3F10.9/3(8F10.9/))
END

```



## B.9 Datafile A.21: Eq. A8

```

C *****
C *
C *      #This program is to prepare part of the data for A.21(Eq.A8)
C *      according to ADINAT users manual input section X.3 & X.4
C *      Variables: DIVU = load increment / number of material input points
C *                  E    = void ratio
C *                  EO   = initial void ratio
C *                  K    = variable of "k*(1+eo) / 1+e"
C *                  MV   = variable of "gamma*mv*(1+e) / (1+eo)"
C *                  P    = effective stress
C *                  POO  = variable of "-lambda * initial effective stress"
C *                  U    = excess pore pressure
C *                  XL   = reduced depth: Eq.31 from Gibson etal (1981)
C *      To use:  RUN *FORTGTEST SCARDS=file SPRINT=-1 PAR=id,source,map
C *              DEBUG -LOAD# 5=data 6=output
C *****
IMPLICIT REAL*8(A-H,O-Z)
REAL*8 DEXP,K(500),MV(500),XL(500)
DIMENSION U(500)
A=1.D0
XL(1)=.0521D0
XL(2)=.1587D0
XL(3)=.2684D0
XL(4)=.3813D0
XL(5)=.4974D0
XL(6)=.6168D0
XL(7)=.7396D0
XL(8)=.8658D0
XL(9)=0.995D0
XL(10)=1.128D0
XL(11)=1.265D0
XL(12)=1.405D0
XL(13)=1.548D0
XL(14)=1.695D0
XL(15)=1.845D0
XL(16)=1.998D0
XL(17)=2.155D0
XL(18)=2.315D0
XL(19)=2.477D0
XL(20)=2.642D0
DIVU=10.D0/20.D0
U(1)=0.D0
DO10II=1,20
10 U(II+1)=U(II)+DIVU
DO20I=1,20
POO=-.332D0*1.74D0*XL(I)
EO=1.49D0+2.34D0*DEXP(POO)
EO1=1.D0+EO
DO30J=1,21
P=-.332D0*(1.74D0*XL(I)+10.D0-U(J))
E=1.49D0+2.34D0*DEXP(P)
E1=1.D0+E
AM=.332D0*(E-1.49D0)
MV(J)=AM/EO1
K(J)=AM*.001944D0*EO1/(E1*E1)
30 CONTINUE
WRITE(6,60)I,A
WRITE(6,61)(U(LL),LL=1,21),(K(JL),JL=1,21)
WRITE(6,62)(U(LL),LL=1,21),(MV(JL),JL=1,21)
20 CONTINUE
STOP
60 FORMAT(3X,I2,F10.0)
61 FORMAT(2(8F10.1/),5F10.1,3E11.6/3(8E11.6/))
62 FORMAT(2(8F10.1/),5F10.1,3F10.9/3(8F10.9/))
END

```



## B.10 Datafile A.21: Eq. A9

```

C *****
C *
C *      # This program is to prepare part of the data for A.21(Eq.A9) *
C *      according to ADINAT users manual input section X.3 & X.4 *
C * Variables: DIVU = load increment / number of interpolation points *
C *              DIVZ = reduced layer depth / number of elements *
C *              E    = void ratio, e *
C *              K    = variable of "k / 1+e" *
C *              MV   = variable of "gamma*mv*(1+e)" *
C *              P    = effective stress *
C *              U    = excess pore pressure *
C *              Z    = reduced coordinate, 'b' *
C * To use:  RUN *FORTGTEST SCARDS=file SPRINT=-1 PAR=id,source,map *
C *          DEBUG -LOAD# 5=data 6=output *
C *****
IMPLICIT REAL*8(A-H,O-Z)
REAL*8 DEXP,K(500),MV(500)
DIMENSION U(500)
A=1.DO
DIVU=10.DO/20.DO
U(1)=0.DO
D01OI=1,20
10 U(II+1)=U(II)+DIVU
DIVZ=2.728DO/20.DO
Z=DIVZ/2.DO
D02OI=1,20
D03OJ=1,21
P=-.332DO*(1.74DO*Z+10.DO-U(J))
E=1.49DO+2.34DO*DEXP(P)
E1=1.DO+E
MV(J)=.332DO*(E-1.49DO)
K(J)=MV(J)*.001944DO/(E1*E1)
30 CONTINUE
WRITE(6,60)I,A
WRITE(6,61)(U(LL),LL=1,21),(K(JJ),JJ=1,21)
WRITE(6,62)(U(LL),LL=1,21),(MV(JJ),JJ=1,21)
Z=Z+DIVZ
20 CONTINUE
STOP
60 FORMAT(3X,I2,F10.0)
61 FORMAT(2(8F10.2/),5F10.2,3E11.6/3(8E11.6/))
62 FORMAT(2(8F10.2/),5F10.2,3F10.9/3(8F10.9/))
END

```

















**B30422**